

Describing Gluons

Axel Maas

28th of July 2009
Pathways to Confinement
Rio de Janeiro
Brazil

Overview

- *Gauge freedom* in Yang-Mills theory

Supported by the FWF

Overview

- Gauge freedom in Yang-Mills theory
- Non-perturbative gauges

Supported by the FWF

Overview

- Gauge freedom in Yang-Mills theory
- Non-perturbative gauges
- Comparing lattice and continuum

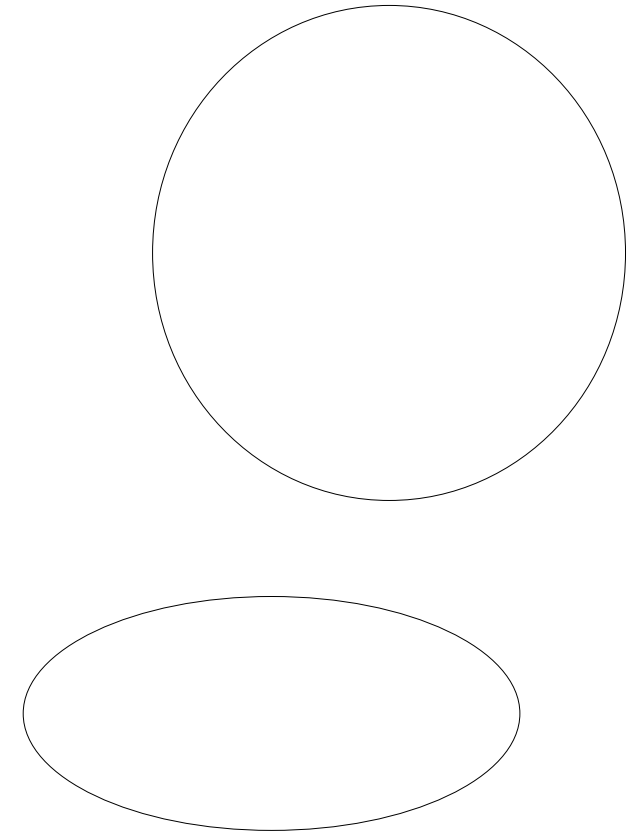
Supported by the FWF

Overview

- Gauge freedom in Yang-Mills theory
- Non-perturbative gauges
- Comparing lattice and continuum
- Correlation functions and physics

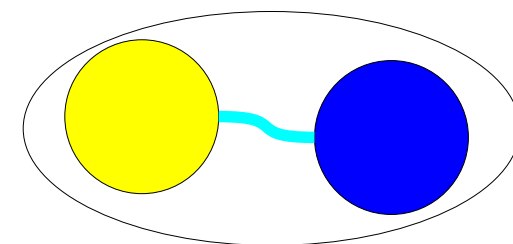
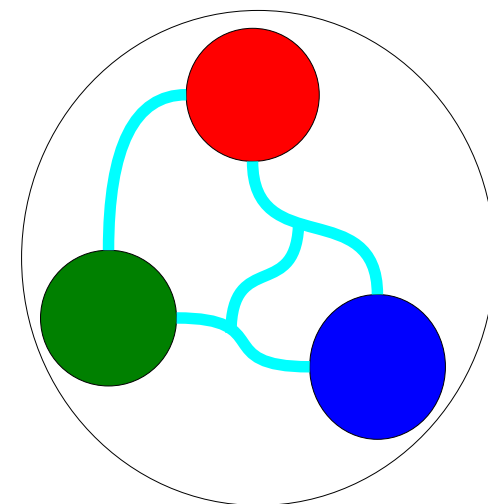
Supported by the FWF

Strong interactions



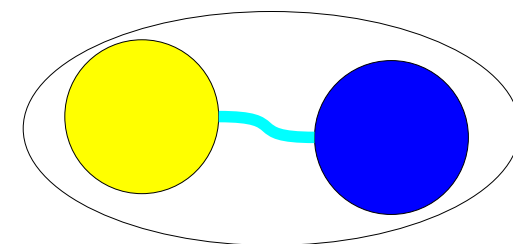
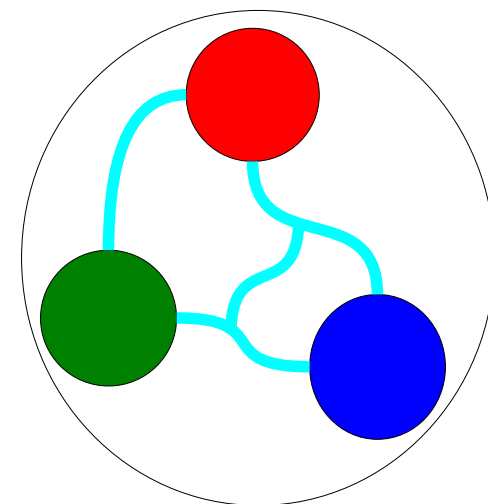
Strong interactions

- Substructure can be described by QCD
 - Degrees of freedom are quarks and gluons



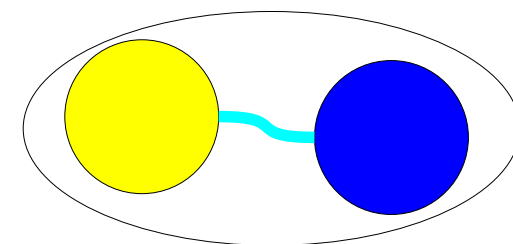
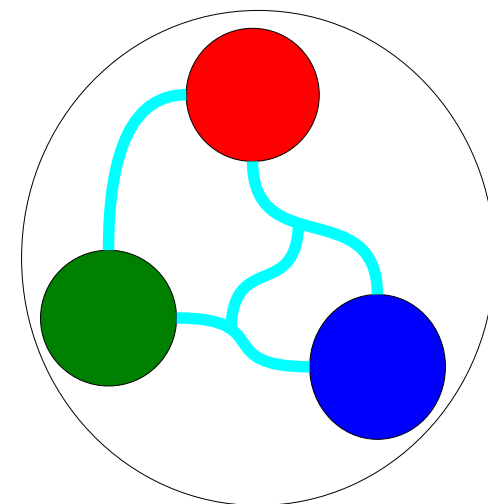
Strong interactions

- Substructure can be described by QCD
 - Degrees of freedom are quarks and gluons
 - Color charge carried by both
 - Confinement: Colored objects unobservable



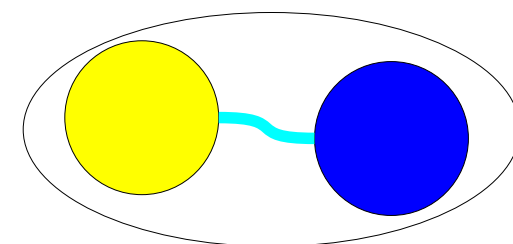
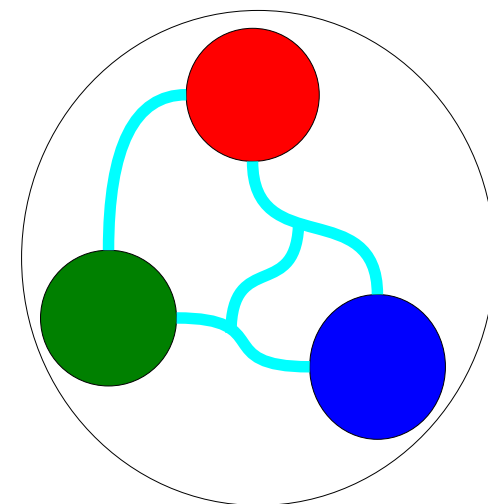
Strong interactions

- Substructure can be described by QCD
 - Degrees of freedom are quarks and gluons
 - Color charge carried by both
 - Confinement: Colored objects unobservable
 - Gauge theory – like QED



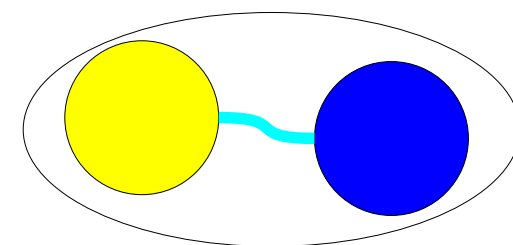
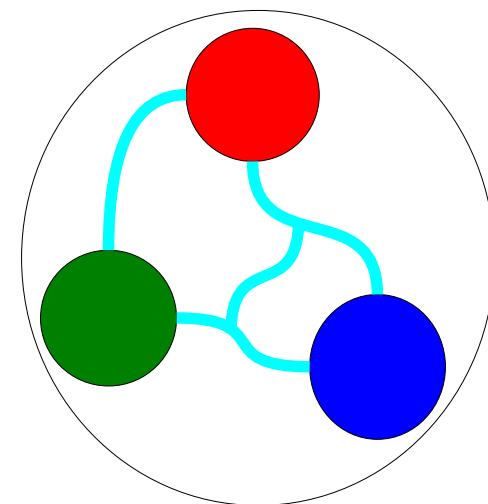
Strong interactions

- Substructure can be described by QCD
 - Degrees of freedom are quarks and gluons
 - Color charge carried by both
 - Confinement: Colored objects unobservable
 - Gauge theory – like QED
 - Color charges in QCD not gauge-invariant
 - Different from QED
 - No concept of a local gauge-invariant color charge distribution
 - Similar to energy density in general relativity



Strong interactions

- Substructure can be described by QCD
 - Degrees of freedom are quarks and gluons
 - Color charge carried by both
 - Confinement: Colored objects unobservable
 - Gauge theory – like QED
 - Color charges in QCD not gauge-invariant
 - Different from QED
 - No concept of a local gauge-invariant color charge distribution
 - Similar to energy density in general relativity
- Reduce complexity and ignore Quarks: Yang-Mills theory
 - Affects gluon properties only quantitatively



Yang-Mills Theory

- Lagrangian (after Wick rotation to Euclidean space-time):

$$L = -\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - gf^{abc} A_\mu^b A_\nu^c$$

- Degrees of freedom:

Gluons: A_μ^a

- g is the coupling constant, giving the strength of coupling
- f^{abc} are numbers, depending on the **gauge group**, $SU(3)$ for QCD:
quarks & gluon are organized in multiplets, just as with spin

Gauge-fixing

- Yang-Mills theory is a **gauge theory**

- **Gauge transformations** $A_\mu^a \rightarrow A_\mu^a + (\delta^{ab} \partial_\mu - g f^{abc} A_\mu^c) \phi^b(x)$

with arbitrary $\phi^a(x)$ change the gauge fields, but leave physics invariant

Gauge-fixing

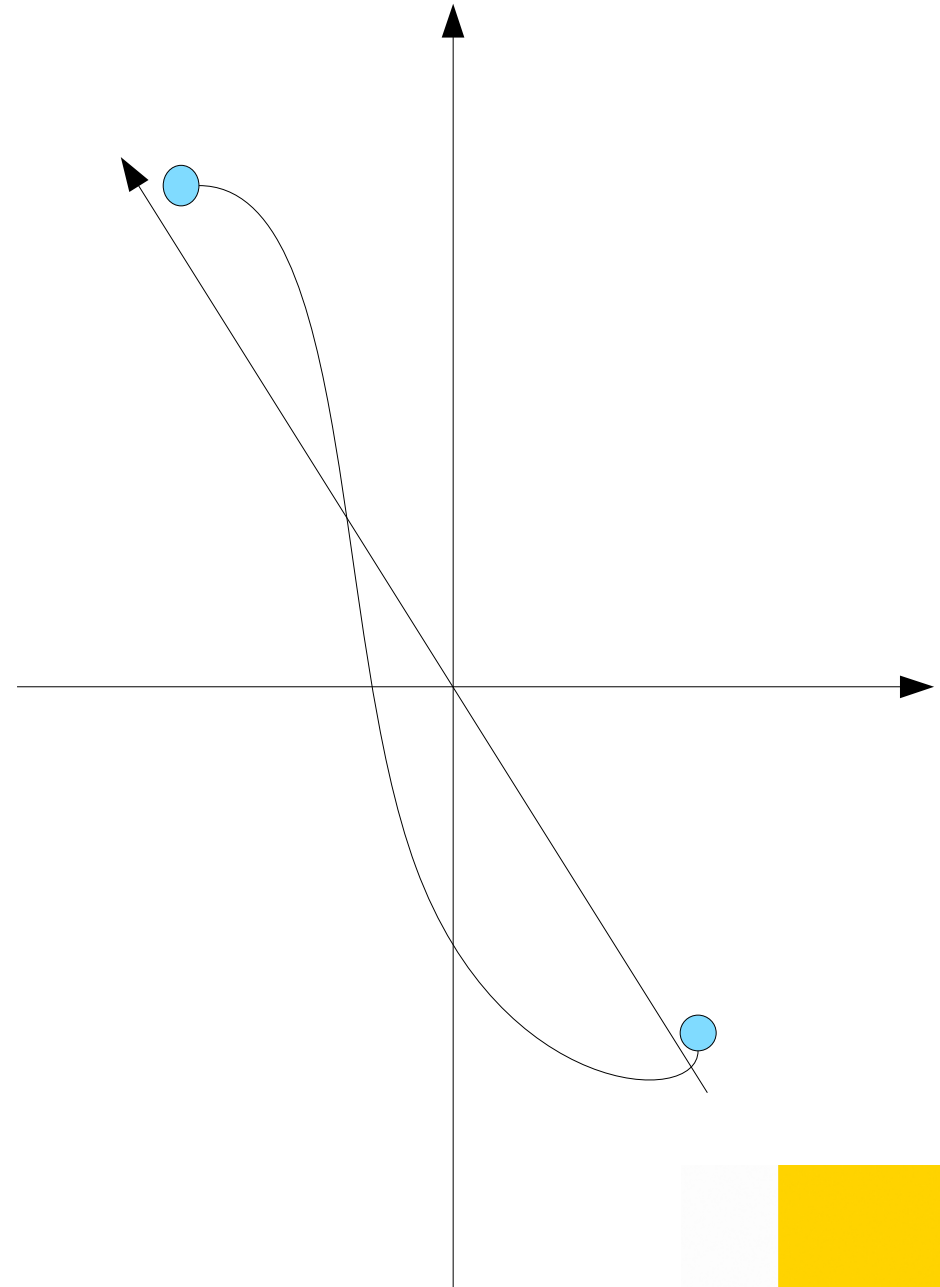
- Yang-Mills theory is a **gauge theory**
 - **Gauge transformations** $A_\mu^a \rightarrow A_\mu^a + (\delta^{ab} \partial_\mu - g f^{abc} A_\mu^c) \phi^b(x)$
with arbitrary $\phi^a(x)$ change the gauge fields, but leave physics invariant
- **Correlation functions** are in general **gauge-dependent**
 - **Gauge-fixing** is required

Gauge-fixing

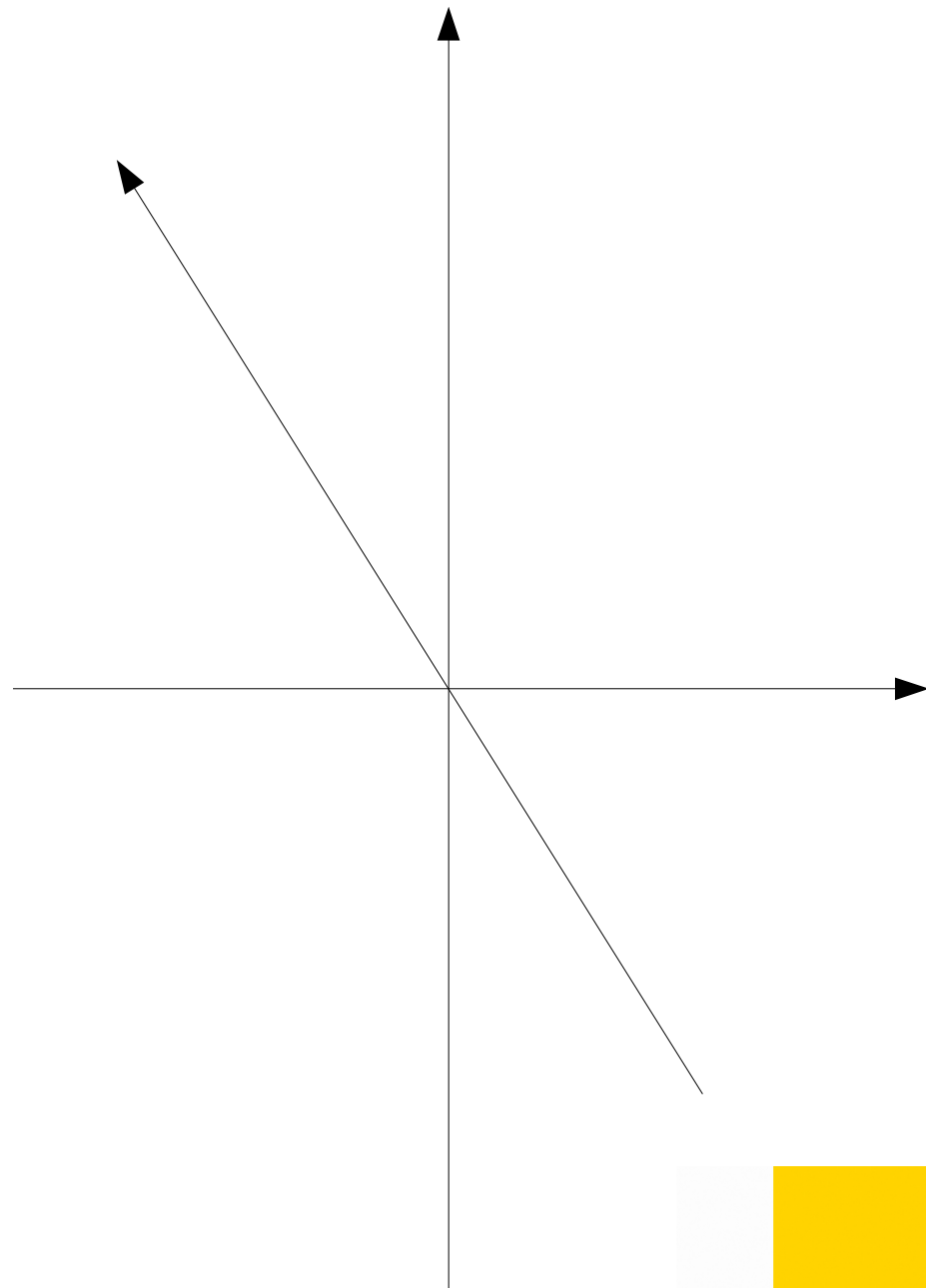
- Yang-Mills theory is a **gauge theory**
 - **Gauge transformations** $A_\mu^a \rightarrow A_\mu^a + (\delta^{ab} \partial_\mu - g f^{abc} A_\mu^c) \phi^b(x)$
with arbitrary $\phi^a(x)$ change the gauge fields, but leave physics invariant
- **Correlation functions** are in general **gauge-dependent**
 - **Gauge-fixing is required**
- **Example: Landau gauge condition** $\partial_\mu A_\mu^a = 0$
 - Here only Landau gauge results will be considered
 - Many other gauges have been studied

Configuration space (artist's view)

- Gauge fields not unique
 - Gauge transformation does not change physics

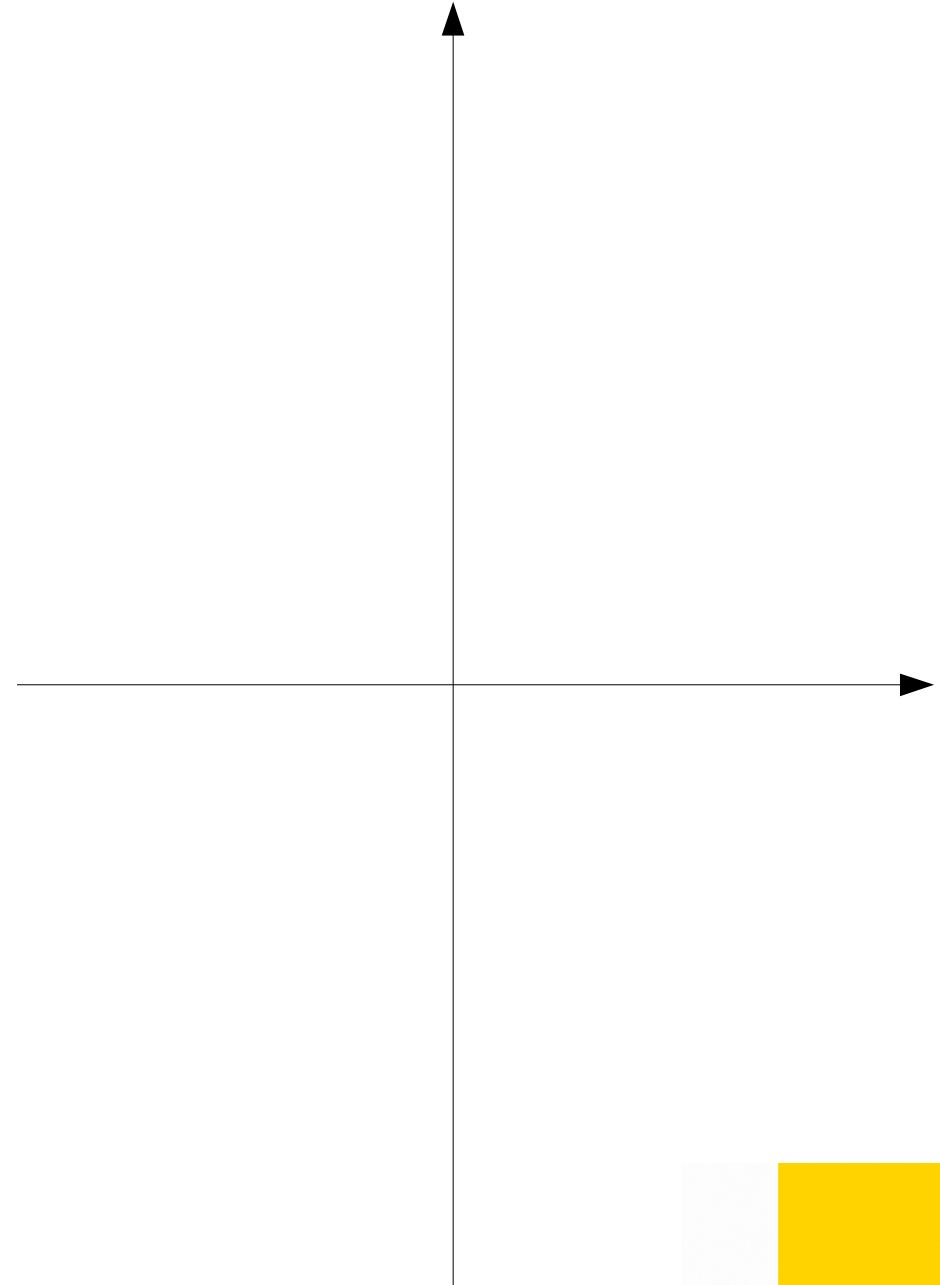


Configuration space (artist's view)



Configuration space (artist's view)

- Impose **Landau gauge** condition
 - Reduces configuration space to a hypersurface



Unique gauge-fixing [For an introduction: Sobreiro & Sorella, 2005]

- Local gauge condition
 - Landau gauge: $\partial_\mu A_\mu^a = 0$

Unique gauge-fixing

[For an introduction: Sobreiro & Sorella, 2005]

- Local gauge condition
 - Landau gauge: $\partial_\mu A_\mu^a = 0$
- This condition can be implemented using auxiliary fields, the so-called ghost fields
 - No physical objects: Pure mathematical convenience

(Perturbative) Landau gauge

- Lagrangian:

$$L = -\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a - \bar{c}^a \partial_\mu D_\mu^{ab} c^b$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - gf^{abc} A_\mu^b A_\nu^c$$

$$D_\mu^{ab} = \delta^{ab} \partial_\mu - gf^{abc} A_\mu^c$$
- Degrees of freedom:
 - Gluons: A_μ^a
 - Ghosts: \bar{c}^a, c^a
- Ghosts interact with gluons: They have to be included

Unique gauge-fixing [For an introduction: Sobreiro & Sorella, 2005]

- Local gauge condition
 - Landau gauge: $\partial_\mu A_\mu^a = 0$
- Sufficient for **perturbation theory**

Proceeding

- Once the gauge is fixed, all kind of (perturbative) calculations can be done
- Use **Green's or correlation functions** as basic entities

Green's Functions or correlation functions [Alkofer & von Smekal PR 2001]

- **Green's functions**, or correlation functions, describe a theory completely

Green's Functions or correlation functions [Alkofer & von Smekal PR 2001]

- **Green's functions**, or correlation functions, describe a theory completely
- Expectation values of a product of field operators
 - Build from the fields, here gluons and ghost
 - E.g.: $\langle \bar{c} c \rangle$

Green's Functions or correlation functions [Alkofer & von Smekal PR 2001]

- **Green's functions**, or correlation functions, describe a theory completely
- Expectation values of a product of field operators
 - Build from the fields, here gluons and ghost
 - E.g.: $\langle \bar{c} c \rangle$
- Full **Green's functions** contain all information
- There are an infinite number of them

Green's Functions or correlation functions [Alkofer & von Smekal PR 2001]

- **Green's functions**, or correlation functions, describe a theory completely
- Expectation values of a product of field operators
 - Build from the fields, here gluons and ghost
 - E.g.: $\langle \bar{c} c \rangle$
- Full **Green's functions** contain all information
- There are an infinite number of them
- If having a non-vanishing color charge they change under **gauge transformation**

Green's Functions or correlation functions [Alkofer & von Smekal PR 2001]

- **Green's functions**, or correlation functions, describe a theory completely
- Expectation values of a product of field operators
 - Build from the fields, here gluons and ghost
 - E.g.: $\langle \bar{c} c \rangle$
- Full **Green's functions** contain all information
- There are an infinite number of them
- If having a non-vanishing color charge they change under **gauge transformation**
- Simplest non-zero **Green's functions: 2-point functions** or **propagators**
 - Expectation values of products of two field operators
 - **1-point functions** vanish

Propagators

- In Landau gauge: Gluon and one auxiliary field: Ghost
- Gluon:

$$D_{\mu\nu}^{ab}(x-y) = \langle A_{\mu}^a(x) A_{\nu}^b(y) \rangle$$

$$D_{\mu\nu}(p) = \left(\delta_{\mu\nu} - \frac{p_{\mu} p_{\nu}}{p^2} \right) \frac{Z(p)}{p^2}$$

Propagators

- In Landau gauge: Gluon and one auxiliary field: Ghost

- Gluon:

$$D_{\mu\nu}^{ab}(x-y) = \langle A_{\mu}^a(x) A_{\nu}^b(y) \rangle$$

$$D_{\mu\nu}(p) = \left(\delta_{\mu\nu} - \frac{p_{\mu} p_{\nu}}{p^2} \right) \frac{Z(p)}{p^2}$$

- Ghost:

$$D_G^{ab}(x-y) = \langle \bar{c}^a(x) c^b(y) \rangle$$

$$D_G(p) = \frac{-G(p)}{p^2}$$

Propagators

- In Landau gauge: Gluon and one auxiliary field: Ghost

- Gluon:

$$D_{\mu\nu}^{ab}(x-y) = \langle A_\mu^a(x) A_\nu^b(y) \rangle$$

$$D_{\mu\nu}(p) = \left(\delta_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right) \frac{Z(p)}{p^2}$$

- Ghost:

$$D_G^{ab}(x-y) = \langle \bar{c}^a(x) c^b(y) \rangle$$

$$D_G(p) = \frac{-G(p)}{p^2}$$

- **Ghost propagator** can be expressed as a gluon operator, the inverse Faddeev-Popov operator

$$D_G^{ab}(x-y) \sim \langle (\partial_\mu D_\mu^{ab})^{-1} \rangle = \langle (\partial_\mu (\delta^{ab} \partial_\mu - g f^{abc} A_\mu^c))^{-1} \rangle$$

Proceeding

- Once the gauge is fixed, all kind of (perturbative) calculations can be done
- Use **Green's or correlation functions** as basic entities
- Combination of **gauge-variant Green's functions** yield gauge-invariant results
 - E.g. scattering cross-sections

Proceeding

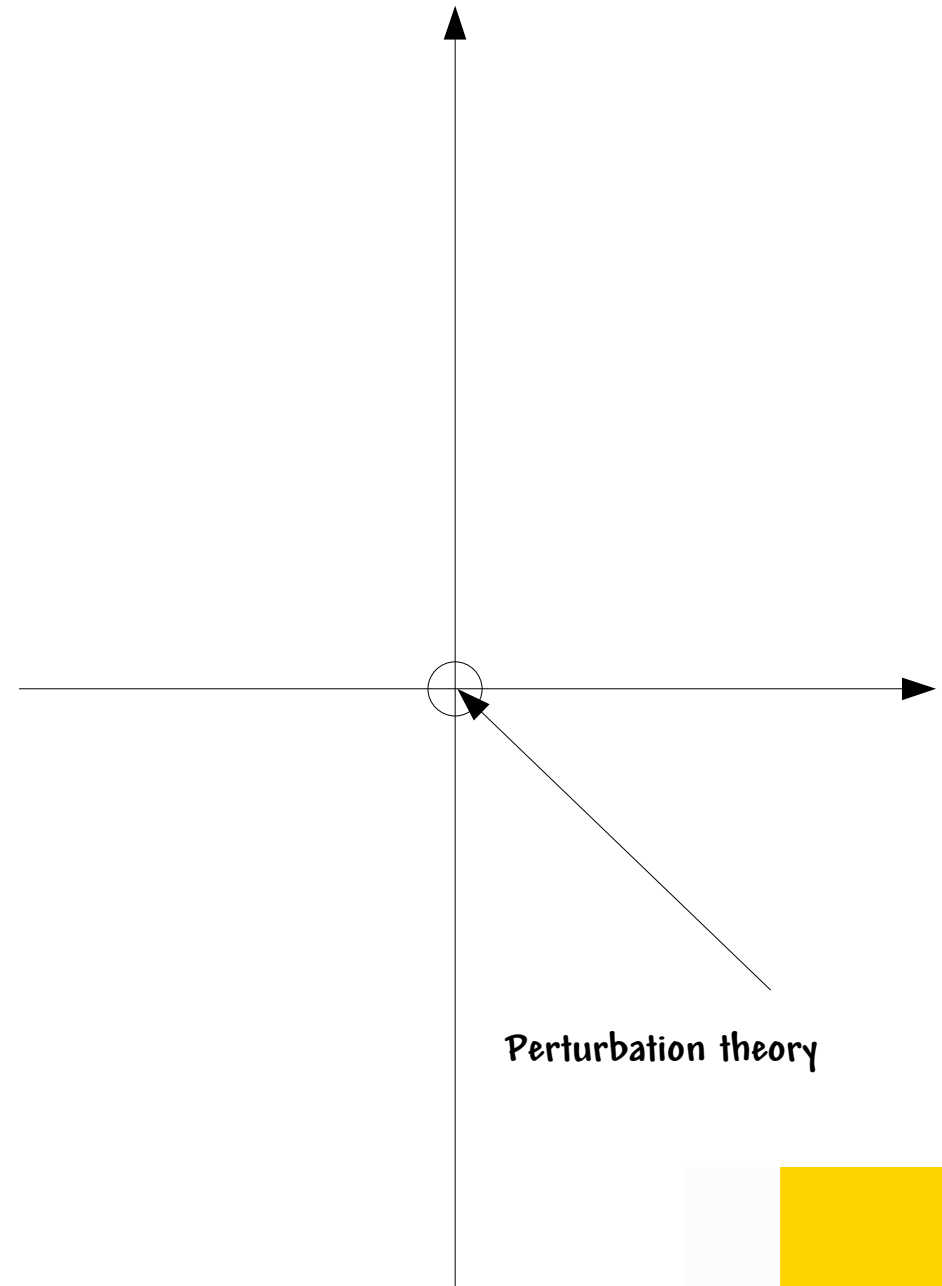
- Once the gauge is fixed, all kind of (perturbative) calculations can be done
- Use **Green's or correlation functions** as basic entities
- Combination of **gauge-variant Green's functions** yield gauge-invariant results
 - E.g. scattering cross-sections
- Almost all **perturbative calculations** proceed via **gauge-variant correlation functions**

Proceeding

- Once the gauge is fixed, all kind of (perturbative) calculations can be done
- Use **Green's or correlation functions** as basic entities
- Combination of **gauge-variant Green's functions** yield gauge-invariant results
 - E.g. scattering cross-sections
- Almost all **perturbative calculations** proceed via **gauge-variant correlation functions**
- Also for the non-perturbative physics?

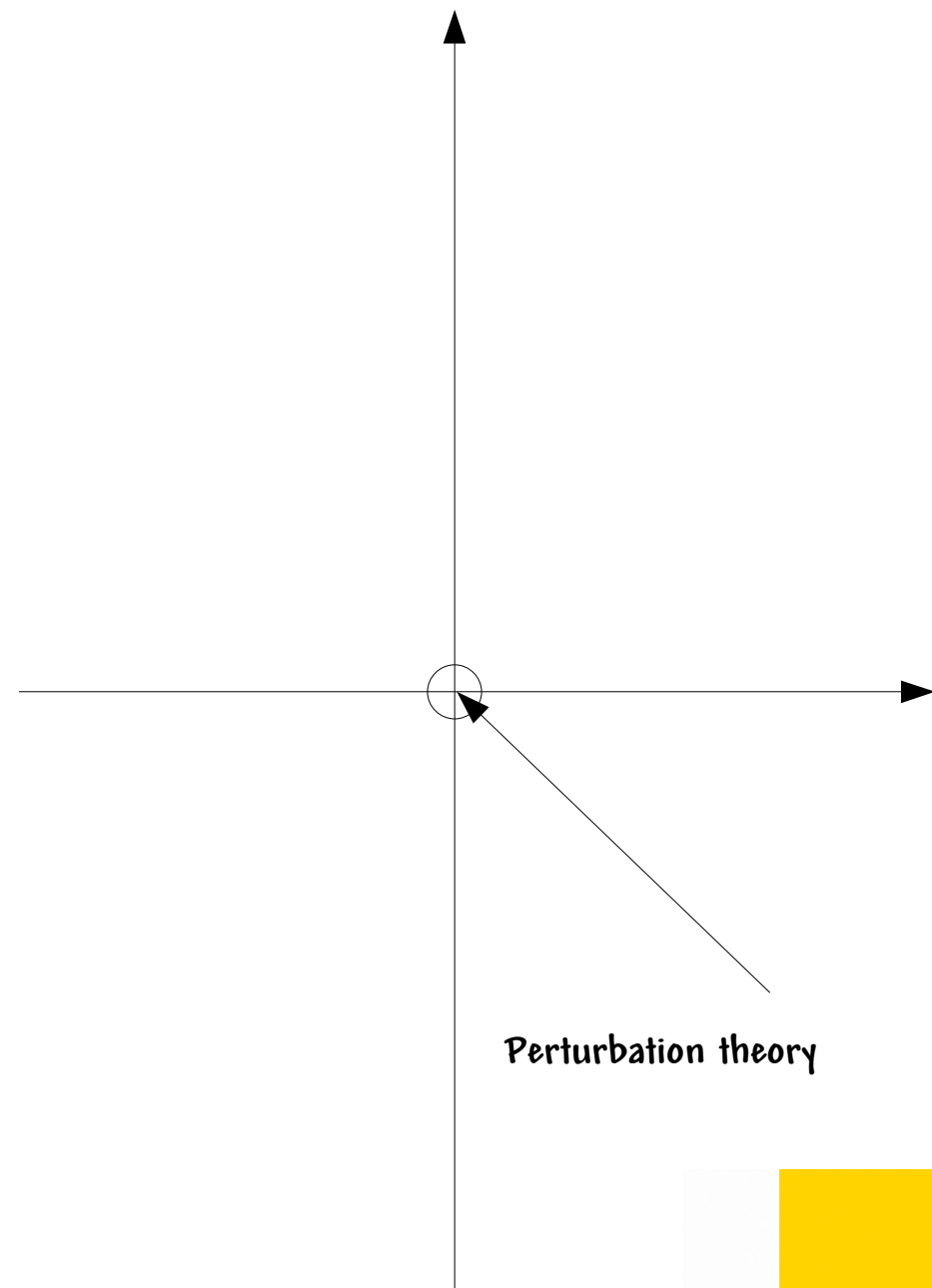
Configuration space (artist's view)

- **Perturbation theory** is applicable close to the origin



Configuration space (artist's view)

- **Perturbation theory** is applicable close to the origin
- Non-perturbative physics probes the complete hypersurface



Unique gauge-fixing [For an introduction: Sobreiro & Sorella, 2005]

- Local gauge condition
 - Landau gauge: $\partial_\mu A_\mu^a = 0$
- Sufficient for perturbation theory
- Insufficient beyond perturbation theory
 - There are gauge-equivalent configurations which obey the same local gauge-condition: Gribov copies [Gribov 1978]

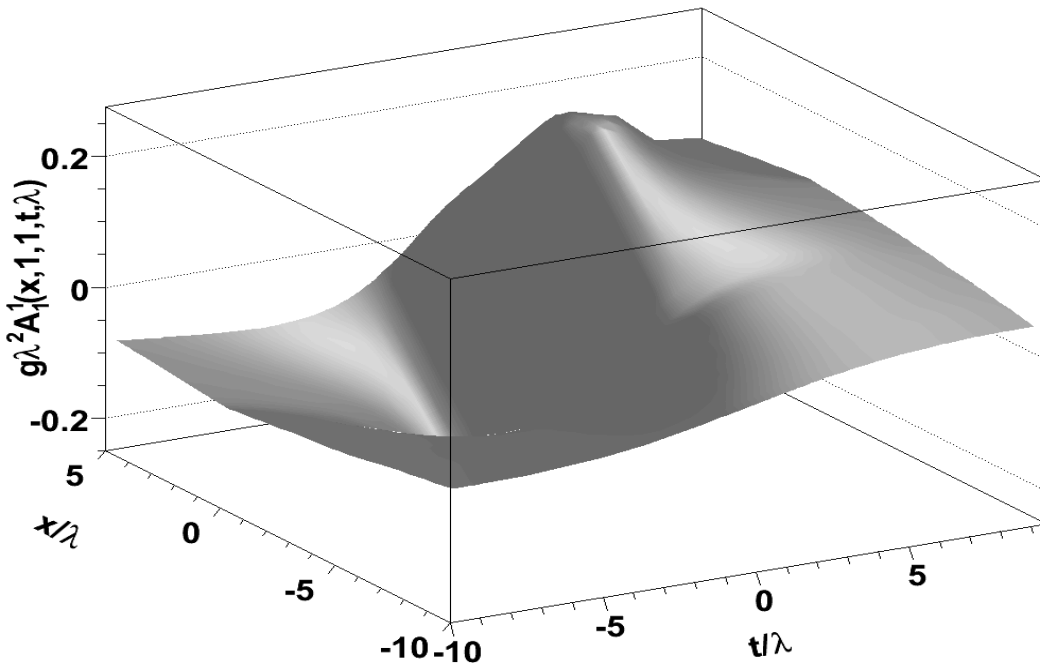
Unique gauge-fixing [For an introduction: Sobreiro & Sorella, 2005]

- Local gauge condition
 - Landau gauge: $\partial_\mu A_\mu^a = 0$
- Sufficient for perturbation theory
- Insufficient beyond perturbation theory
 - There are gauge-equivalent configurations which obey the same local gauge-condition: Gribov copies [Gribov 1978]
- There are no local gauge conditions known, which select a unique gauge field configuration [Singer 1978]
 - Non-local conditions possible

Example: Instanton

[Maas, EPJC 2006]

Instanton field

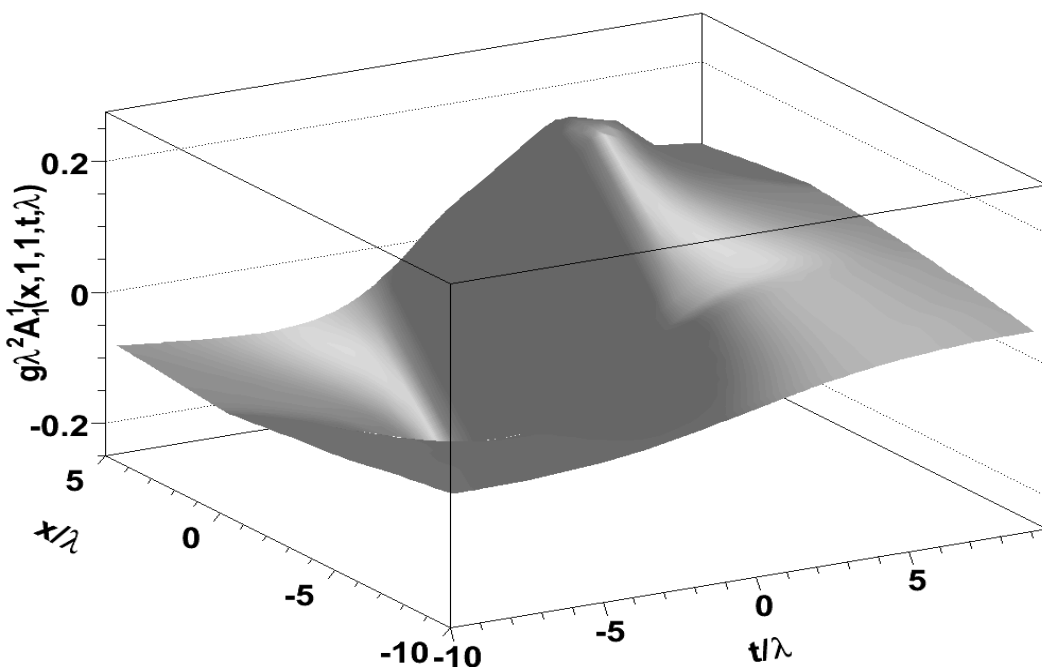


- Instanton field configuration is $A_{\mu}^a(r, \lambda) = 2r_{\nu} \eta_{\nu\mu}^a / (g(r^2 + \lambda^2))$

Example: Instanton

[Maas, EPJC 2006]

Instanton field

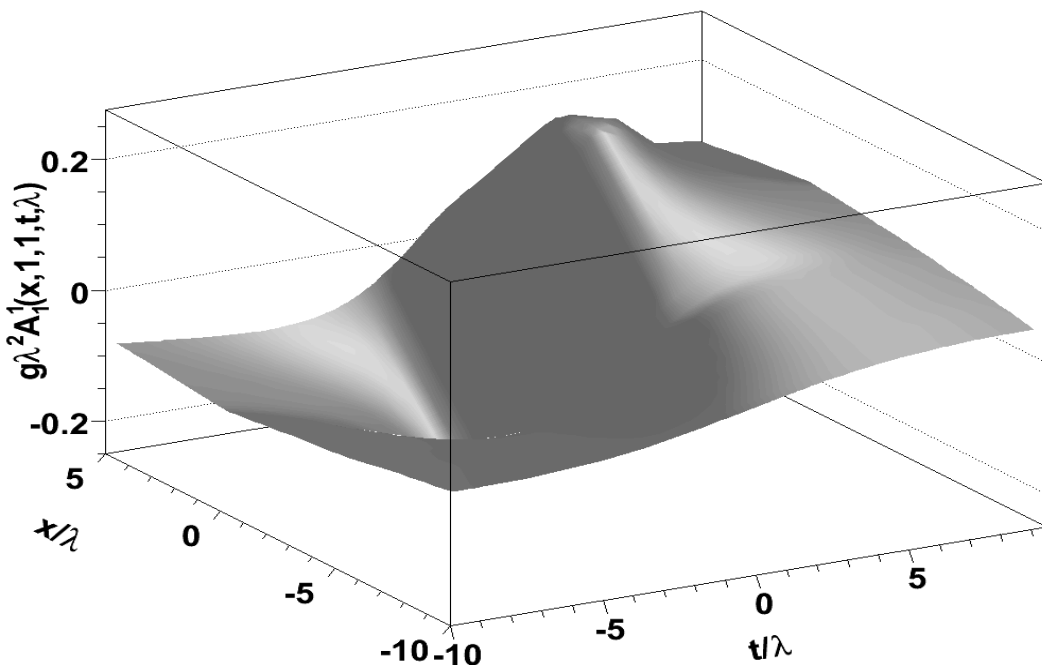


- Instanton field configuration is $A_{\mu}^a(r, \lambda) = 2r_{\nu} \eta_{\nu\mu}^a / (g(r^2 + \lambda^2))$
- It is a Landau-gauge configuration, satisfying $\partial_{\mu} A_{\mu}^a = 0$

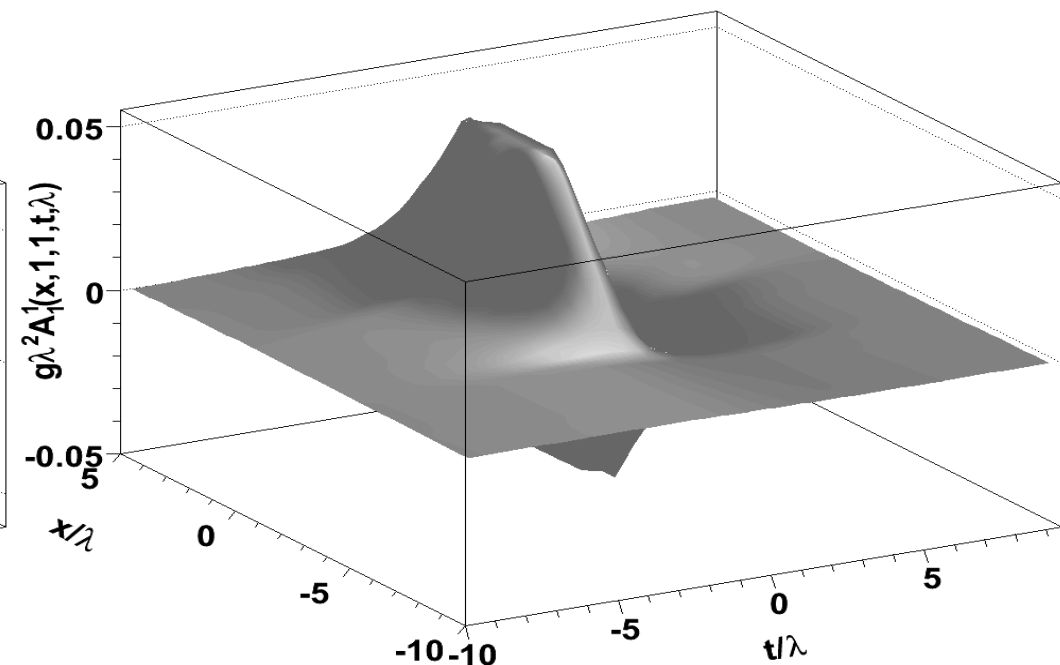
Example: Instanton

[Maas, EPJC 2006]

Instanton field



Instanton field

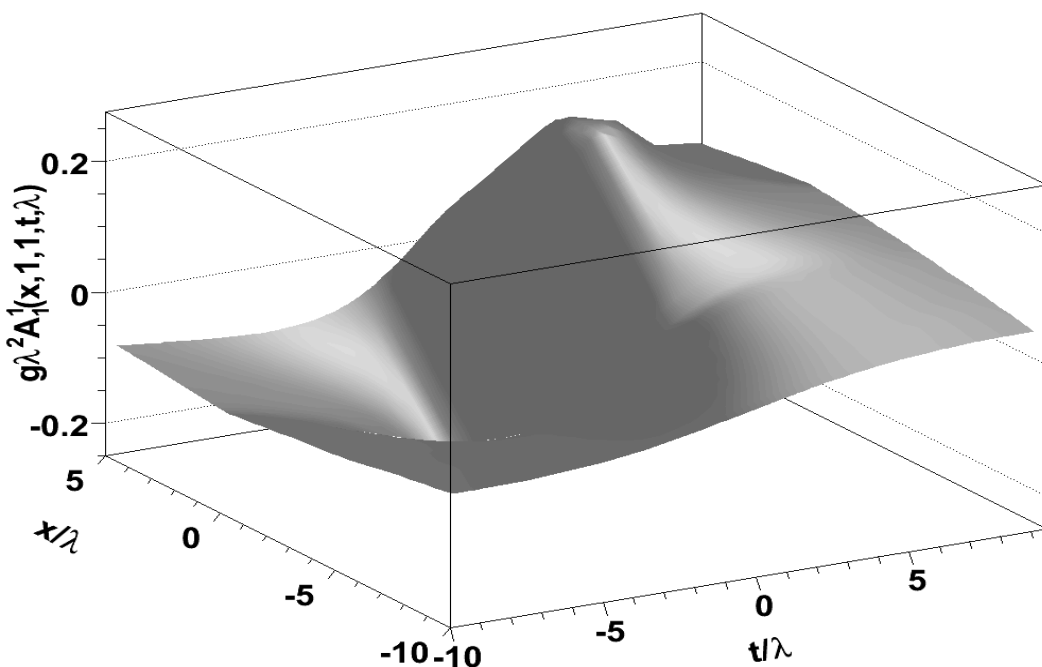


- Gauge transformation to $A_\mu^a(r, \lambda) = -2r_\nu \eta_{\nu\mu}^a \lambda^2 / (gr^2(r^2 + \lambda^2))$

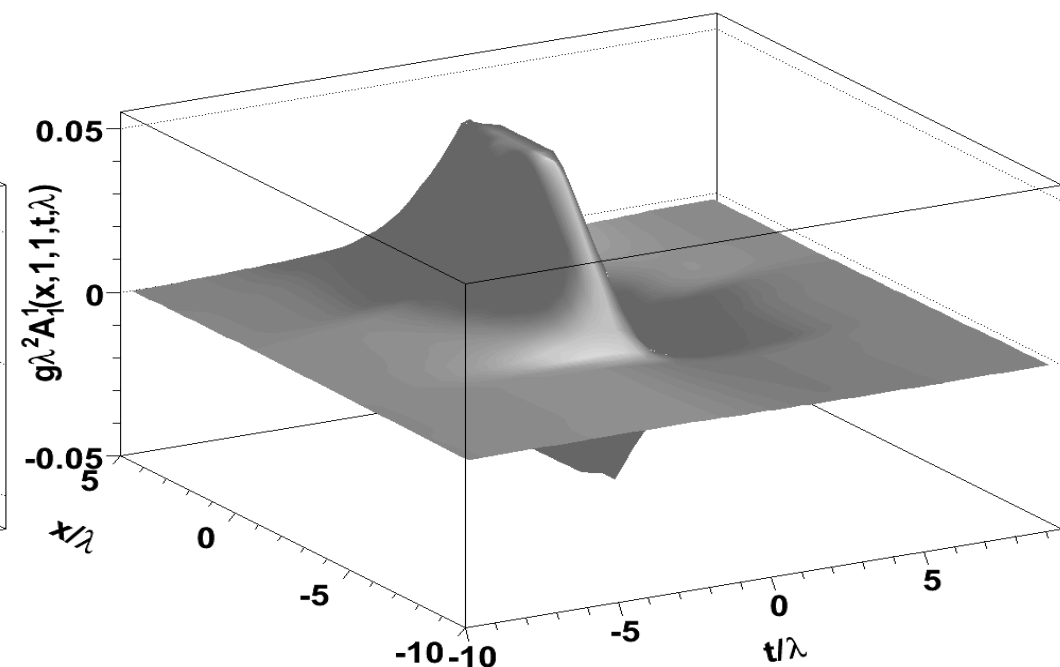
Example: Instanton

[Maas, EPJC 2006]

Instanton field



Instanton field

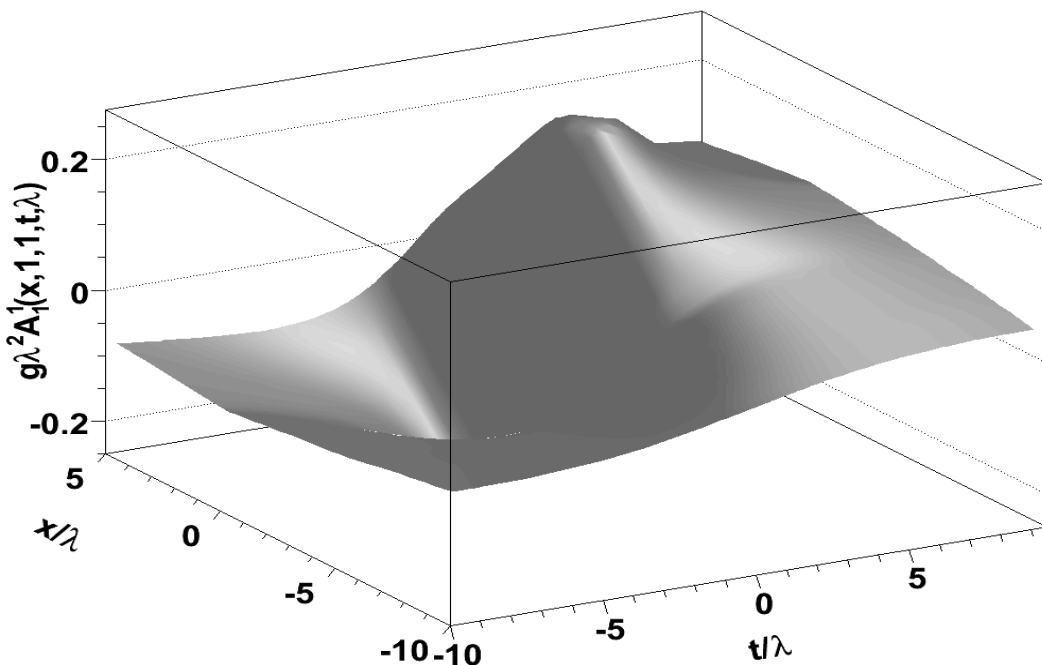


- **Gauge transformation** to $A_\mu^a(r, \lambda) = -2r_\nu \eta_{\nu\mu}^a \lambda^2 / (gr^2(r^2 + \lambda^2))$
- It is still a Landau gauge configuration!

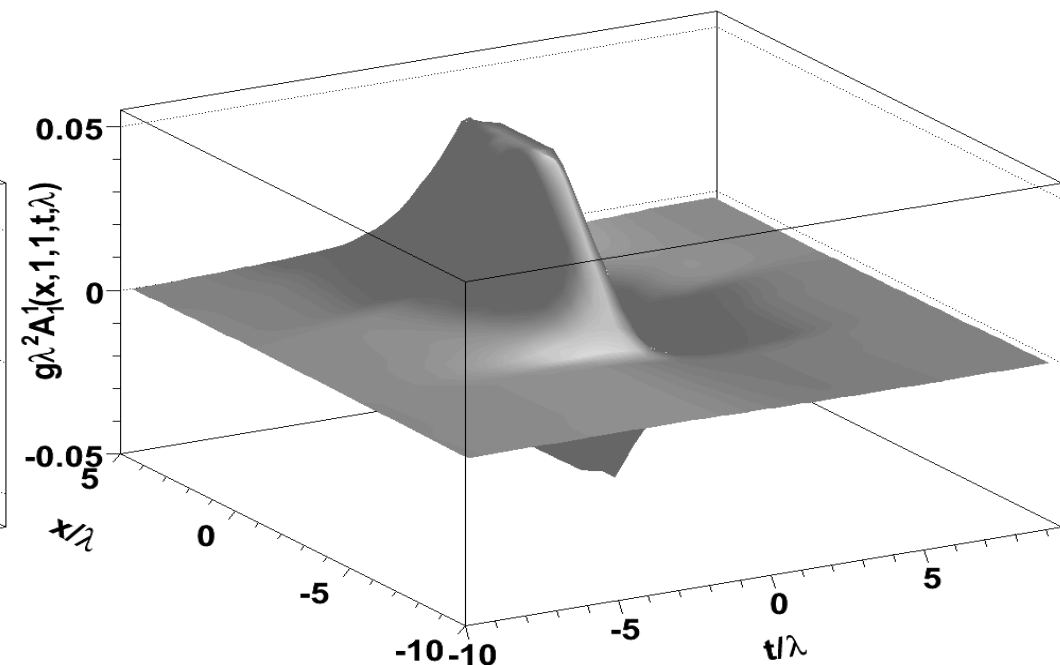
Example: Instanton

[Maas, EPJC 2006]

Instanton field



Instanton field



- **Gauge transformation** to $A_{\mu}^a(r, \lambda) = -2r_{\nu} \eta_{\nu\mu}^a \lambda^2 / (gr^2(r^2 + \lambda^2))$
 - It is still a Landau gauge configuration!
 - **Gribov copy**
 - Non-perturbative: Depends on $1/g$

Residual freedom

- Impose **Landau gauge** condition
 - Reduces configuration space to a hypersurface

Residual freedom

- Impose **Landau gauge** condition
 - Reduces configuration space to a hypersurface
- Leaves the non-perturbative gauge freedom to choose between **gauge-equivalent Gribov copies**
 - **Residual gauge orbit**
 - Set of **Gribov copies**, in general not continuous

Residual freedom

- Impose **Landau gauge** condition
 - Reduces configuration space to a hypersurface
- Leaves the non-perturbative gauge freedom to choose between **gauge-equivalent Gribov copies**
 - **Residual gauge orbit**
 - Set of **Gribov copies**, in general not continuous
- A definite prescription is required for an unambiguous result
 - There is no unique prescription how to complete Landau gauge non-perturbatively
 - There is no possibility using a local condition

Residual freedom

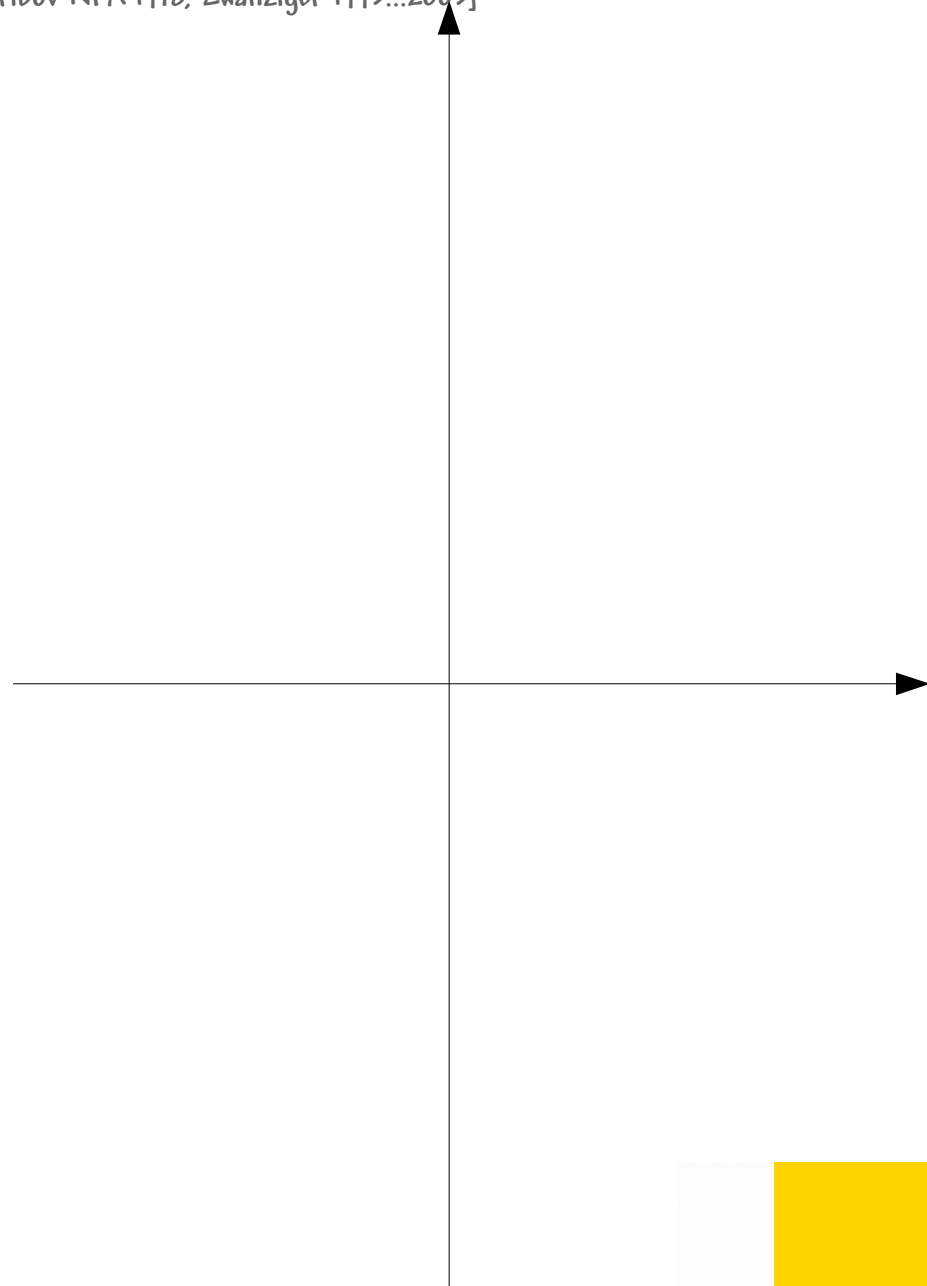
- Impose **Landau gauge** condition
 - Reduces configuration space to a hypersurface
- Leaves the non-perturbative gauge freedom to choose between **gauge-equivalent Gribov copies**
 - **Residual gauge orbit**
 - Set of **Gribov copies**, in general not continuous
- A definite prescription is required for an unambiguous result
 - There is no unique prescription how to complete Landau gauge non-perturbatively
 - There is no possibility using a local condition
- Construct a non-local condition instead to solve the problem

Residual freedom

- Impose **Landau gauge** condition
 - Reduces configuration space to a hypersurface
- Leaves the non-perturbative gauge freedom to choose between **gauge-equivalent Gribov copies**
 - **Residual gauge orbit**
 - Set of **Gribov copies**, in general not continuous
- A definite prescription is required for an unambiguous result
 - There is no unique prescription how to complete Landau gauge non-perturbatively
 - There is no possibility using a local condition
- Construct a non-local condition instead to solve the problem
- Choice: Leave the global color symmetry unfixed

Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

- Absolute Landau gauge



Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

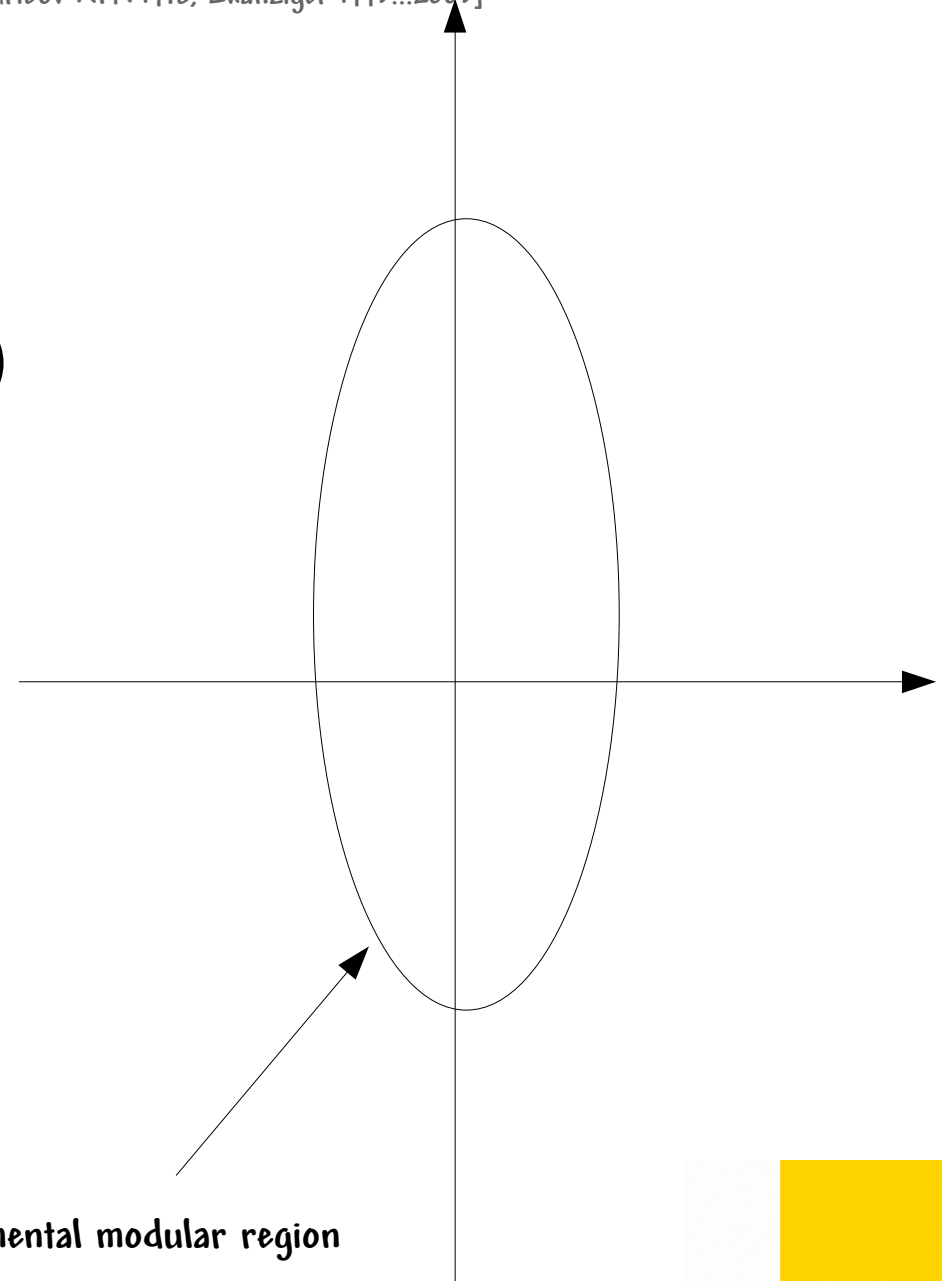
- Absolute Landau gauge

- A global minimum of

$$-\int d^d x A_\mu^a(x) A_\mu^a(x) \sim -\int d^d p D_{\mu\mu}^{aa}(p)$$

defines the **fundamental modular region**

- This region is bounded and convex



Fundamental modular region

Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

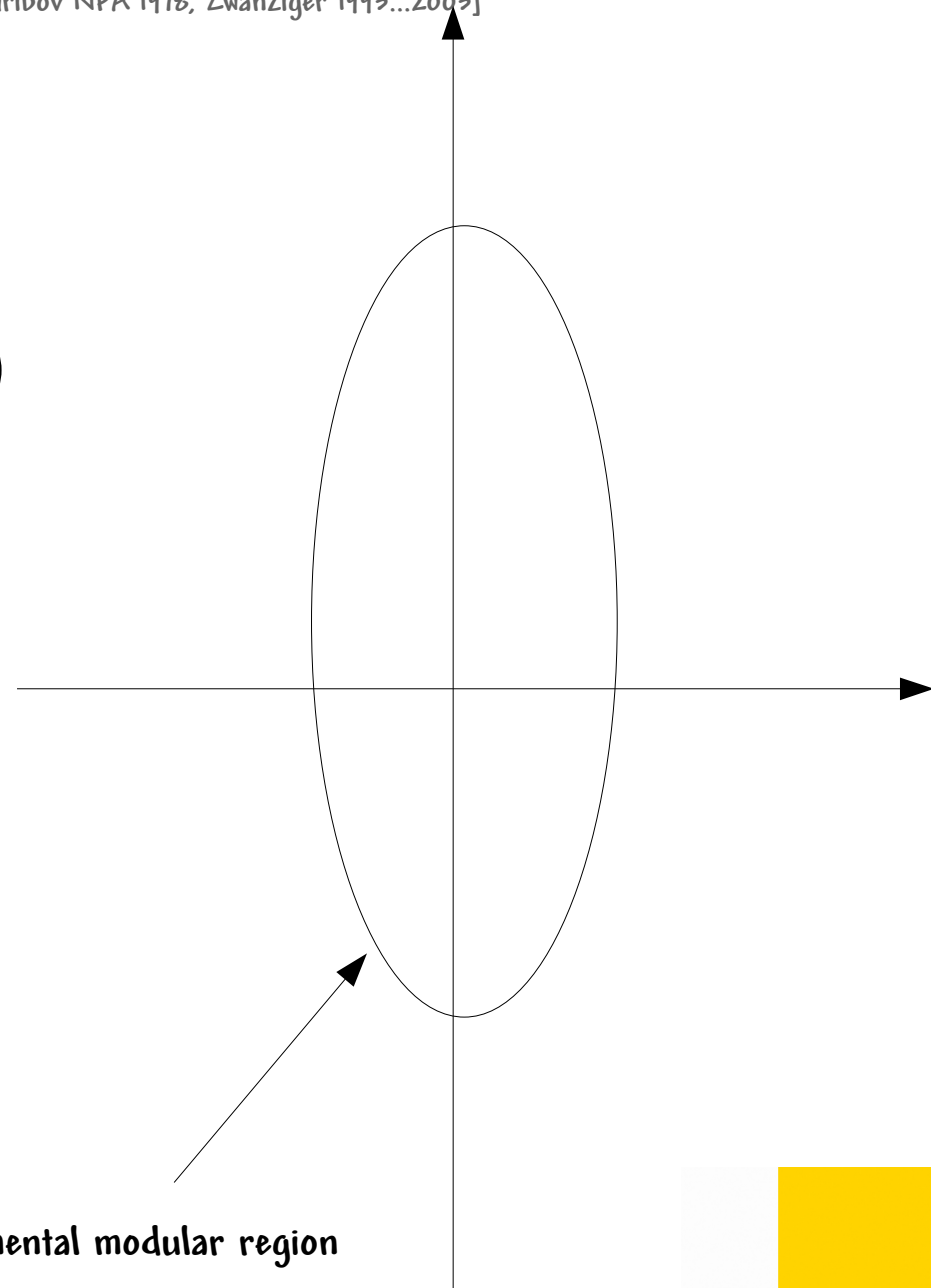
- Absolute Landau gauge

- A global minimum of

$$-\int d^d x A_\mu^a(x) A_\mu^a(x) \sim -\int d^d p D_{\mu\mu}^{aa}(p)$$

defines the **fundamental modular region**

- This region is bounded and convex
- Singles out exactly one gauge copy
 - Unique: Absolute Landau gauge



Fundamental modular region

Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

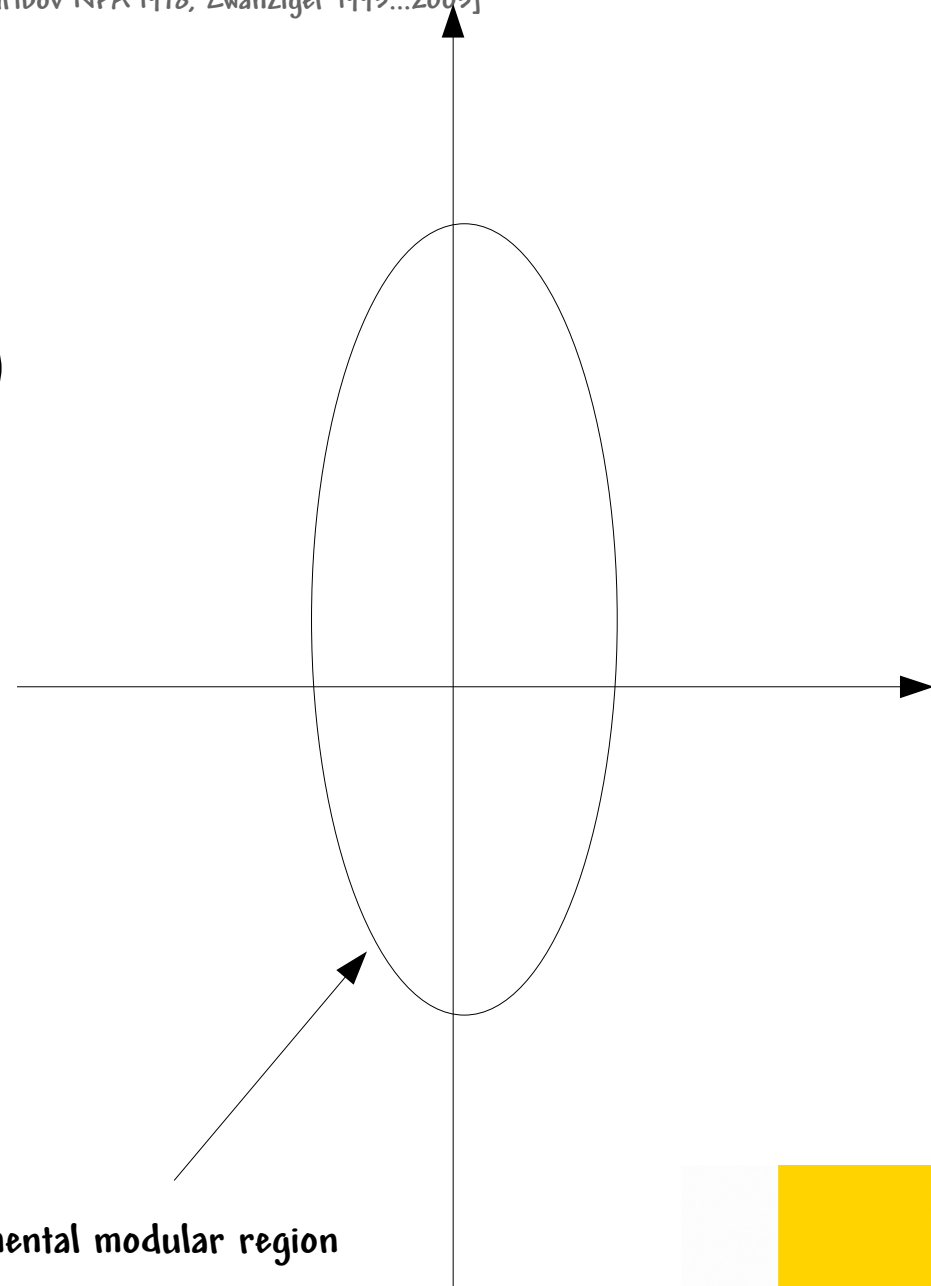
- Absolute Landau gauge

- A global minimum of

$$-\int d^d x A_\mu^a(x) A_\mu^a(x) \sim -\int d^d p D_{\mu\mu}^{aa}(p)$$

defines the **fundamental modular region**

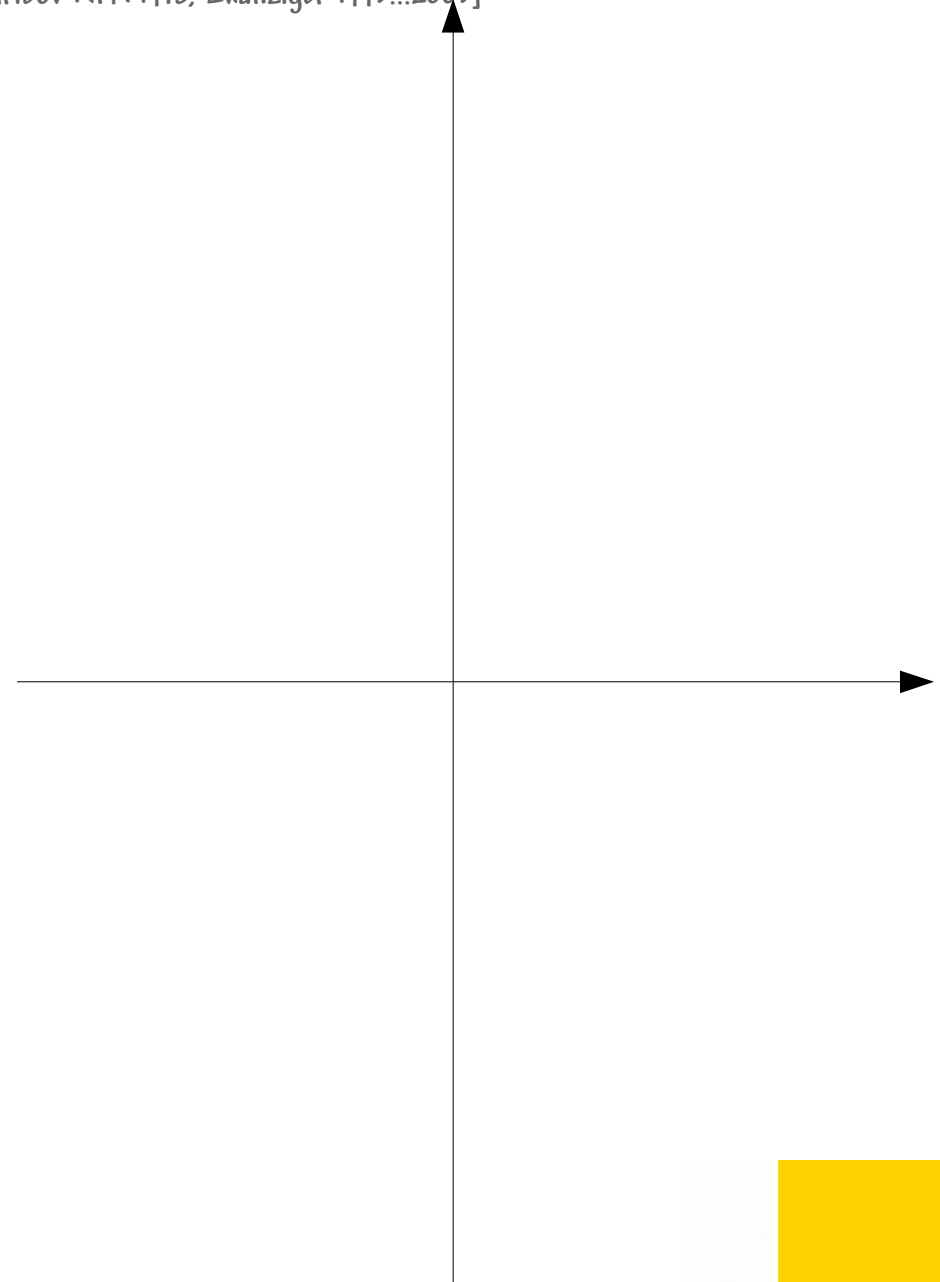
- This region is bounded and convex
- Singles out exactly one gauge copy
 - Unique: Absolute Landau gauge
- Drawback: Badly divergent and regularization may introduce degeneracies



Fundamental modular region

Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

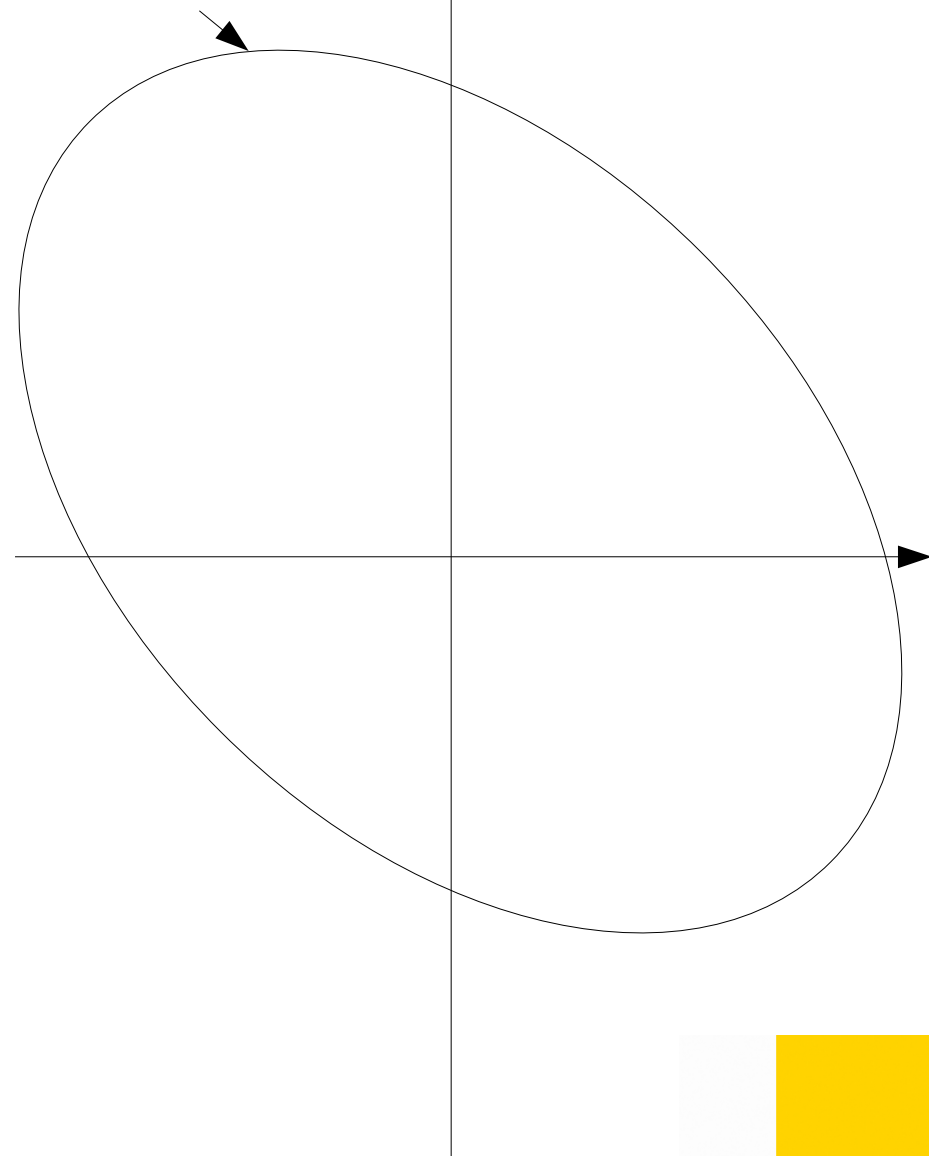
- Different approach: Enlarge the search space



Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

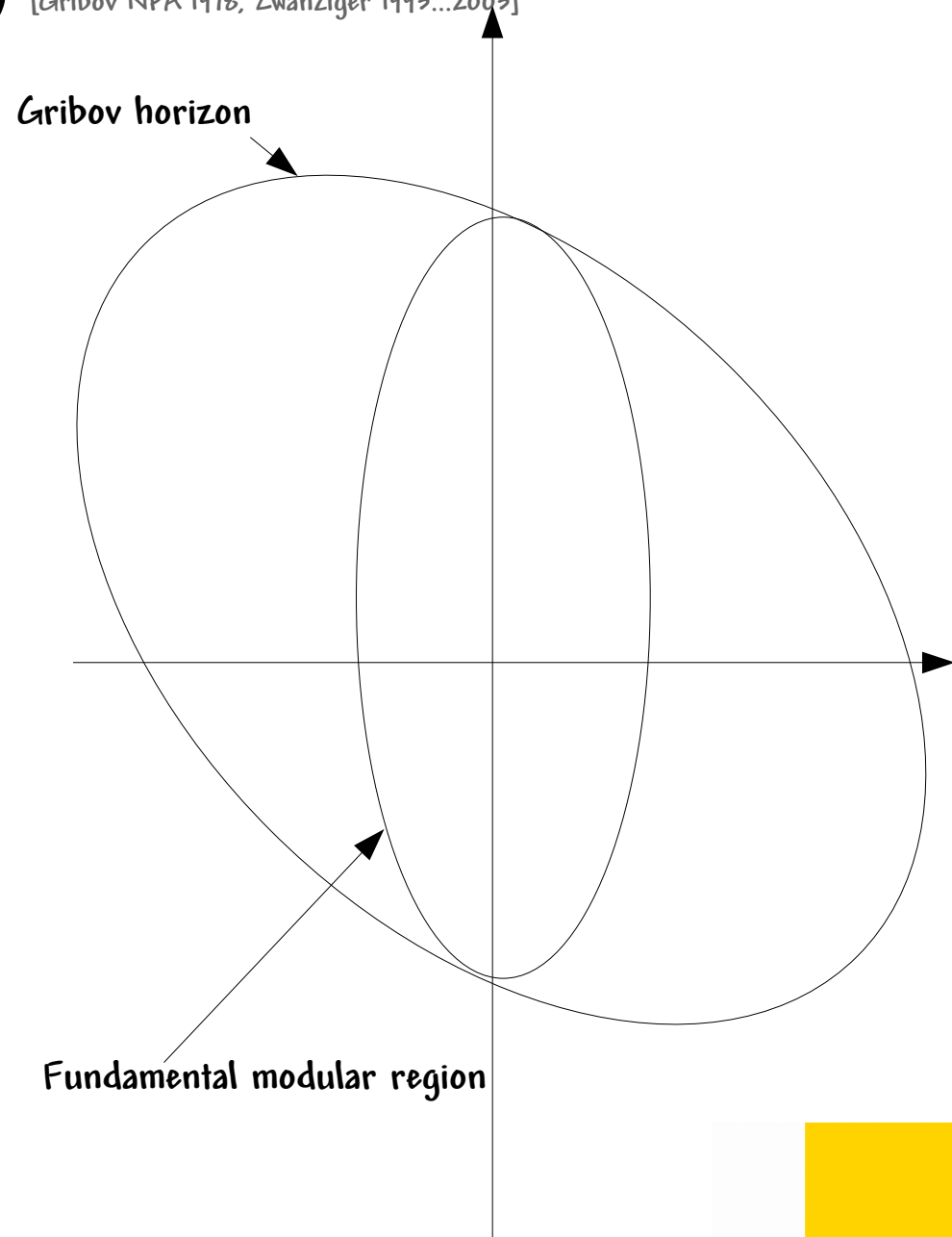
- Different approach: Enlarge the search space
- **Gribov horizon** encloses all field configurations with positive Faddeev-Popov operator $(-\partial_\mu D_\mu)$

Gribov horizon



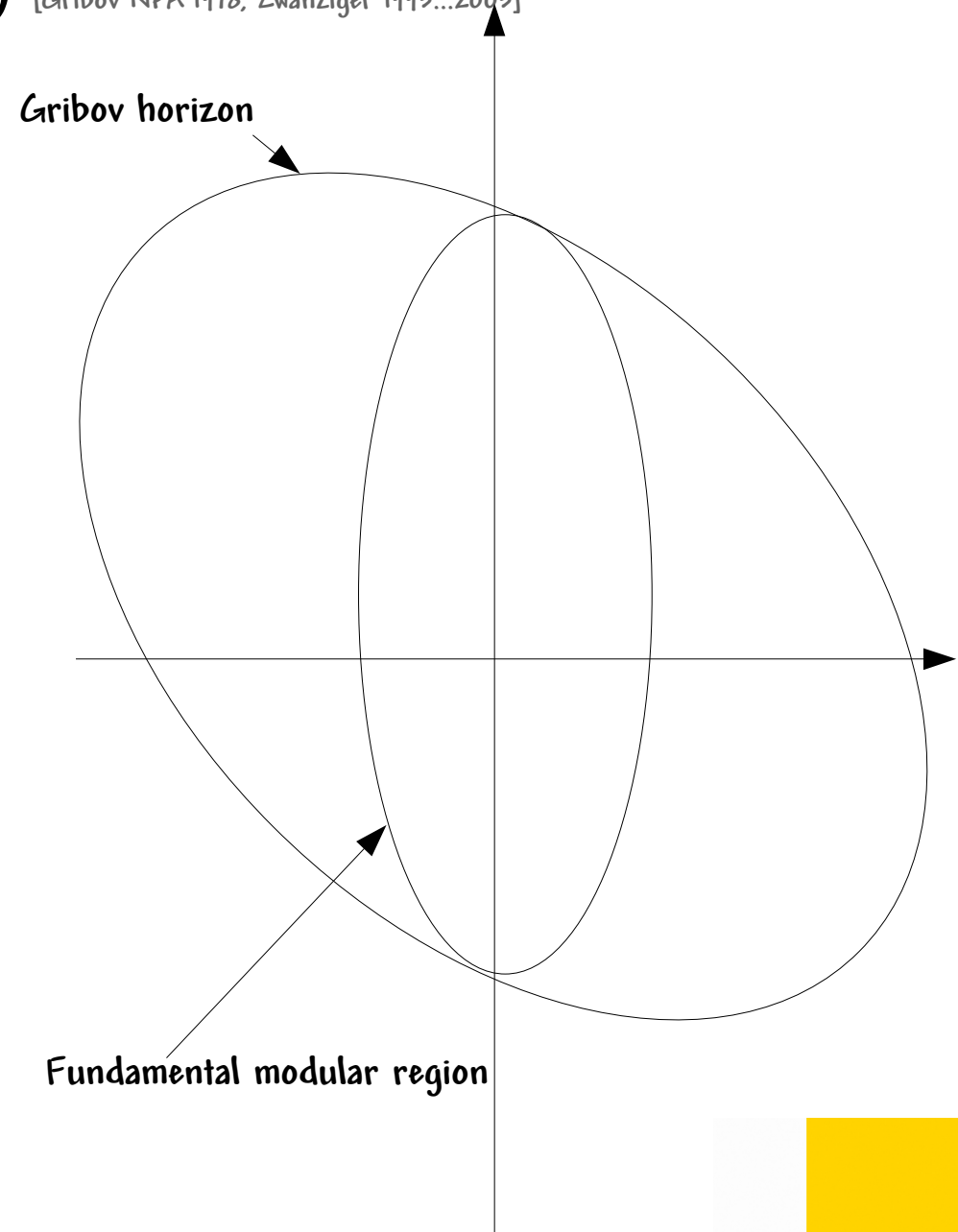
Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

- Different approach: Enlarge the search space
- **Gribov horizon** encloses all field configurations with positive Faddeev-Popov operator $(-\partial_\mu D_\mu)$
- Includes the **fundamental modular region**



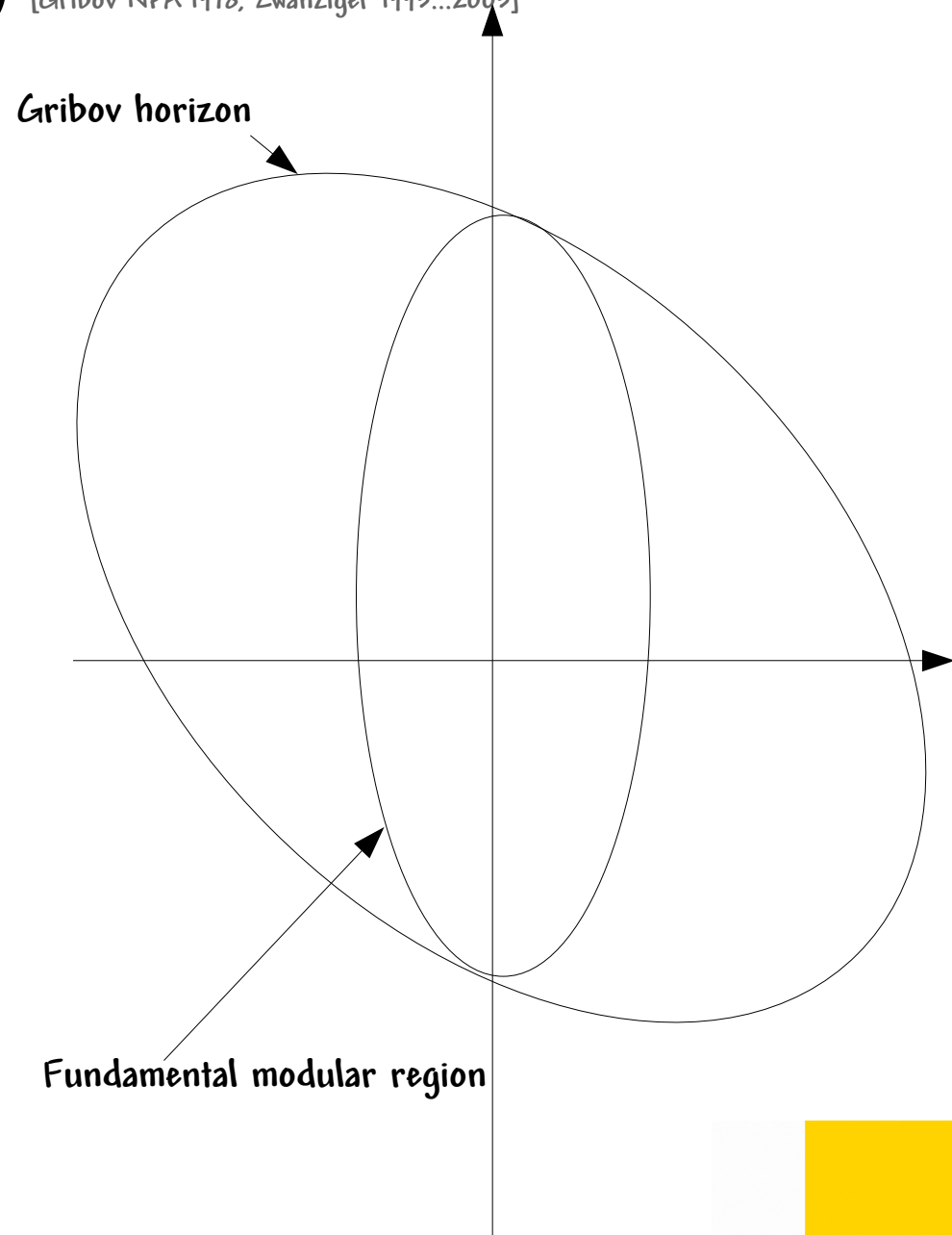
Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

- Different approach: Enlarge the search space
- **Gribov horizon** encloses all field configurations with positive Faddeev-Popov operator $(-\partial_\mu D_\mu)$
- Includes the **fundamental modular region**
- All gauge orbits pass through this region



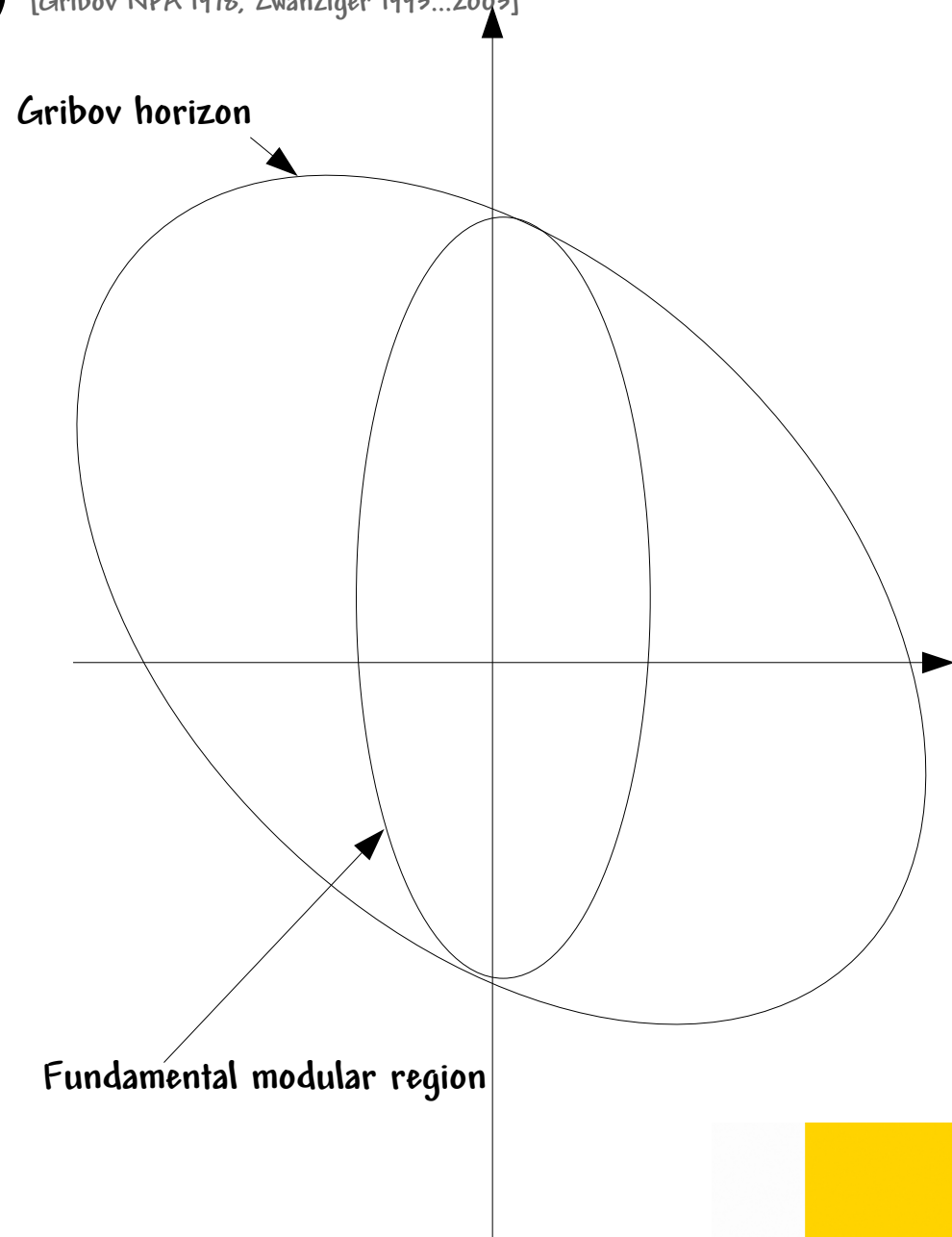
Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

- Different approach: Enlarge the search space
- **Gribov horizon** encloses all field configurations with positive Faddeev-Popov operator $(-\partial_\mu D_\mu)$
- Includes the **fundamental modular region**
- All gauge orbits pass through this region
 - Many **Gribov copies** for each
 - How many is many?



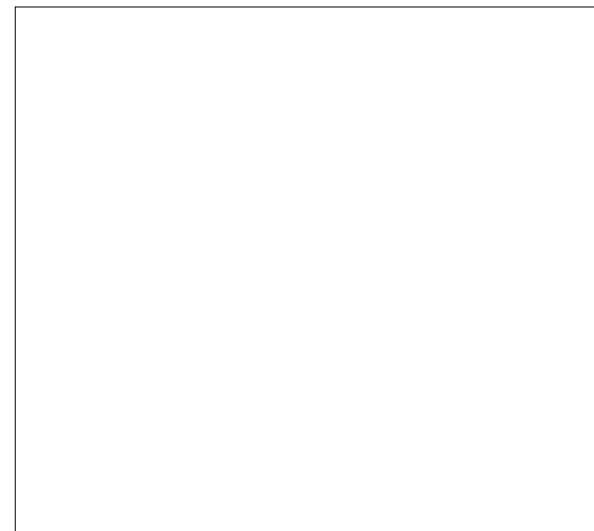
Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

- Different approach: Enlarge the search space
- **Gribov horizon** encloses all field configurations with positive Faddeev-Popov operator $(-\partial_\mu D_\mu)$
- Includes the **fundamental modular region**
- All gauge orbits pass through this region
 - Many **Gribov copies** for each
 - How many is many?
 - Requires a **tool** for investigations



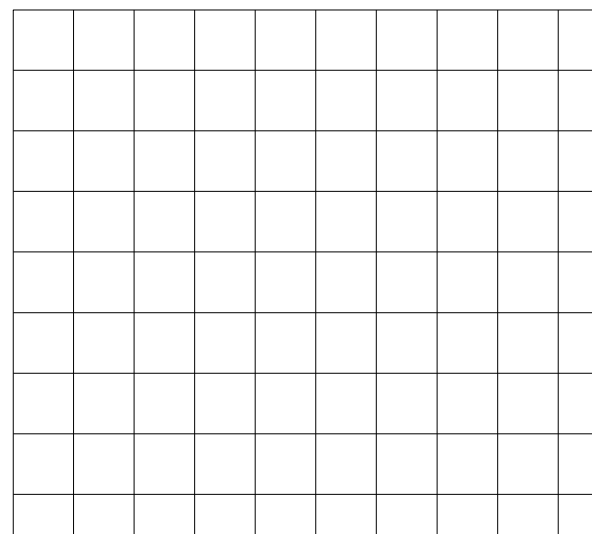
Lattice calculations

- Take a **finite volume** – usually a hypercube



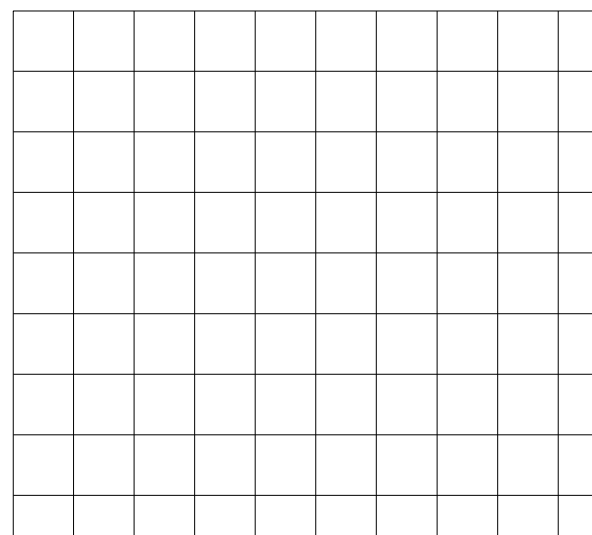
Lattice calculations

- Take a **finite volume** – usually a hypercube
- Discretize it, and get a **finite, hypercubic lattice**



Lattice calculations

- Take a **finite volume** – usually a hypercube
- Discretize it, and get a **finite, hypercubic lattice**
- **Calculate observables using path integration**
 - $\langle \bar{c} c \rangle = \int dA d\bar{c} dc \bar{c} c \exp(-\int d^d x L)$
 - Can be done numerically
 - Uses **Monte-Carlo methods**



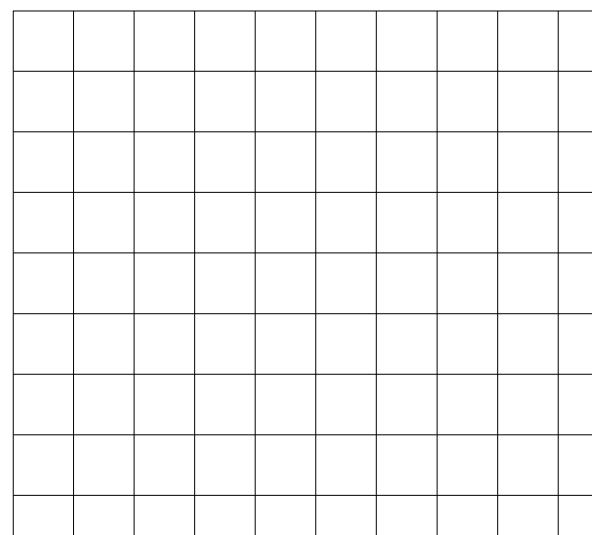
Lattice calculations

- Take a **finite volume** – usually a hypercube
- Discretize it, and get a **finite, hypercubic lattice**
- **Calculate observables using path integration**

- $\langle \bar{c} c \rangle = \int dA d\bar{c} dc \bar{c} c \exp(-\int d^d x L)$

- Can be done numerically
- Uses **Monte-Carlo methods**

- **Artifacts**
 - Finite volume/discretization
 - Zero momentum problematic



Identifying Gribov copies

- Identifying *Gribov copies* is a non-trivial task

Identifying Gribov copies

- Identifying **Gribov copies** is a non-trivial task
 - Problem: No constructive possibility to generate all Gribov copies on the residual gauge orbit

Identifying Gribov copies

- Identifying **Gribov copies** is a non-trivial task
 - Problem: No constructive possibility to generate all Gribov copies on the residual gauge orbit
 - Practical solution: Generate for each residual gauge orbit many random copies and select different ones [Cucchieri, 1997]

Identifying Gribov copies

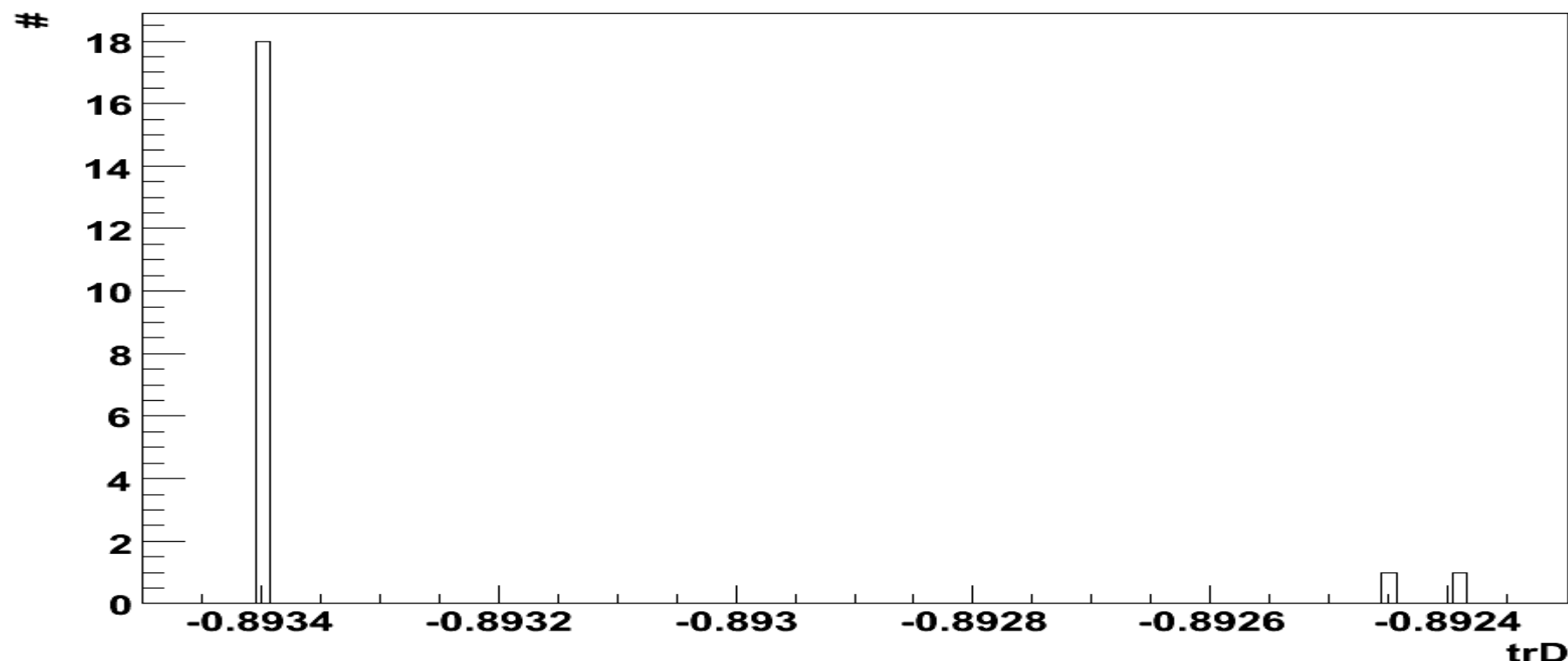
- Identifying **Gribov copies** is a non-trivial task
 - Problem: No constructive possibility to generate all Gribov copies on the residual gauge orbit
 - Practical solution: Generate for each residual gauge orbit many random copies and select different ones [Cucchieri, 1997]
 - How to decide whether two copies are different?

Identifying Gribov copies

- Identifying **Gribov copies** is a non-trivial task
 - Problem: No constructive possibility to generate all Gribov copies on the residual gauge orbit
 - Practical solution: Generate for each residual gauge orbit many random copies and select different ones [Cucchieri, 1997]
 - How to decide whether two copies are different?
 - Non-trivial problem. Requires in principle to show a connection by a non-trivial gauge transformation
 - Practical solution: Require sufficiently different value of $\text{Tr}D$
 - Other are dismissed as **numerical/lattice artifacts**
 - Is this good? Will possibly work only on a **lattice!**

TrD on the residual gauge orbit [3d, beta=3.46, Maas, unpublished]

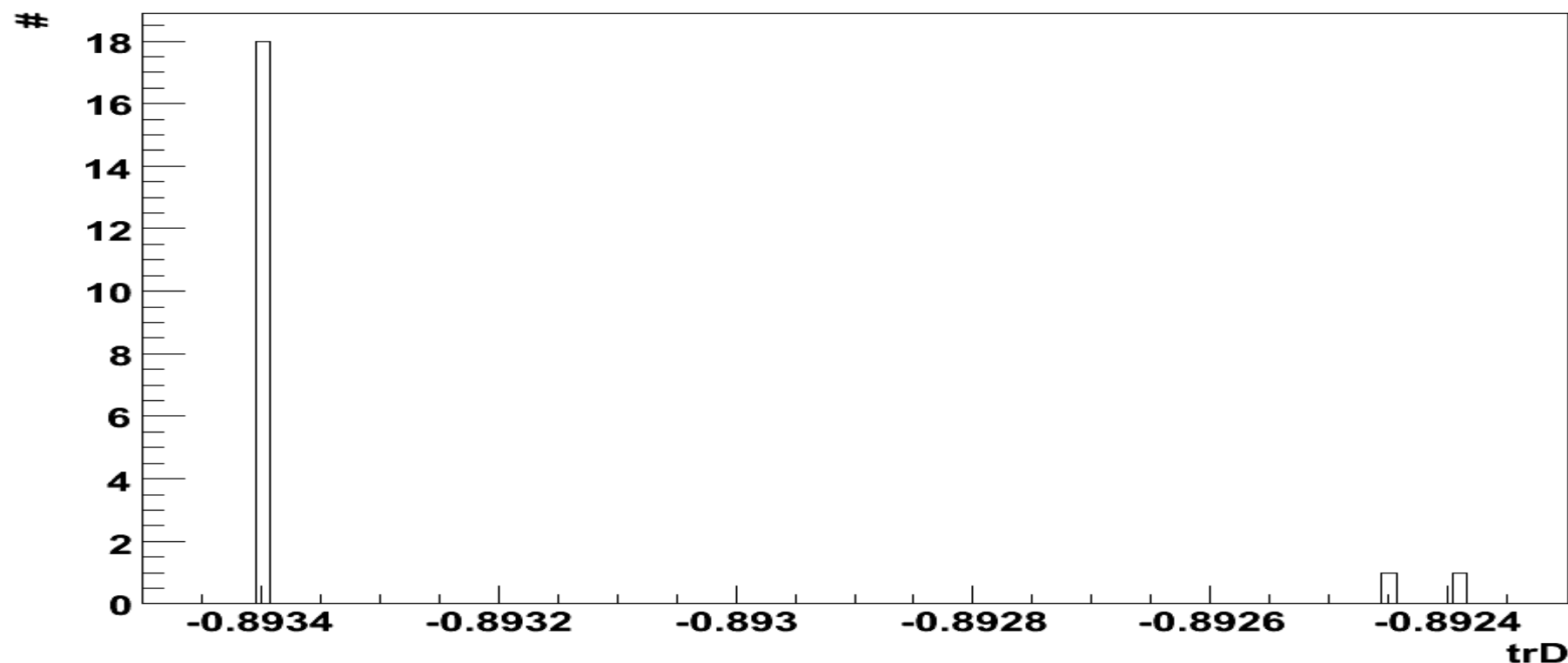
TrD for $V=(3.1 \text{ fm})^3$



- Rather different values of TrD for different **Gribov copies**
- Many artificial copies of one particular copy: Attractive basin in TrD

TrD on the residual gauge orbit [3d, beta=3.46, Maas, unpublished]

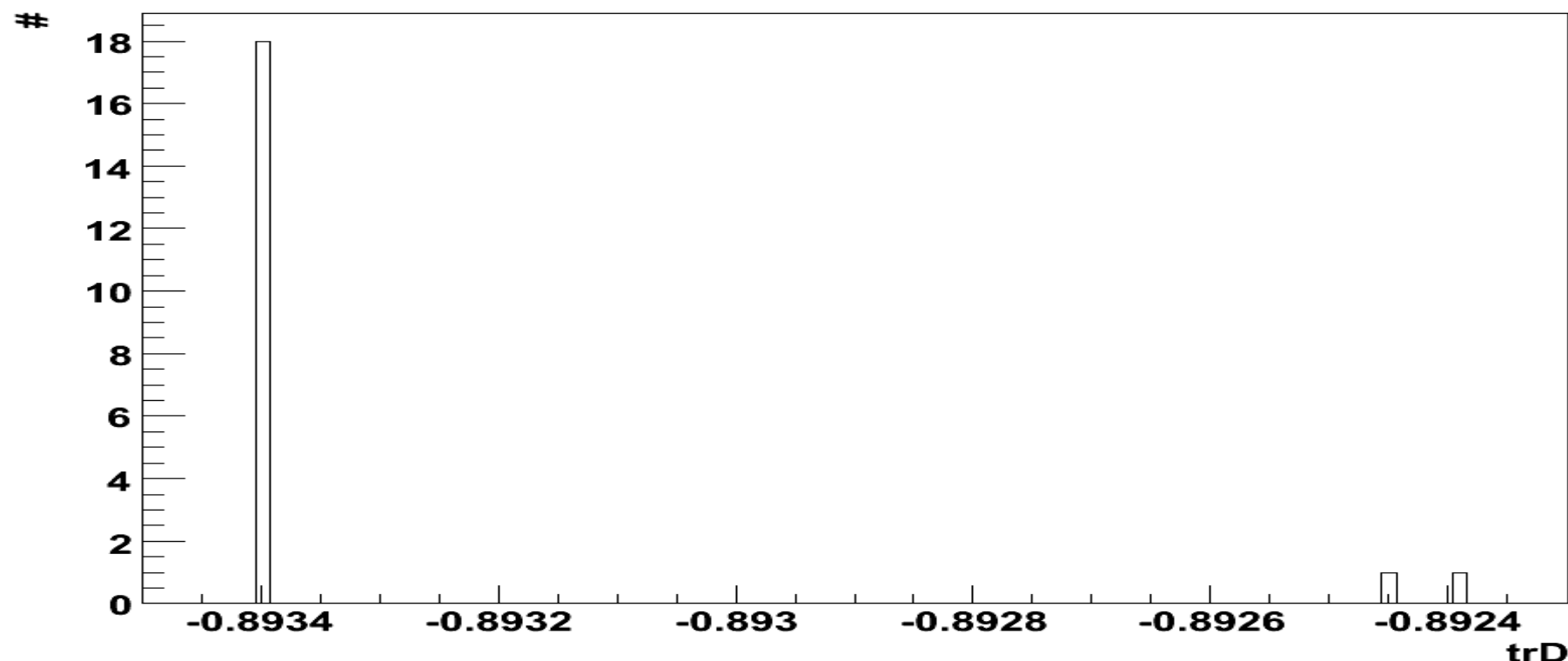
TrD for $V=(3.1 \text{ fm})^3$



- Rather different values of TrD for different **Gribov copies**
- Many artificial copies of one particular copy: Attractive basin in TrD
- Difference in peak is below 10^{-10} : Numerical/lattice artifacts

TrD on the residual gauge orbit [3d, beta=3.46, Maas, unpublished]

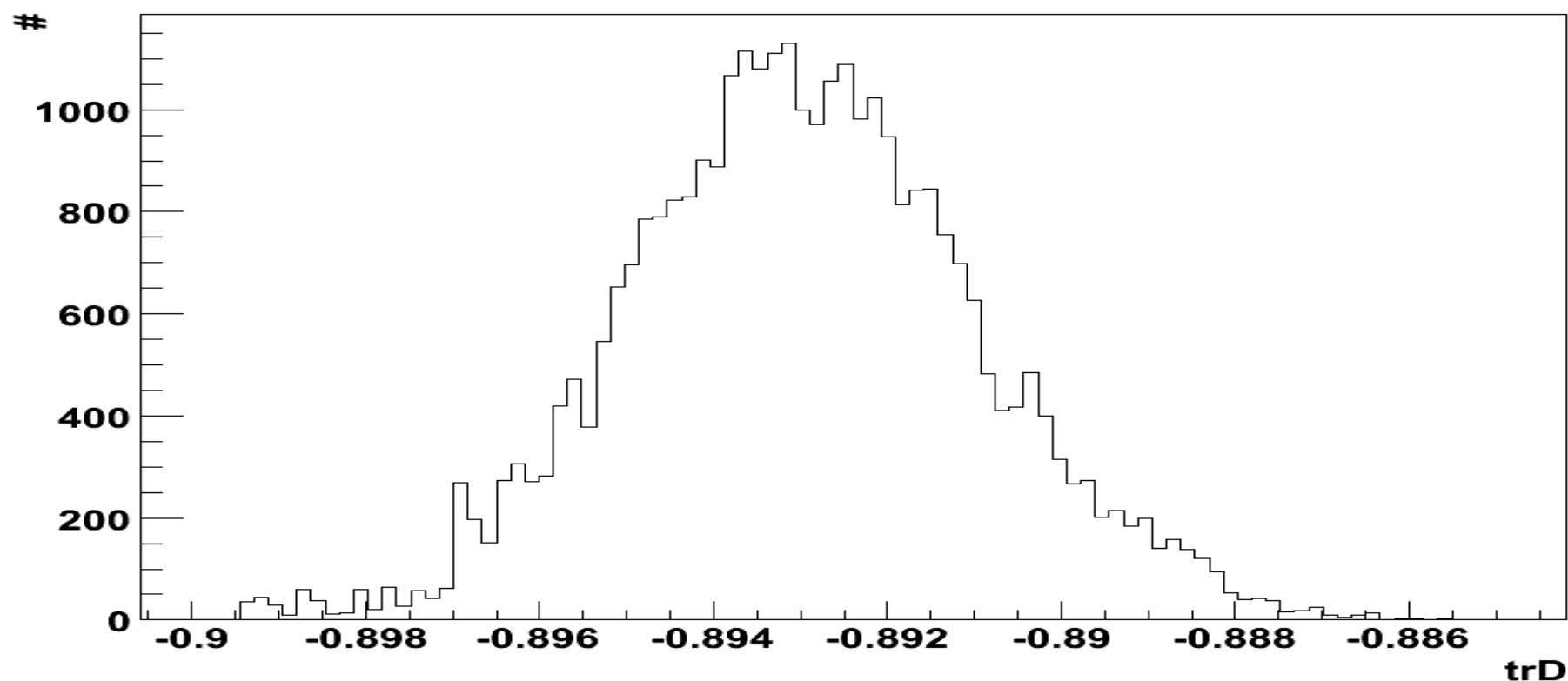
TrD for $V=(3.1 \text{ fm})^3$



- Rather different values of TrD for different **Gribov copies**
- Many artificial copies of one particular copy: Attractive basin in TrD
- Difference in peak is below 10^{-10} : Numerical/lattice artifacts
- Pragmatic definition: Copy requires a separation of 10^{-5} in TrD

Distribution of TrD [3d, beta=3.46, Maas, unpublished]

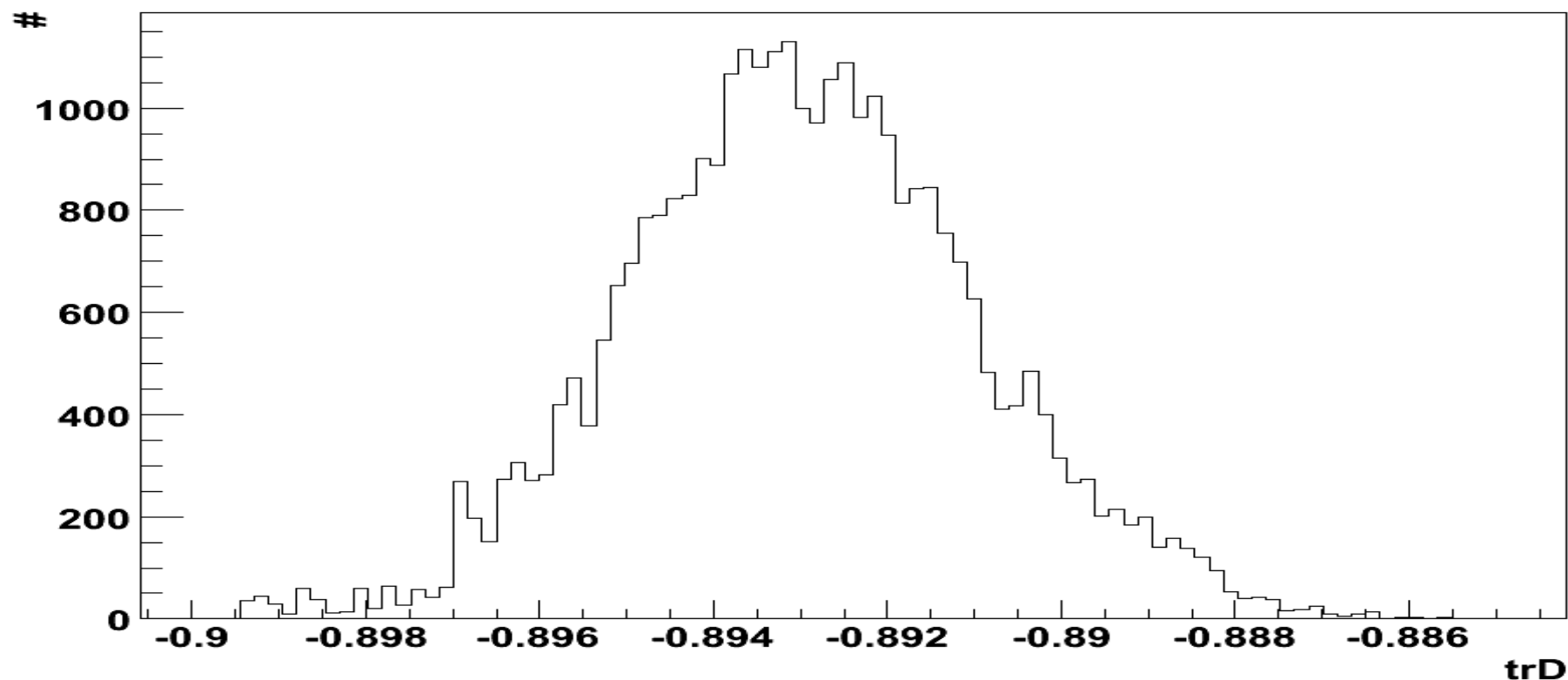
TrD for $V=(3.1 \text{ fm})^3$ from 1622 configurations



- Distribution is almost Gaussian
 - Never a real Gaussian

Distribution of TrD [3d, beta=3.46, Maas, unpublished]

TrD for $V=(3.1 \text{ fm})^3$ from 1622 configurations



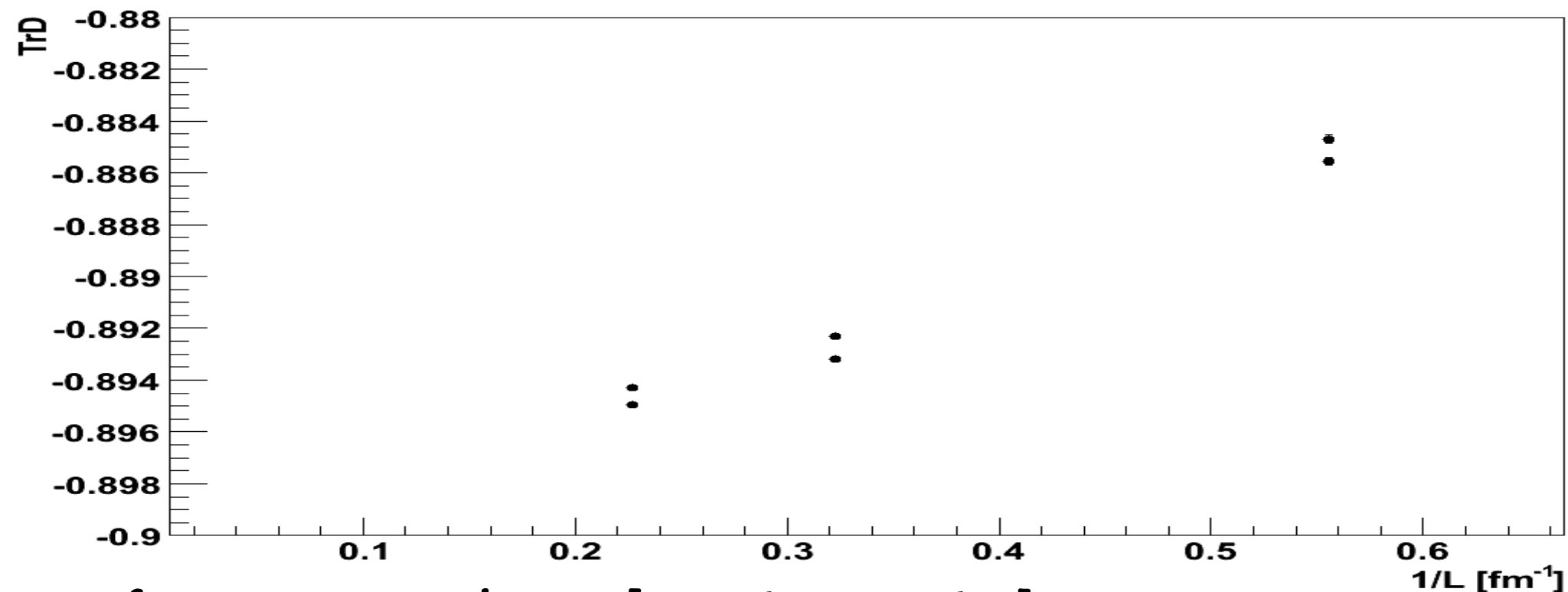
- Distribution is almost Gaussian
 - Never a real Gaussian
- Variation of TrD corresponds to varying gluon propagator

TrD as a gauge parameter [3d, beta=3.45, Maas, unpublished]

- Define gauge corridor as $[\max \text{TrD}, \min \text{TrD}]$

TrD as a gauge parameter [3d, beta=3.45, Maas, unpublished]

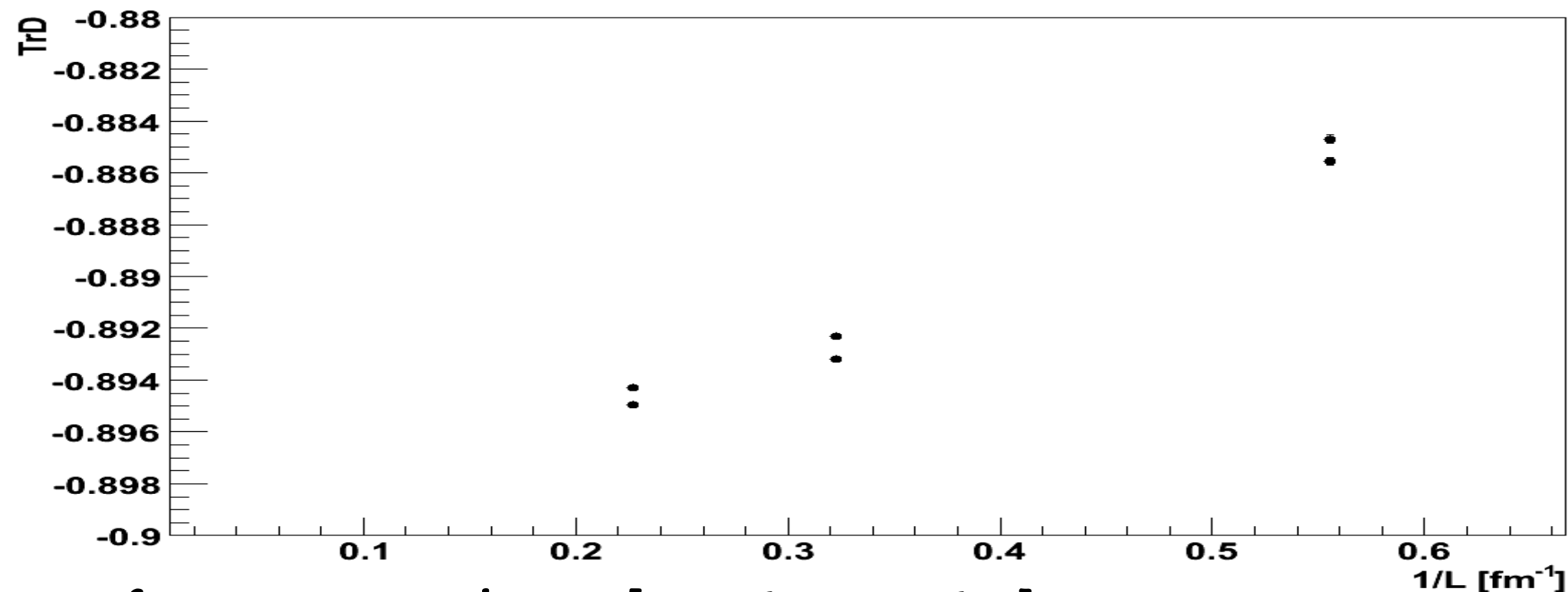
TrD corridor



- Define gauge corridor as $[\max \text{TrD}, \min \text{TrD}]$ [Zwanziger, 2003]

TrD as a gauge parameter [3d, beta=3.45, Maas, unpublished]

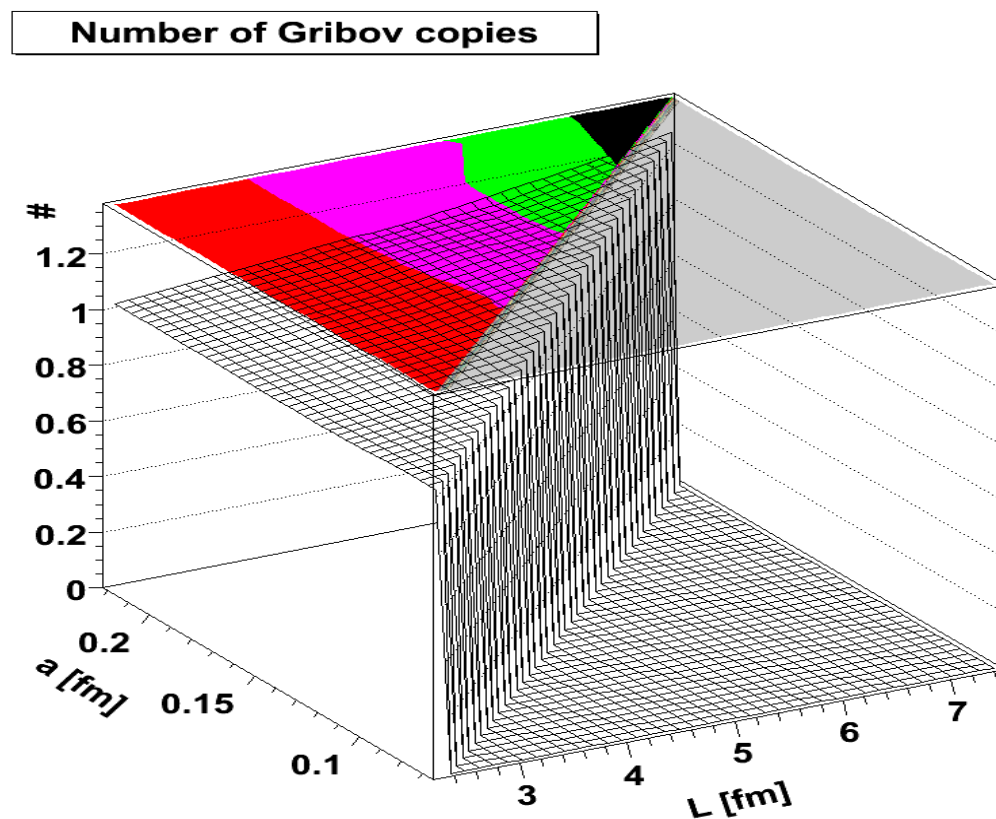
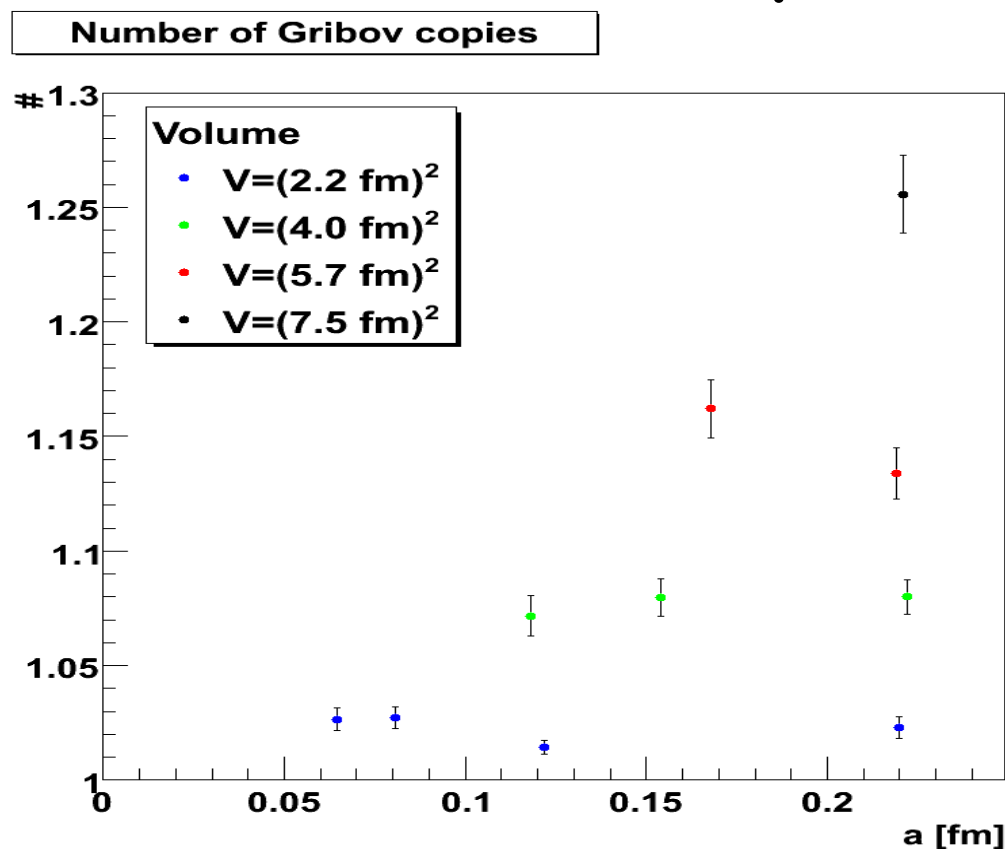
TrD corridor



- Define gauge corridor as [max TrD, min TrD]
- TrD is not showing a strong volume dependence
 - For some theories known to degenerate [de Forcrand, 1994]
 - Degeneracies would support the Zwanziger entropy conjecture

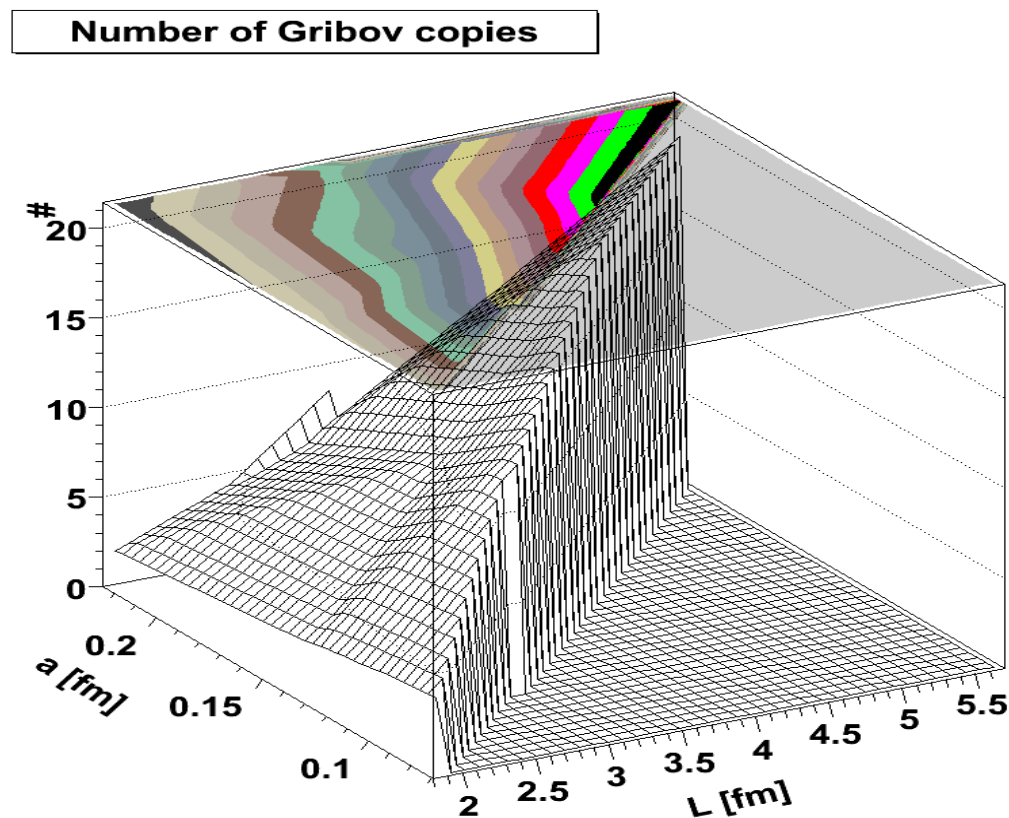
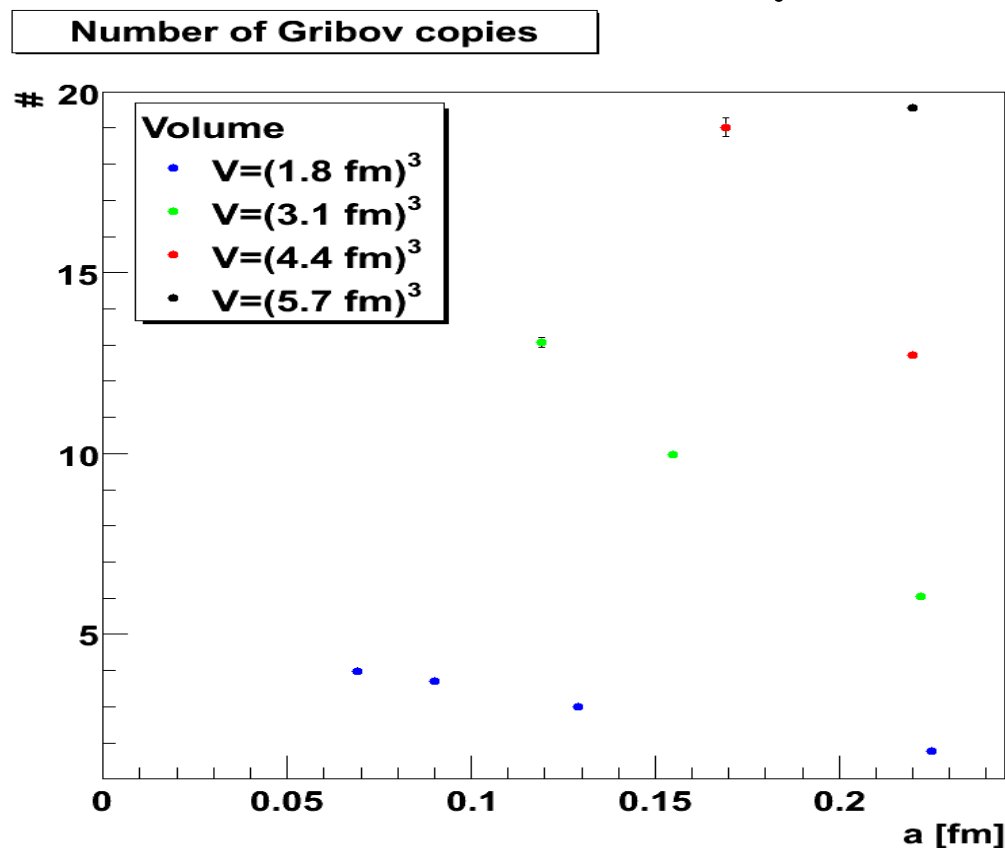
[Zwanziger, 2003]

Number of Gribov copies in two dimensions [Maas, unpublished]



- Very slow evolution with volume and discretization
- Very large volumes will be needed to see an effect

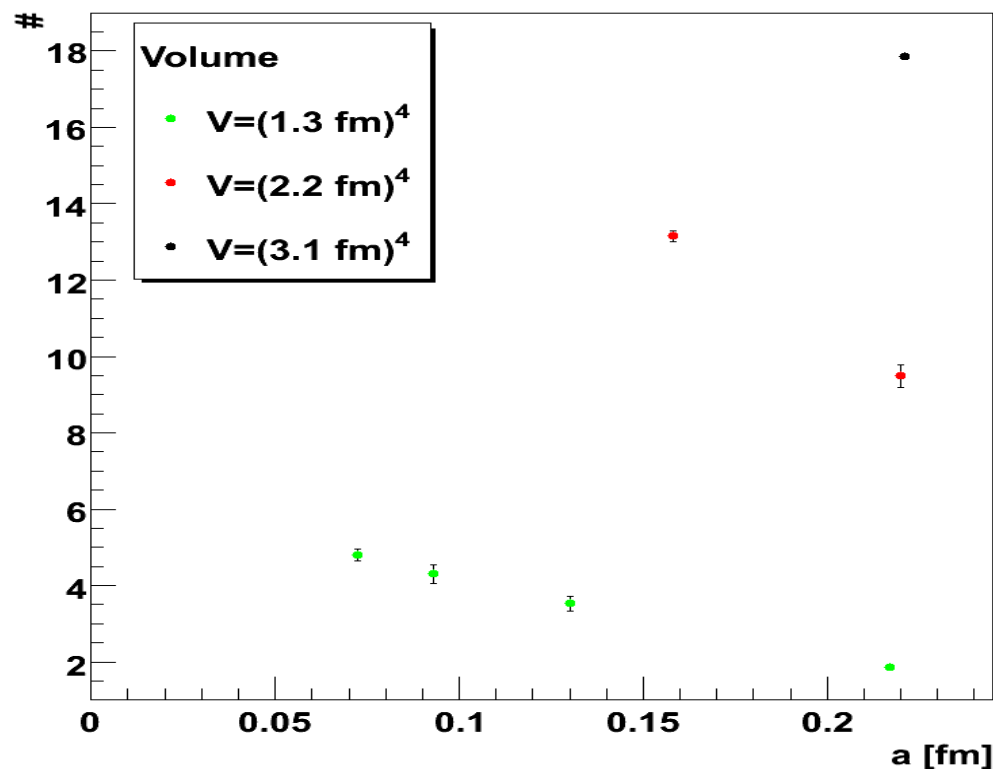
Number of Gribov copies in three dimensions [Maas, unpublished]



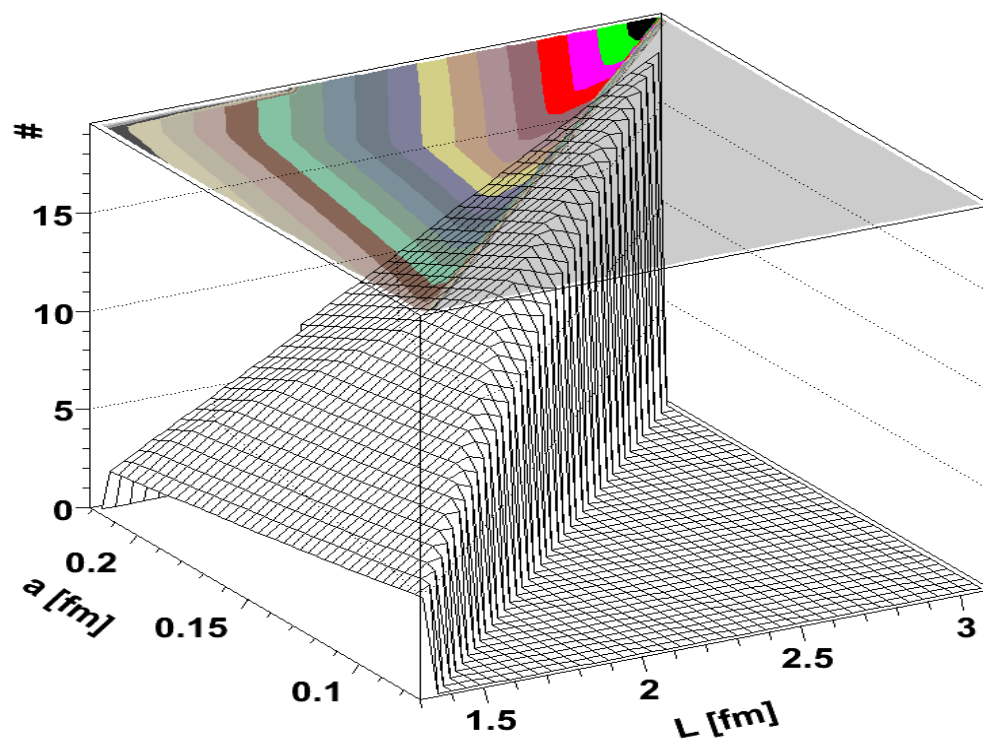
- Evolution with volume and discretization
- Already rather small volumes have a large number of copies

Number of Gribov copies in four dimensions [Maas, unpublished]

Number of Gribov copies



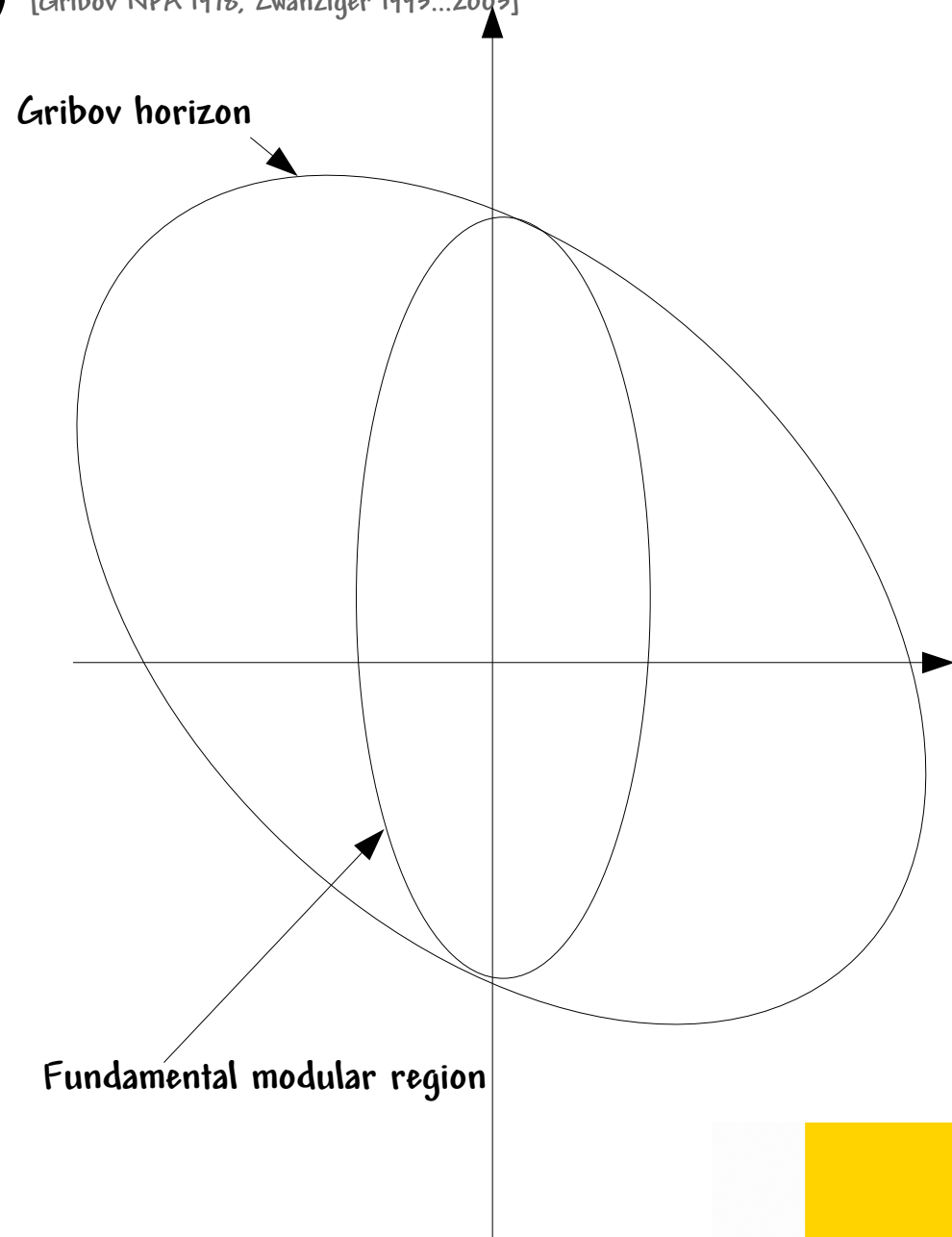
Number of Gribov copies



- Large number of copies even on small volumes
- Big effects can be expected
- Number of Gribov copies rises quickly with the number of dimensions

Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

- Different approach: Enlarge the search space
- **Gribov horizon** encloses all field configurations with positive Faddeev-Popov operator $(-\partial_\mu D_\mu)$
- Includes the **fundamental modular region**
- All gauge orbits pass through this region
 - Many **Gribov copies** for each
- How to select a unique representative?



Constructing a global gauge condition

- At the boundary of the first Gribov horizon the Faddeev-Popov operator develops zero eigenmodes

Constructing a global gauge condition

- At the boundary of the **first Gribov horizon** the Faddeev-Popov operator develops zero eigenmodes
- Lowest eigenvalue of the Faddeev-Popov operator varies along the gauge orbits
- Finite quantity, but very hard to determine in general

Constructing a global gauge condition

- At the boundary of the **first Gribov horizon** the Faddeev-Popov operator develops zero eigenmodes
- Lowest eigenvalue of the Faddeev-Popov operator varies along the gauge orbits
- Finite quantity, but very hard to determine in general
- Is there also an expression in terms of a **correlation function**?

Constructing a global gauge condition

- At the boundary of the **first Gribov horizon** the Faddeev-Popov operator develops zero eigenmodes
- Lowest eigenvalue of the Faddeev-Popov operator varies along the gauge orbits
- Finite quantity, but very hard to determine in general
- Is there also an expression in terms of a **correlation function**?
- Not directly. But the **ghost propagator** is influenced by the eigenspectrum:

$$\frac{-G(p)}{p^2} = D_G^{ab}(p) \sim \langle (\partial_\mu D_\mu^{ab})^{-1} \rangle \sim \sum_i \frac{1}{\lambda_i} |\psi_i|^2$$

Constructing a global gauge condition

- At the boundary of the **first Gribov horizon** the Faddeev-Popov operator develops zero eigenmodes
- Lowest eigenvalue of the Faddeev-Popov operator varies along the gauge orbits
- Finite quantity, but very hard to determine in general
- Is there also an expression in terms of a **correlation function**?
- Not directly. But the **ghost propagator** is influenced by the eigenspectrum:

$$\frac{-G(p)}{p^2} = D_G^{ab}(p) \sim \langle (\partial_\mu D_\mu^{ab})^{-1} \rangle \sim \sum_i \frac{1}{\lambda_i} |\psi_i|^2$$

- **$G(p)$** candidate for a characterization of a copy

Constructing a global gauge condition [Maas, unpublished, Maas, 2008, Fischer et al. 2008]

- Effects will be strongest at the most non-perturbative momenta: The infrared

Constructing a global gauge condition [Maas, unpublished, Maas, 2008, Fischer et al. 2008]

- Effects will be strongest at the most non-perturbative momenta: The infrared
- Is the renormalization-group invariant $b = G(0)/G(P)$ with $P > 0$ fixed a possible second gauge parameter?
 - Fourier-transform at low momentum: Highly non-local
 - Does it vary along the (residual) gauge orbit inside the first Gribov region?

Constructing a global gauge condition [Maas, unpublished, Maas, 2008, Fischer et al. 2008]

- Effects will be strongest at the most non-perturbative momenta: The infrared
- Is the renormalization-group invariant $b = G(0)/G(P)$ with $P > 0$ fixed a possible second gauge parameter?
 - Fourier-transform at low momentum: Highly non-local
 - Does it vary along the (residual) gauge orbit inside the first Gribov region?
- Investigate using the lattice
 - Not possible to access $G(0)$ on the lattice
 - Take lowest momentum instead: Finite-volume Artifact

Investigating the residual gauge orbit

- There are many copies along the residual gauge orbit

Investigating the residual gauge orbit

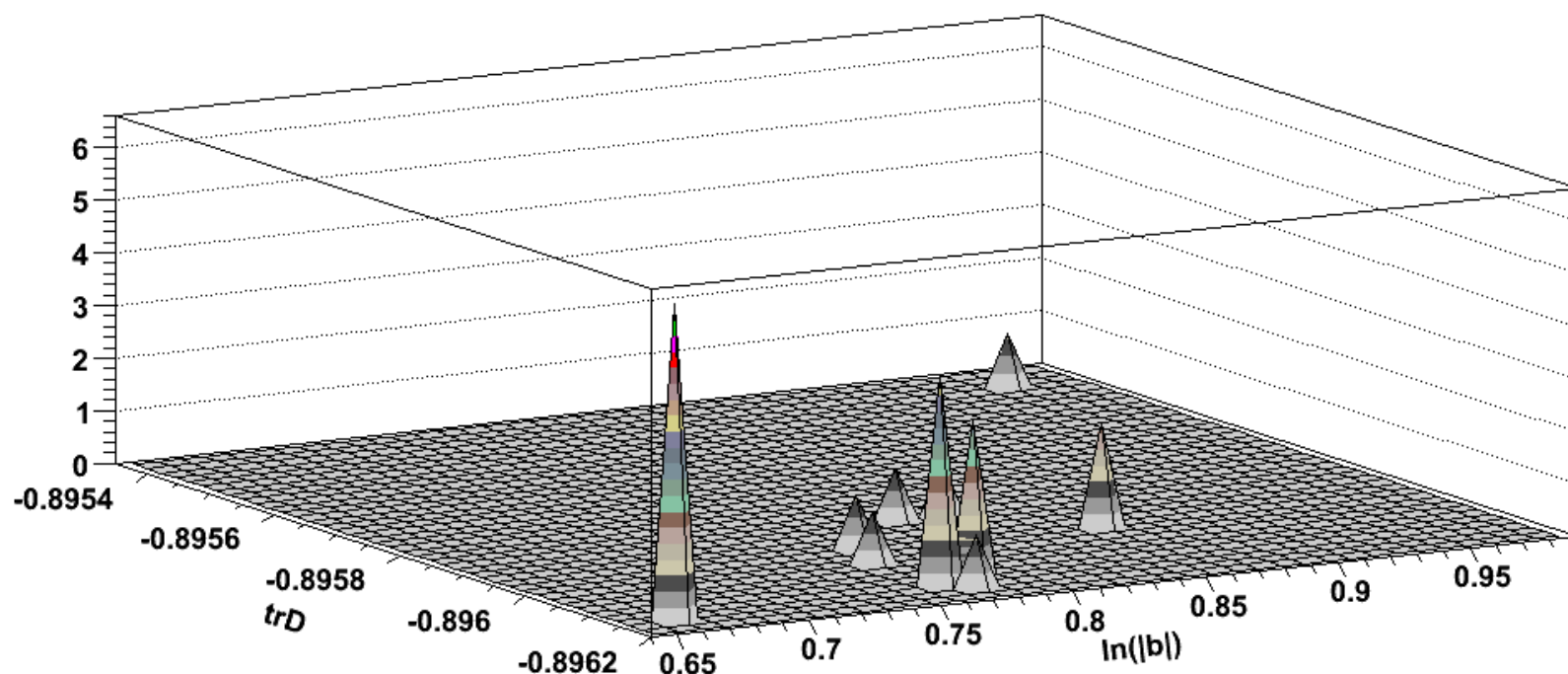
- There are many copies along the residual gauge orbit
- Is b different for each copy?
- If, how is b distributed?
- How is b correlated with the quantity $\text{Tr}D$?
- How evolves b with volume and discretization?
- Is there a possibility to use b as a gauge parameter?

Distribution of b [3d, beta=3.46, Maas, unpublished]

- At small volumes: Small number of *Gribov copies*

Distribution of b [3d, beta=3.46, Maas, unpublished]

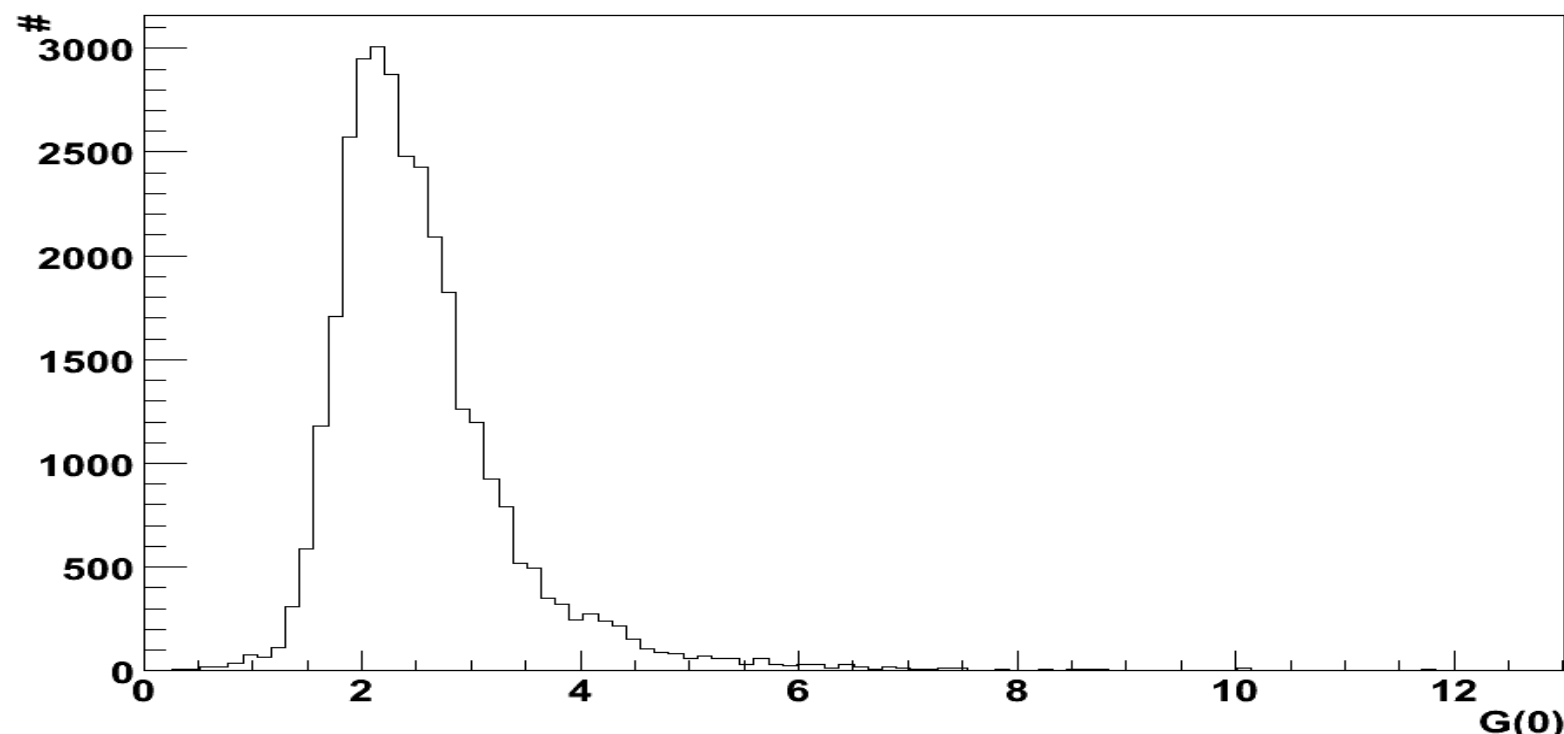
TrD vs. b for $V=(4.4 \text{ fm})^3$ (single configuration)



- At small volumes: Small number of **Gribov copies**
- b significantly different, if TrD is significantly different
 - Correspond to different **continuum Gribov copies**
 - b is relatively more different than TrD

Distribution of b on average [3d, beta=3.46, Maas, unpublished]

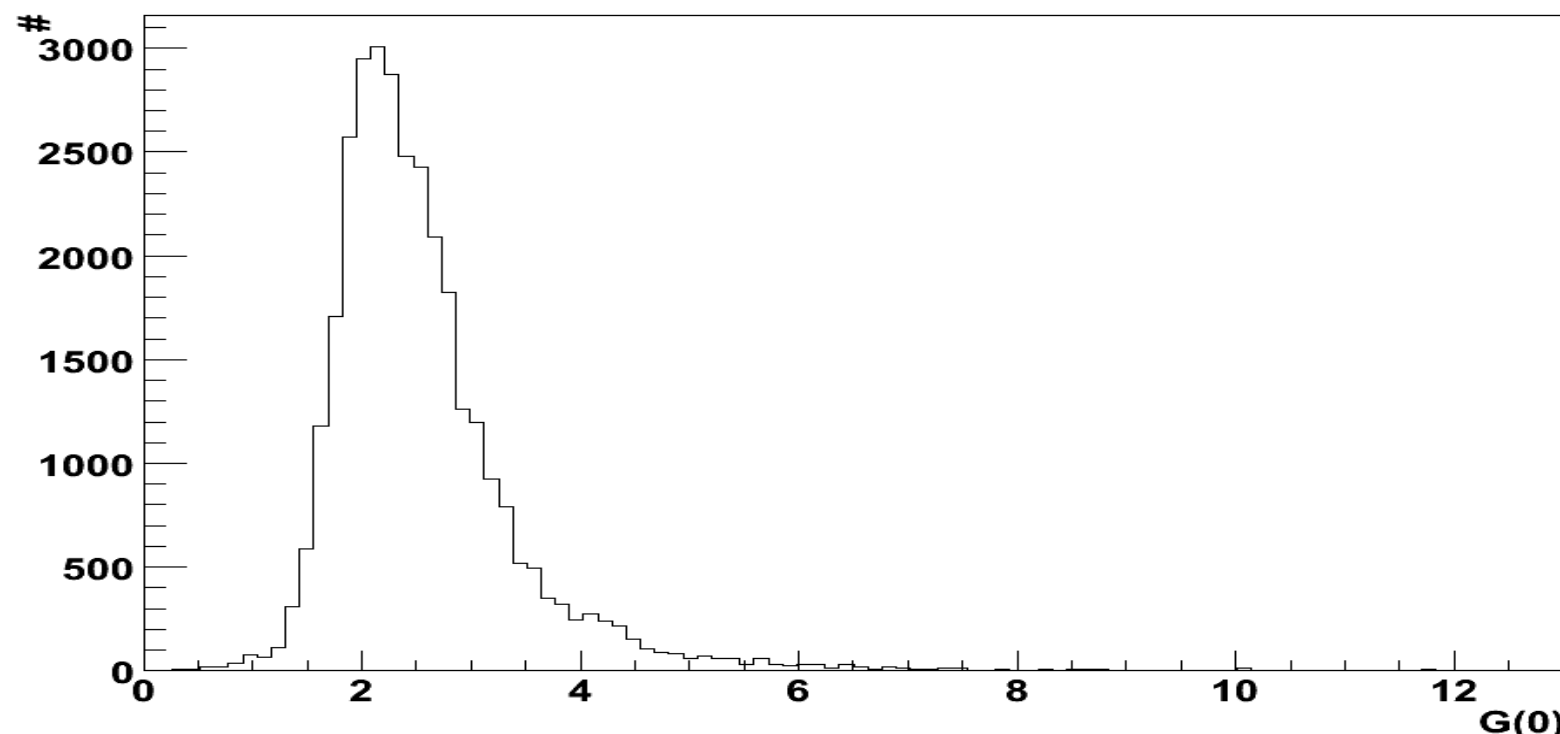
$G(0.395 \text{ GeV})/G(\infty \text{ GeV})$ for $V=(3.1 \text{ fm})^3$ from 1622 configurations



- Distribution is asymmetric with tail to large values
 - Extreme values for the **ghost propagator** are possible

Distribution of b on average [3d, beta=3.46, Maas, unpublished]

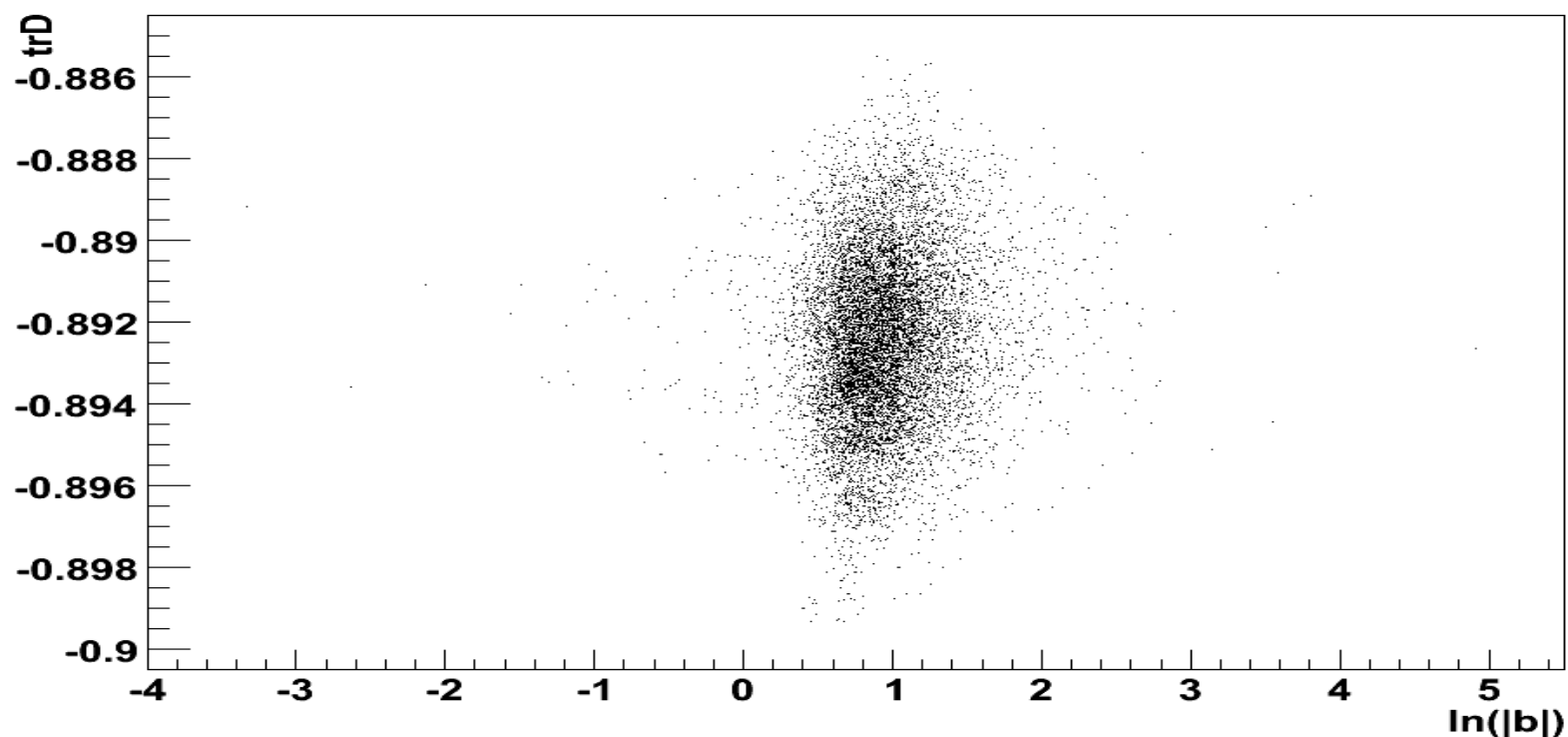
$G(0.395 \text{ GeV})/G(\infty \text{ GeV})$ for $V=(3.1 \text{ fm})^3$ from 1622 configurations



- Distribution is asymmetric with tail to large values
 - Extreme values for the **ghost propagator** are possible
- Is there a unique correlation of **b** and **$\text{Tr}D$** ?

Distribution of b on average [3d, beta=3.46, Maas, unpublished]

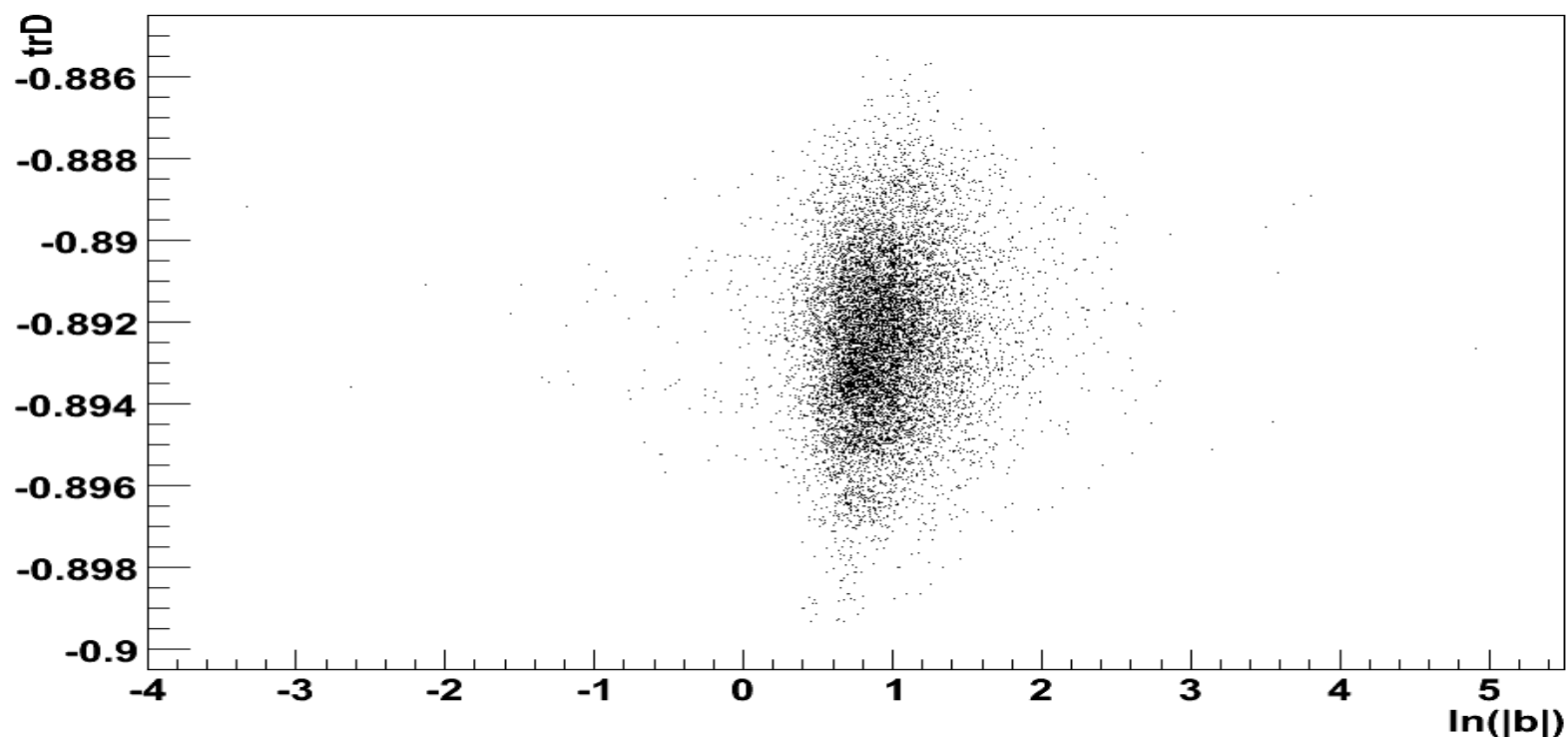
TrD vs. b for $V=(3.1 \text{ fm})^3$



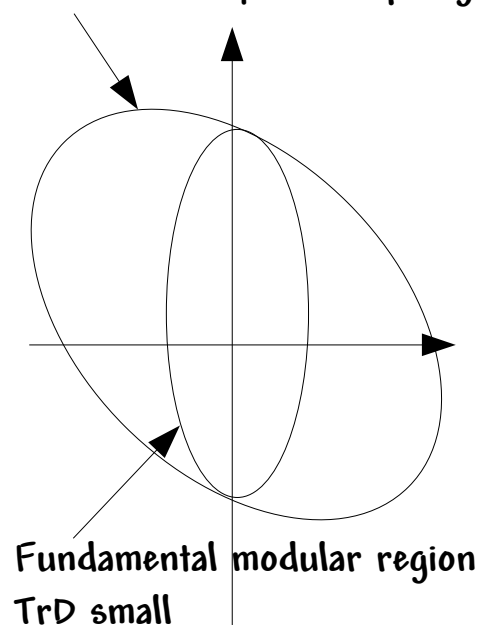
- No obvious correlation between both quantities

Distribution of b on average [3d, beta=3.46, Maas, unpublished]

TrD vs. b for $V=(3.1 \text{ fm})^3$



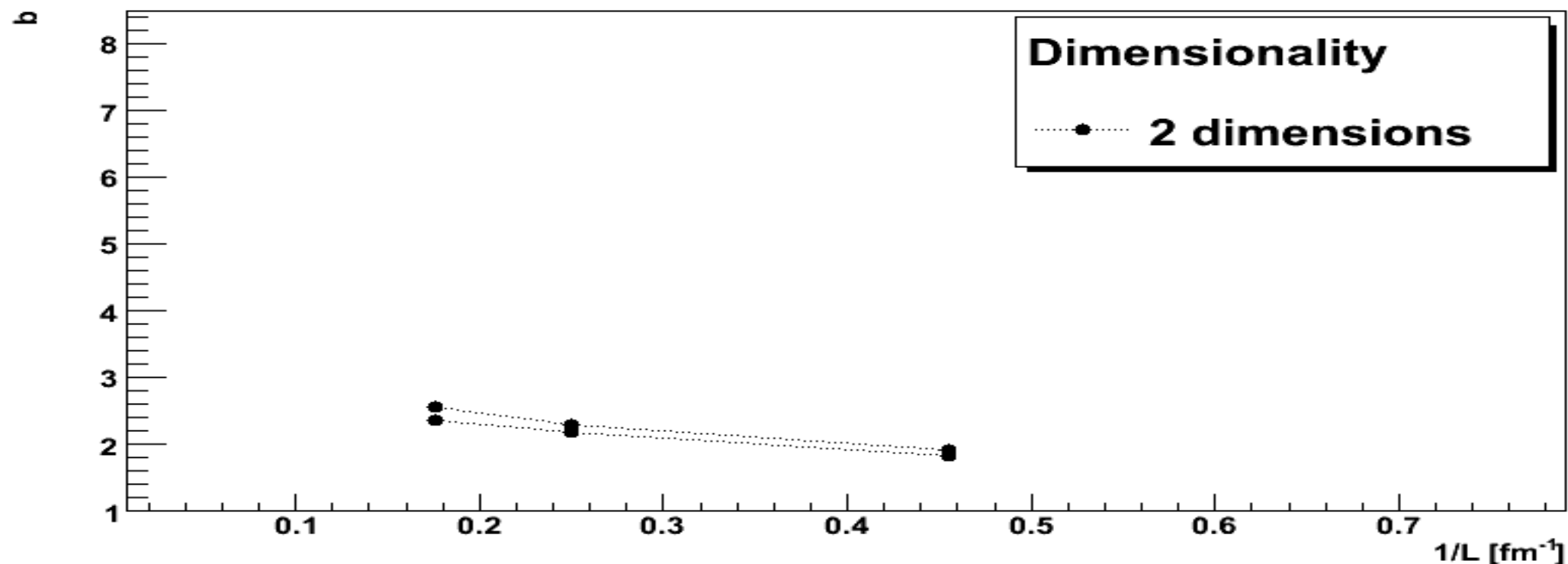
Gribov horizon: b potentially large



- No obvious correlation between both quantities
 - Only a small tendency that small b (far away from the Gribov horizon) correlate with small TrD (fundamental modular region)

The b gauge corridor [Maas, unpublished]

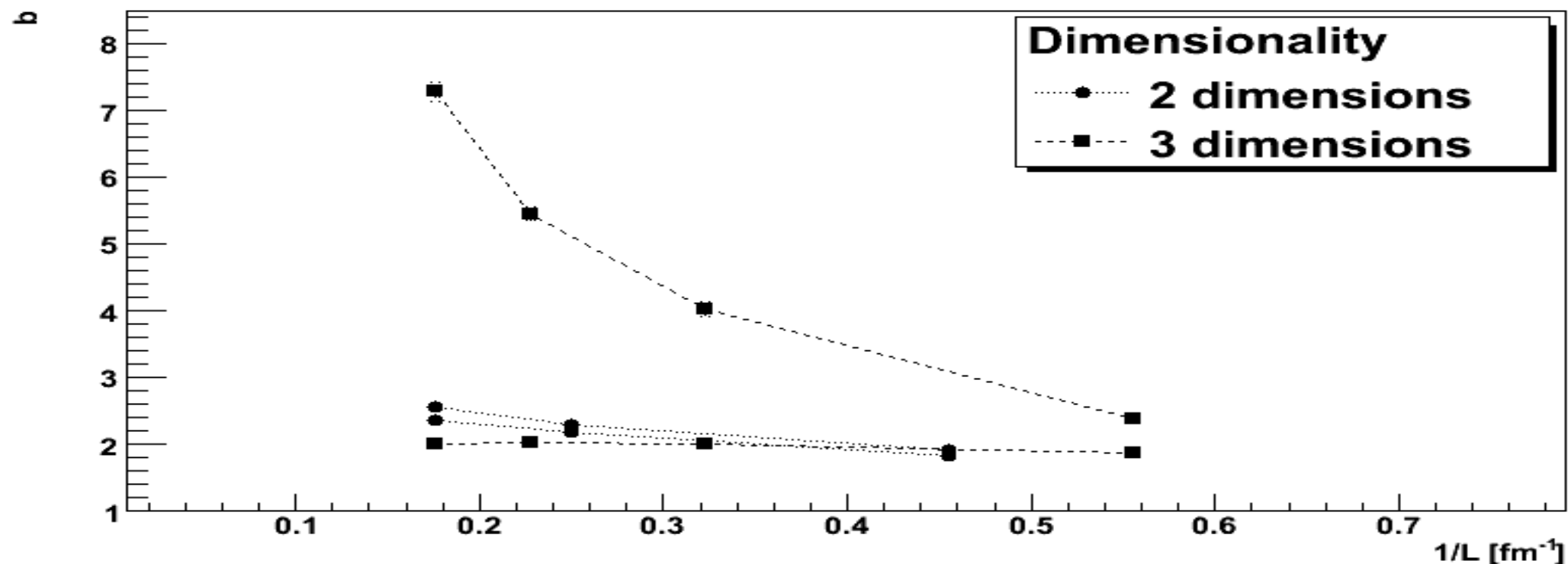
b corridor



- The **corridor** [min b , max b] opens with volume
 - Overall scale partly a **finite-volume effect**
 - **max b /min b** not: Possible range of b values increases
 - Can **max b** diverge in the thermodynamic limit?

The b gauge corridor [Maas, unpublished]

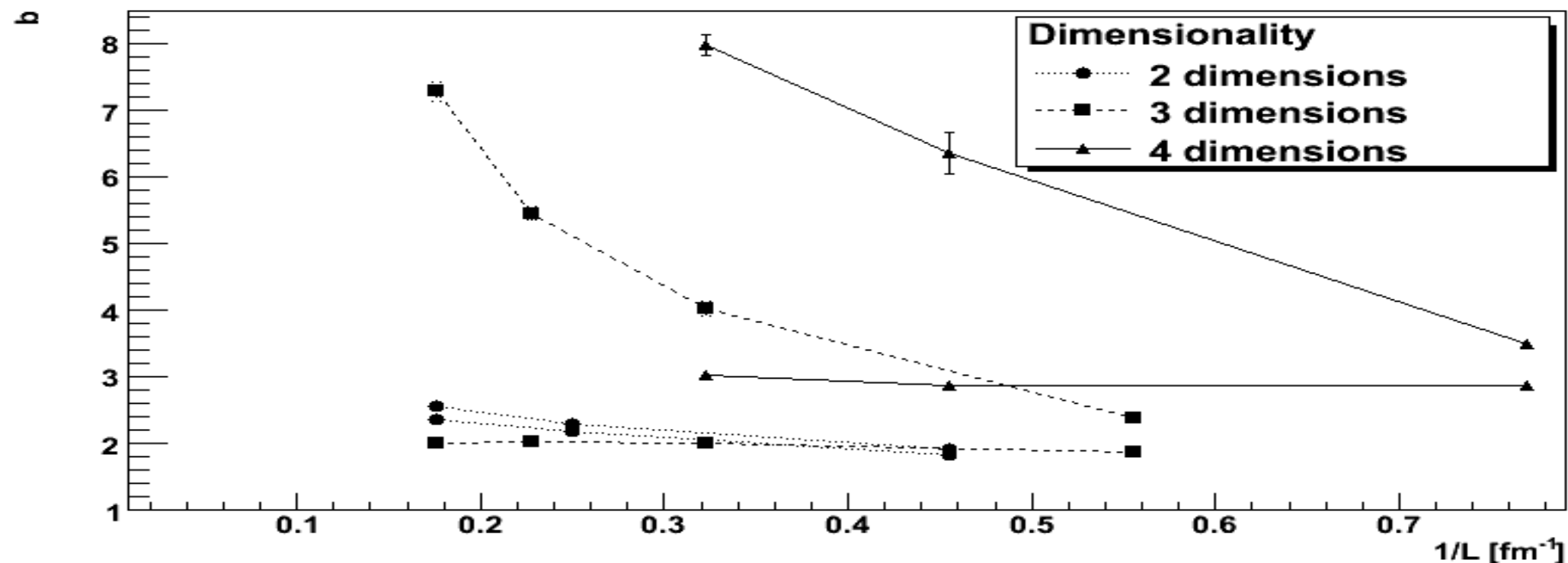
b corridor



- The **corridor** [min b , max b] opens with volume
 - Overall scale partly a **finite-volume effect**
 - **max b /min b** not: Possible range of b values increases
 - Can **max b** diverge in the thermodynamic limit?

The b gauge corridor [Maas, unpublished]

b corridor



- The **corridor** [min b , max b] opens with volume
 - Overall scale partly a **finite-volume effect**
 - **max b /min b** not: Possible range of b values increases
 - Can **max b** diverge in the thermodynamic limit?

Constructing a gauge parameter from b

- Volumes increases: b distribution becomes dense

Constructing a gauge parameter from b

- Volumes increases: b distribution becomes dense
- Constructing a gauge
 - Select from the range $[\min b, \max b]$ a value B (for a given volume)
 - Select on each residual orbit the copy with b closest to B
 - Average over these copies to obtain the gauge-fixed correlation functions in the Landau- B gauge

Constructing a gauge parameter from b

- Volumes increases: b distribution becomes dense
- Constructing a gauge
 - Select from the range $[\min b, \max b]$ a value B (for a given volume)
 - Select on each residual orbit the copy with b closest to B
 - Average over these copies to obtain the gauge-fixed correlation functions in the Landau- B gauge
 - In addition: $\min B$ and $\max B$ gauges, selecting the extremal values of b on each residual gauge orbit

Compare correlation functions between Landau gauges

- Landau- \mathcal{B} gauges
 - Here: $\min\mathcal{B}$ und $\max\mathcal{B}$ -gauges

Compare correlation functions between Landau gauges

- Landau- \mathcal{B} gauges
 - Here: $\min\mathcal{B}$ und $\max\mathcal{B}$ -gauges
- Absolute Landau gauge

Compare correlation functions between Landau gauges

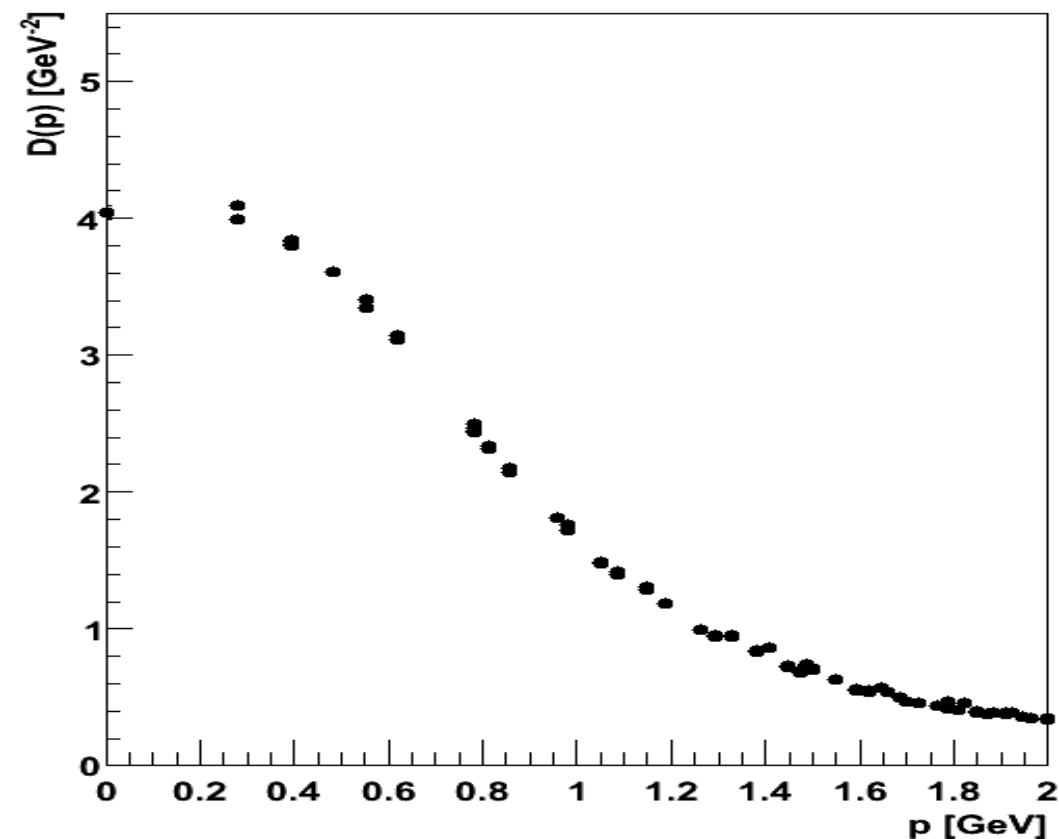
- Landau- \mathcal{B} gauges
 - Here: $\min\mathcal{B}$ und $\max\mathcal{B}$ -gauges
- Absolute Landau gauge
- Inverse Landau gauge
 - Maximize $\text{Tr}D$

Compare correlation functions between Landau gauges

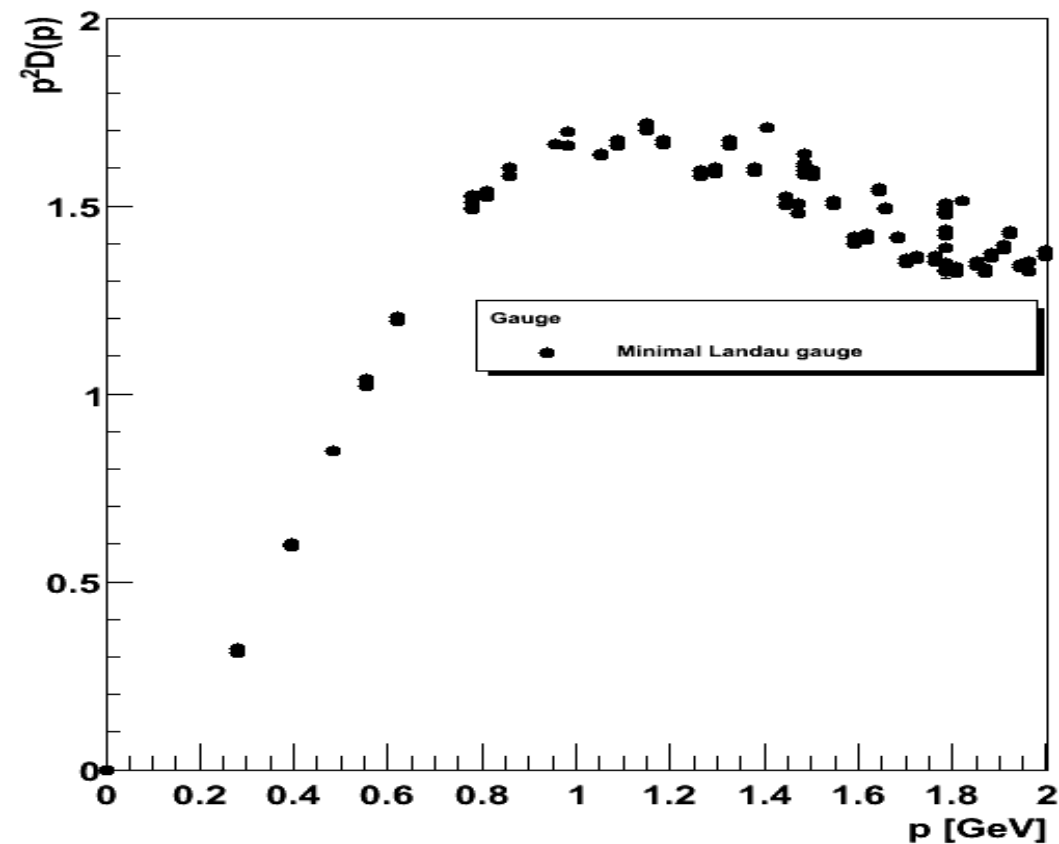
- Landau-B gauges
 - Here: $\min B$ und $\max B$ -gauges
- Absolute Landau gauge
- Inverse Landau gauge
 - Maximize $\text{Tr} D$
- Minimal Landau gauge
 - Select a random copy on each residual gauge orbit
 - Probability distribution is determined by the underlying algorithm
 - Differ for different algorithms, but cheap compared to all other gauges
 - No direct analogue in or translation to continuum field theory
 - Current standard in lattice calculations

Gluon propagator in the minimal Landau gauge [3d, Maas, unpublished]

Gluon propagator



Gluon dressing function



- Mild infrared suppression in 3d (strong in 2d, none in 4d)

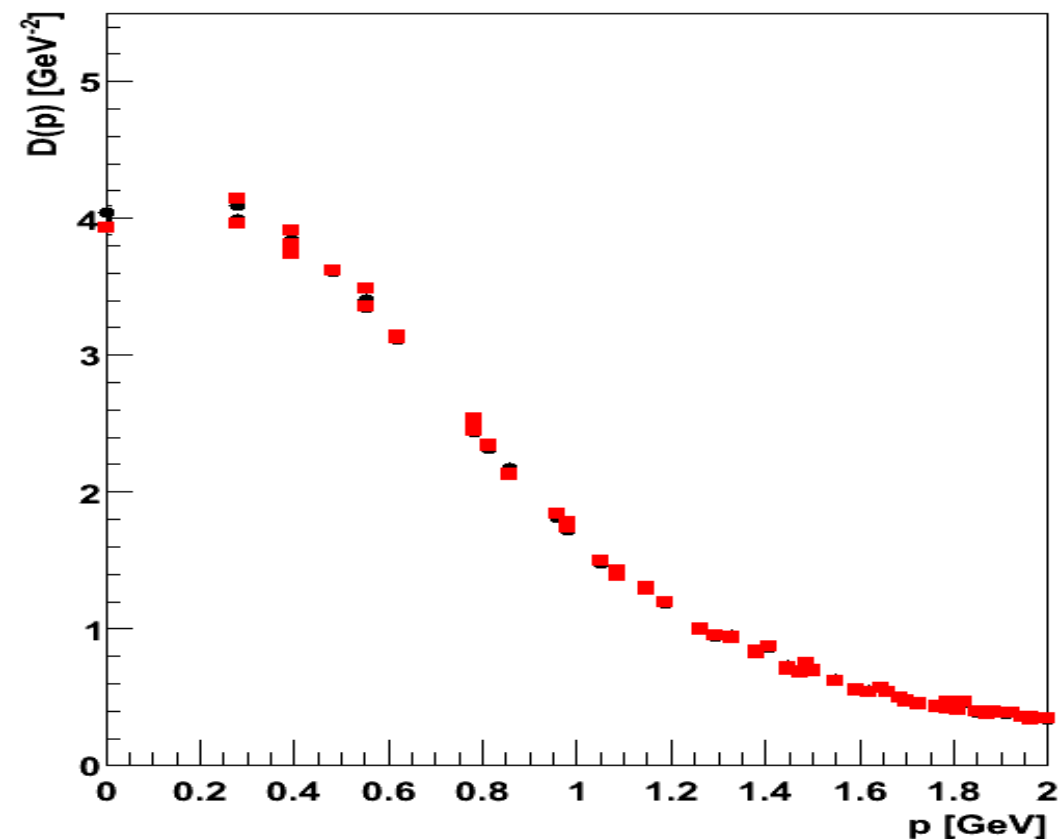
- Infrared constant in the infinite-volume/continuum limit

[Cucchieri et al. 2007/8, von Smekal et al. 2008]

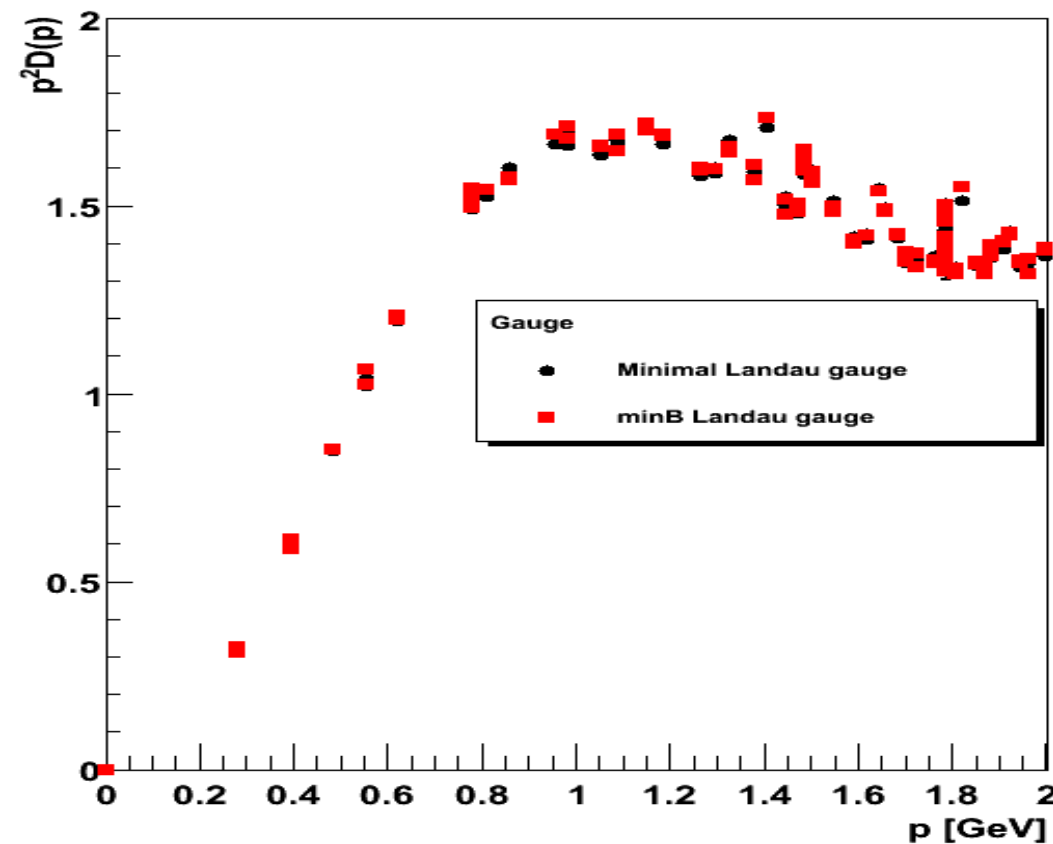
- Possibly vanishing in 2d [Maas, 2007, Cucchieri 2007/8]

Gluon propagator in the minB Landau gauge [3d, Maas, unpublished]

Gluon propagator



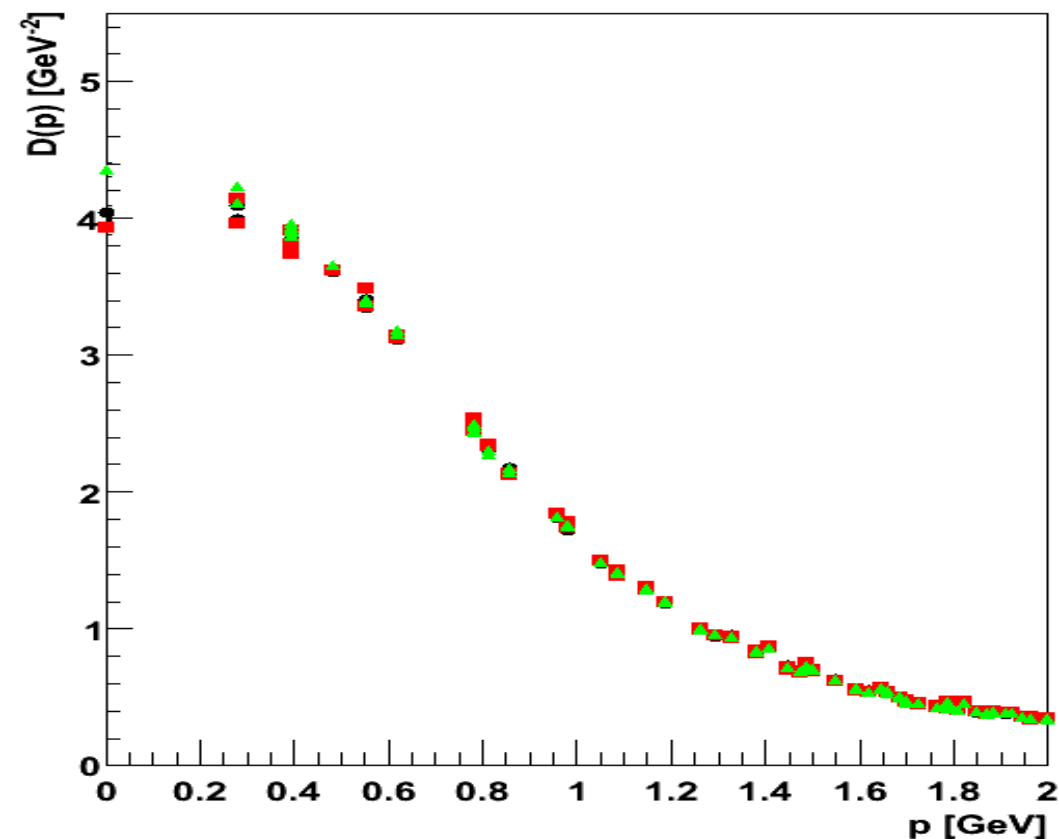
Gluon dressing function



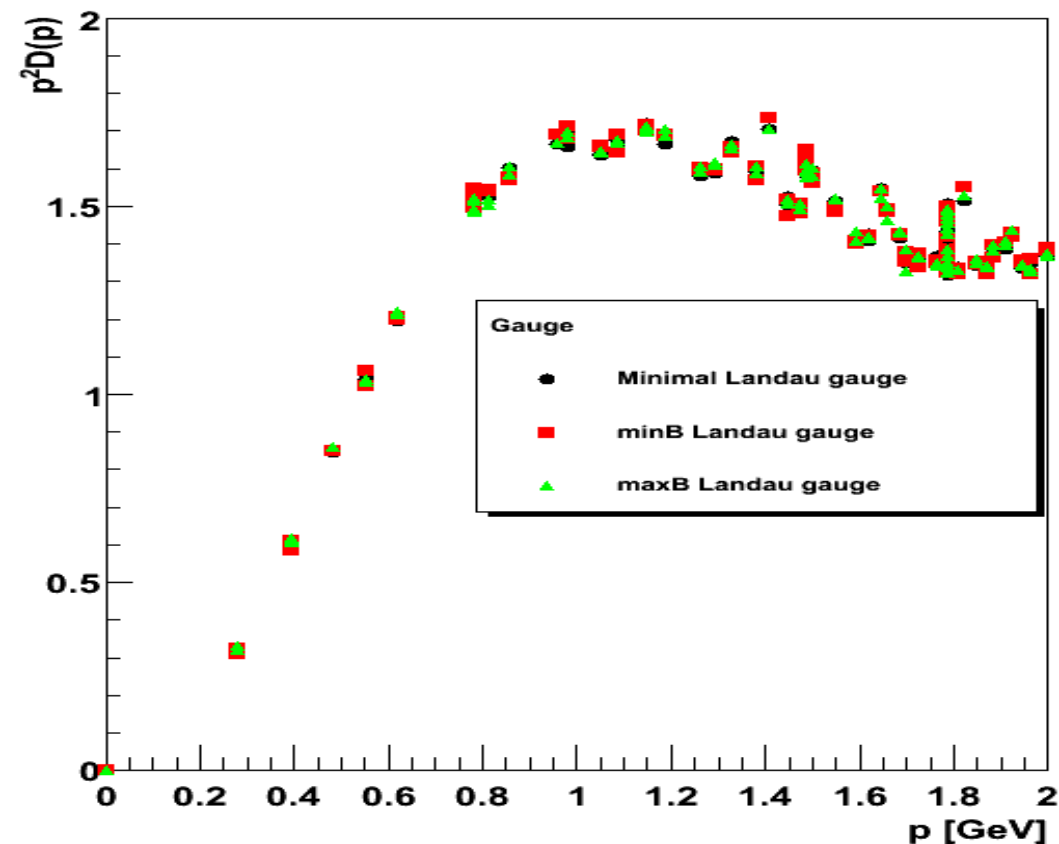
- At current volumes little difference to minimal Landau gauge

Gluon propagator in the maxB Landau gauge [3d, Maas, unpublished]

Gluon propagator



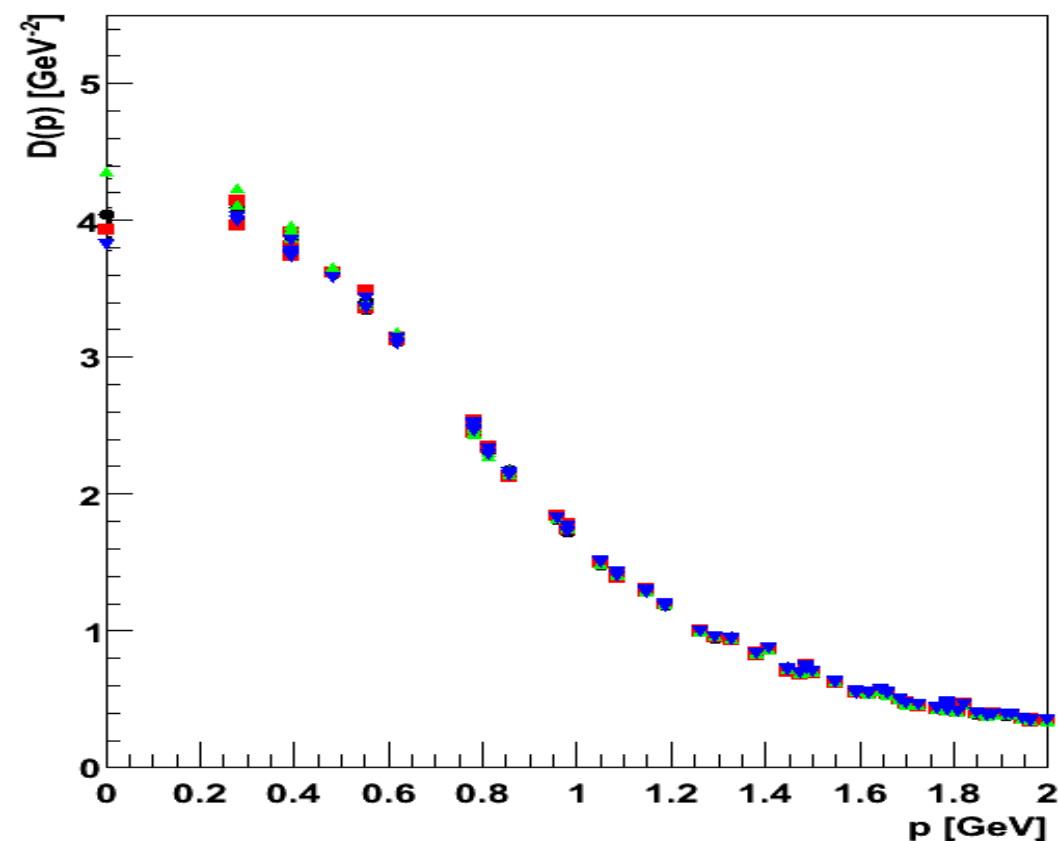
Gluon dressing function



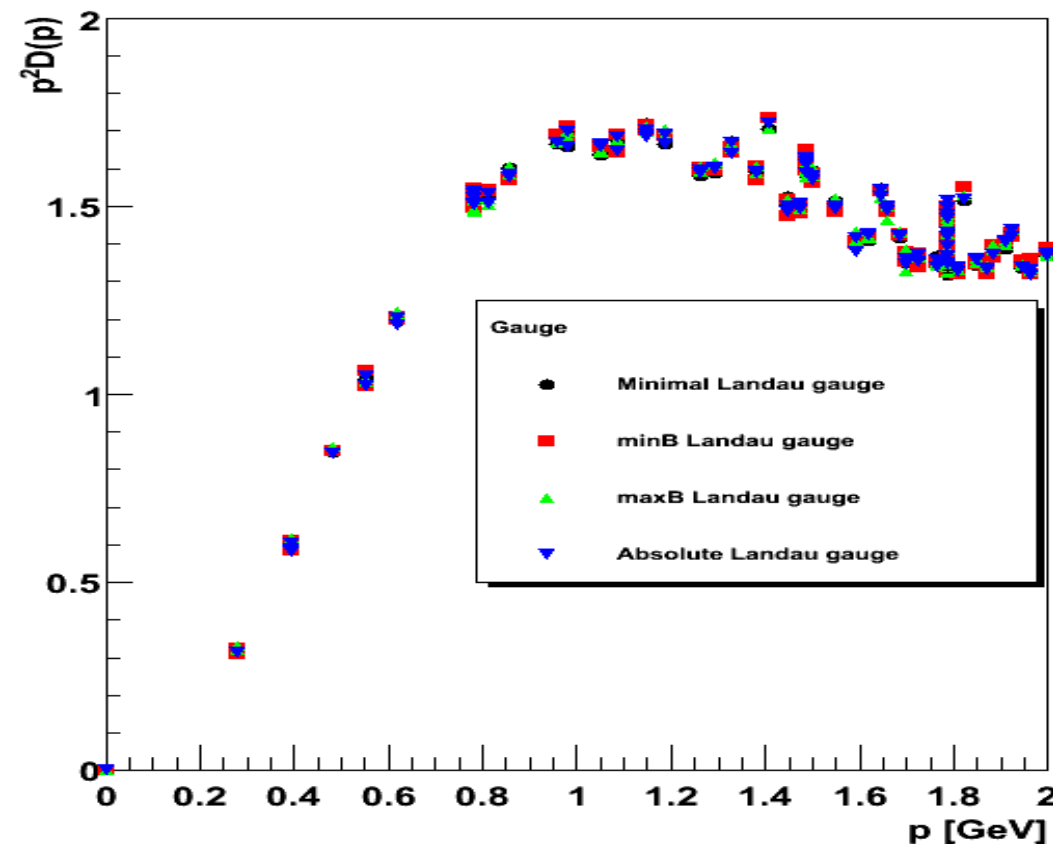
- At current volumes little difference to minimal Landau gauge
 - A little less infrared suppressed

Gluon propagator in the absolute Landau gauge [3d, Maas, unpublished]

Gluon propagator



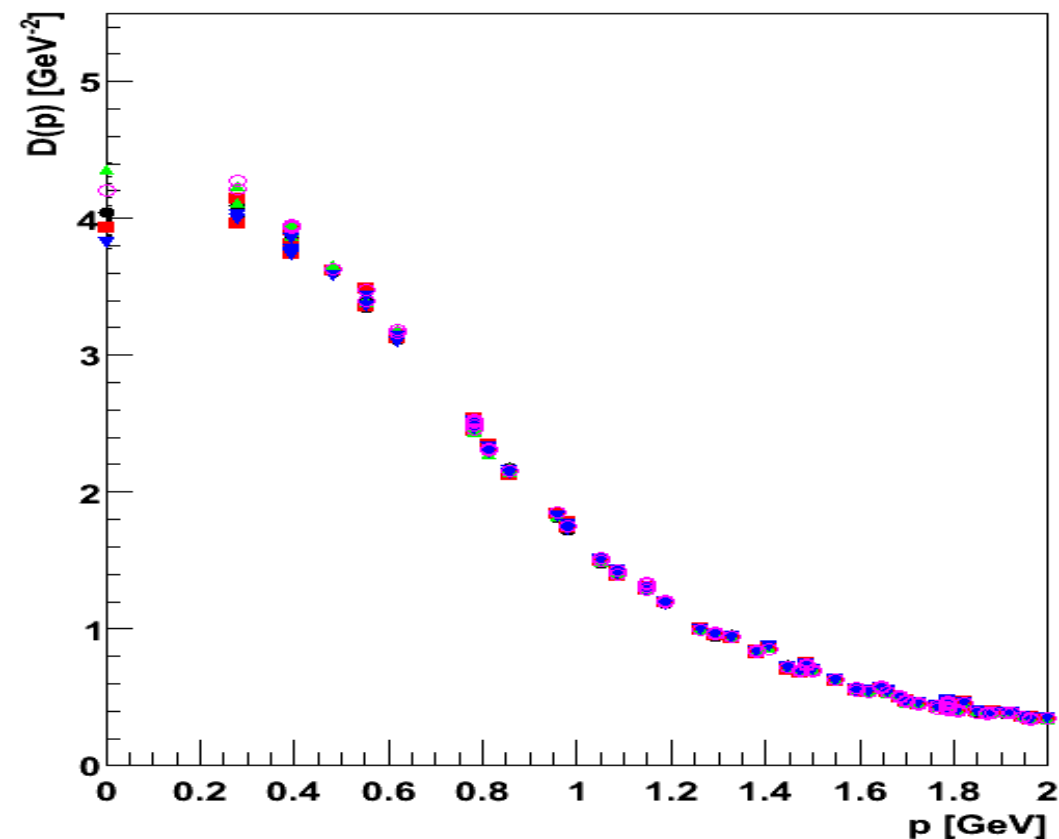
Gluon dressing function



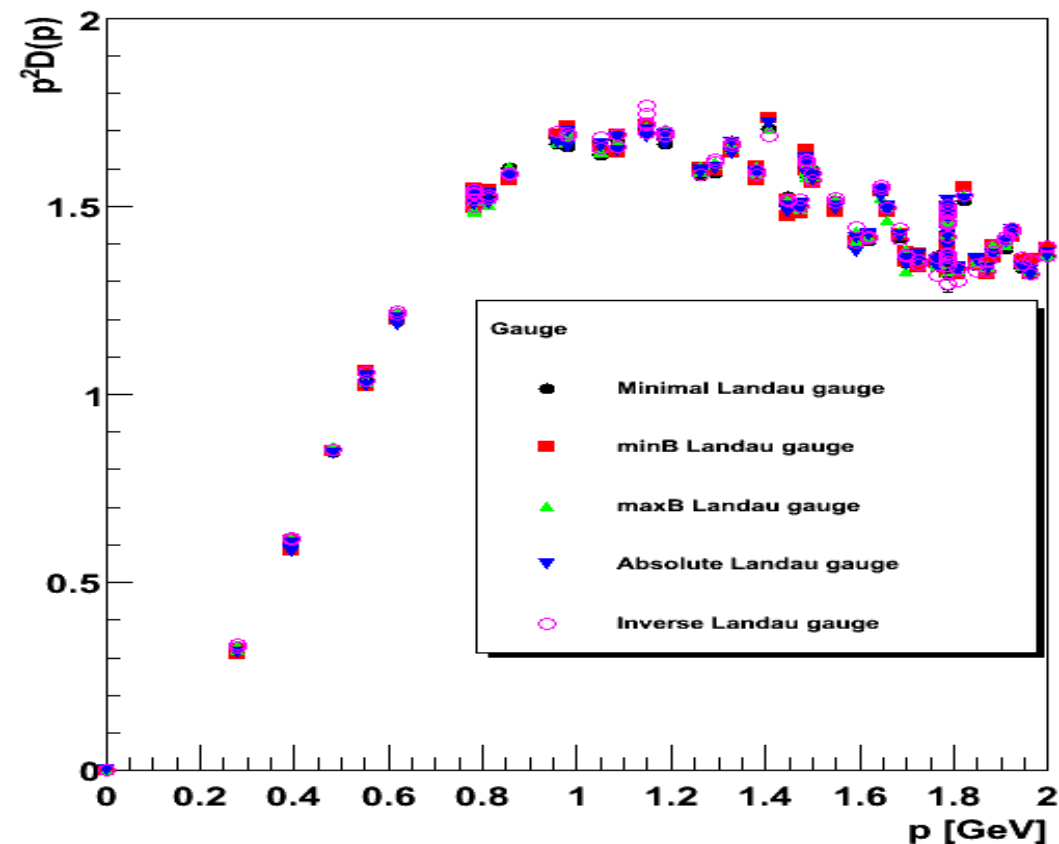
- At current volumes little difference to minimal Landau gauge
 - Somewhat stronger infrared suppressed
 - Unclear whether it vanishes

Gluon propagator in the inverse Landau gauge [3d, Maas, unpublished]

Gluon propagator



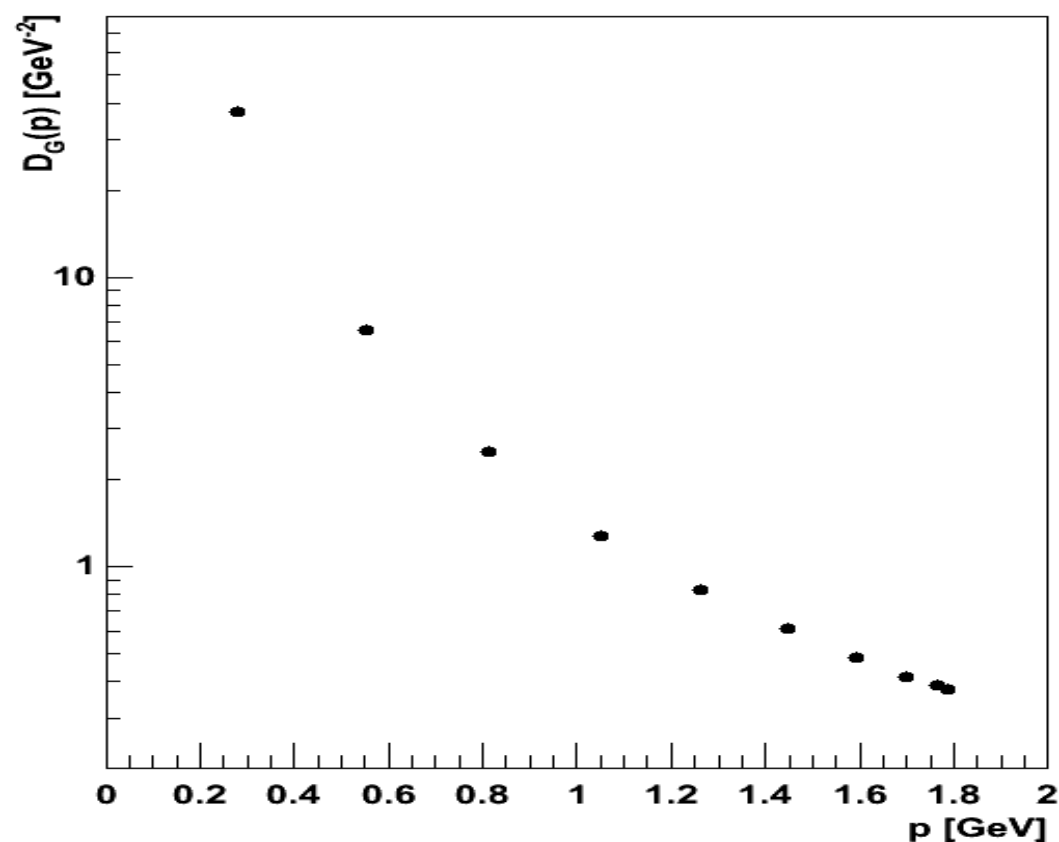
Gluon dressing function



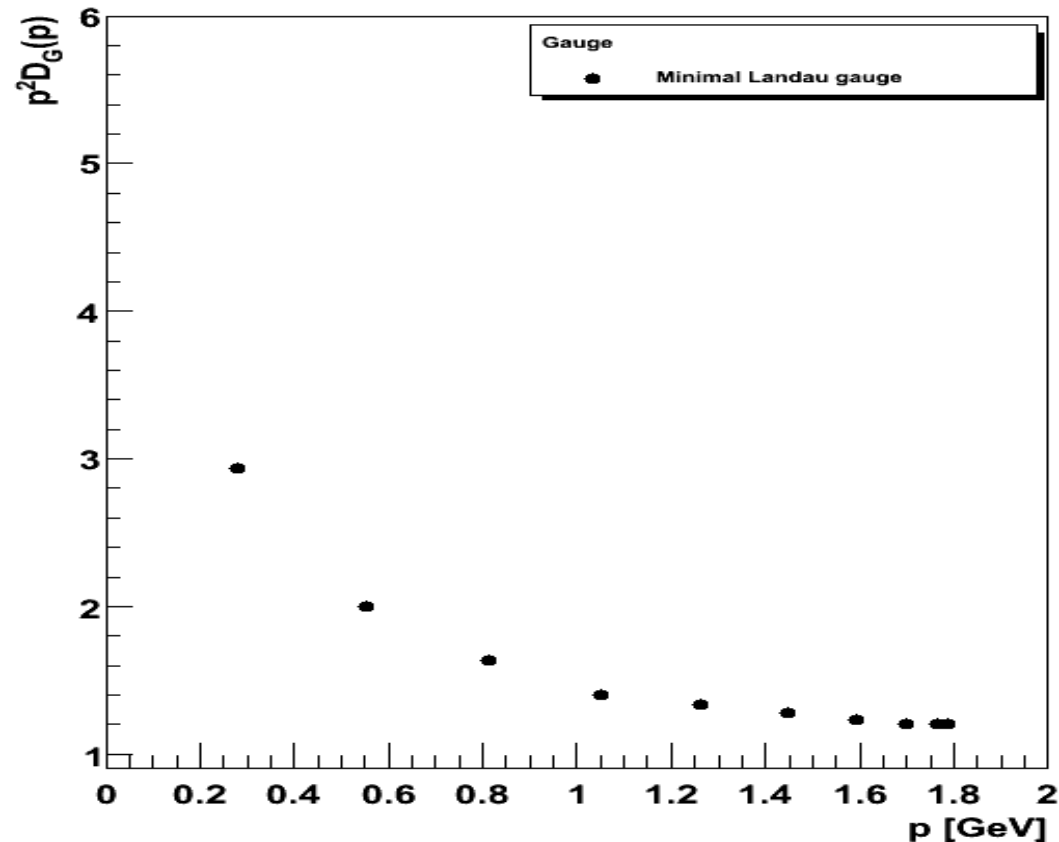
- At current volumes little difference to minimal Landau gauge

Ghost propagator in the minimal Landau gauge [3d, Maas, unpublished]

Ghost propagator



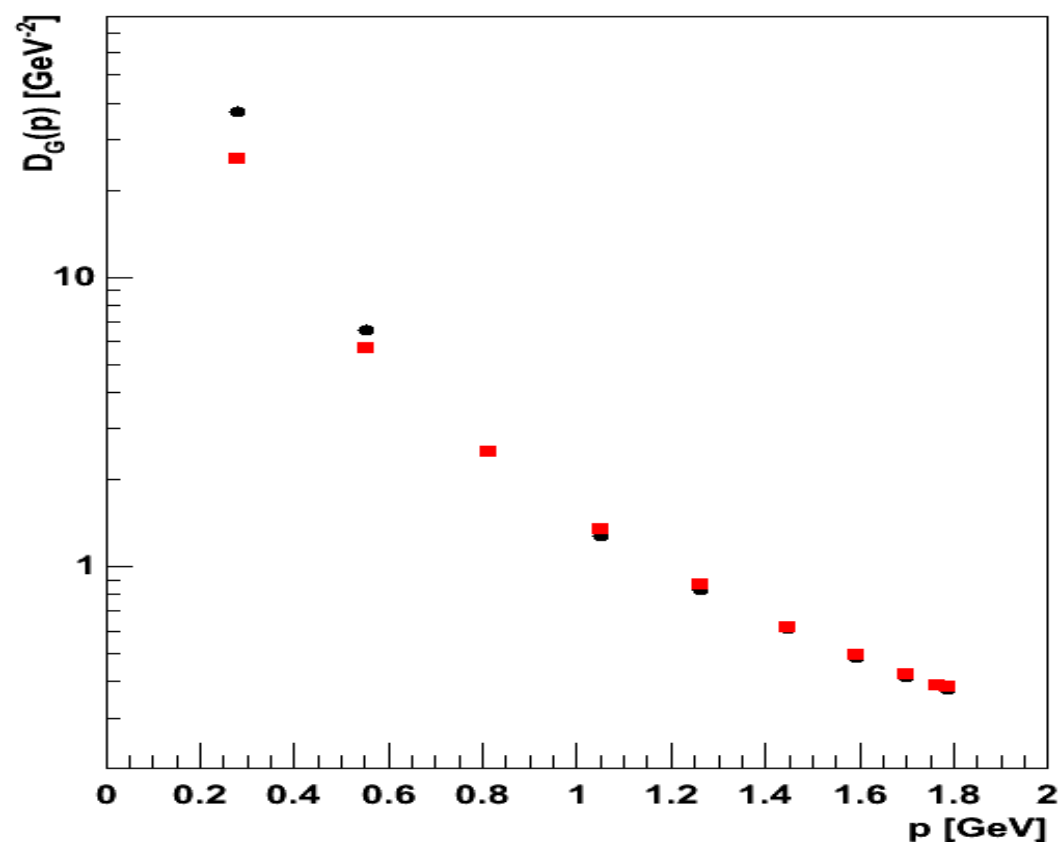
Ghost dressing function



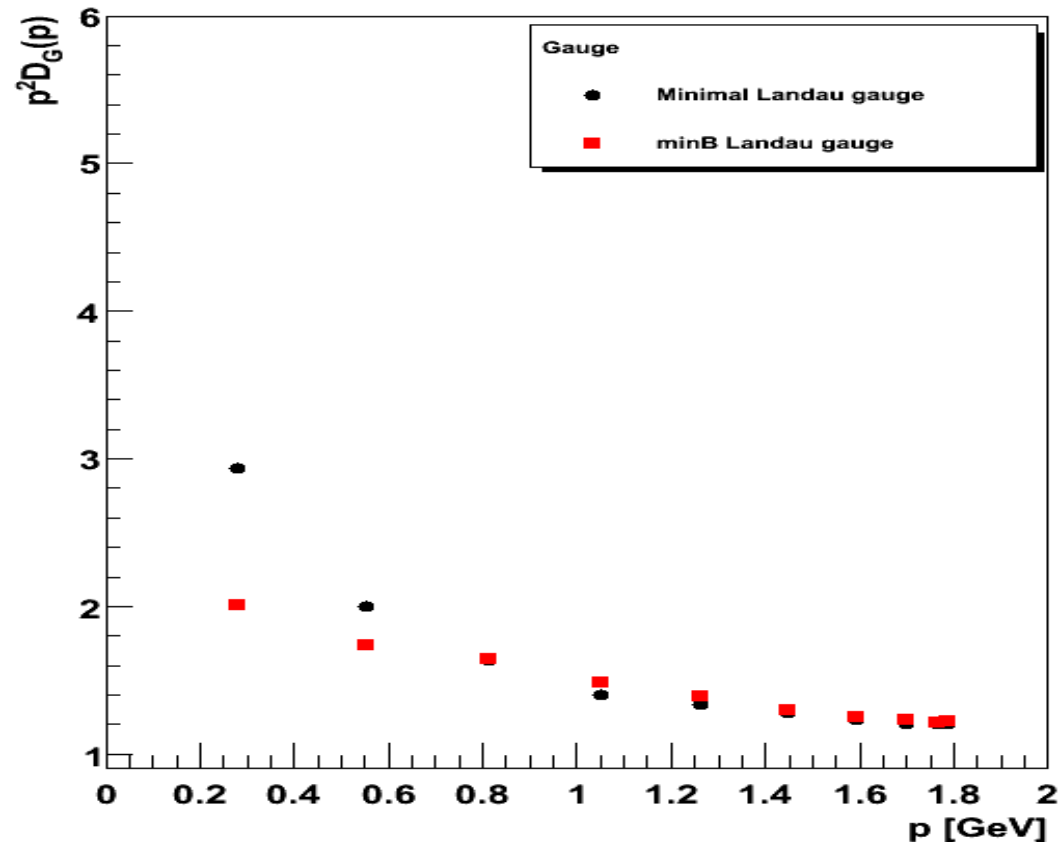
- Infrared enhanced (in all dimensions)
 - In 3d and 4d infrared finite dressing function [Cucchieri et al. 2007/8, von Smekal 2007]
 - In 2d possibly diverging [Maas 2007, Cucchieri et al. 2007/8]

Ghost propagator in the minB Landau gauge [3d, Maas, unpublished]

Ghost propagator



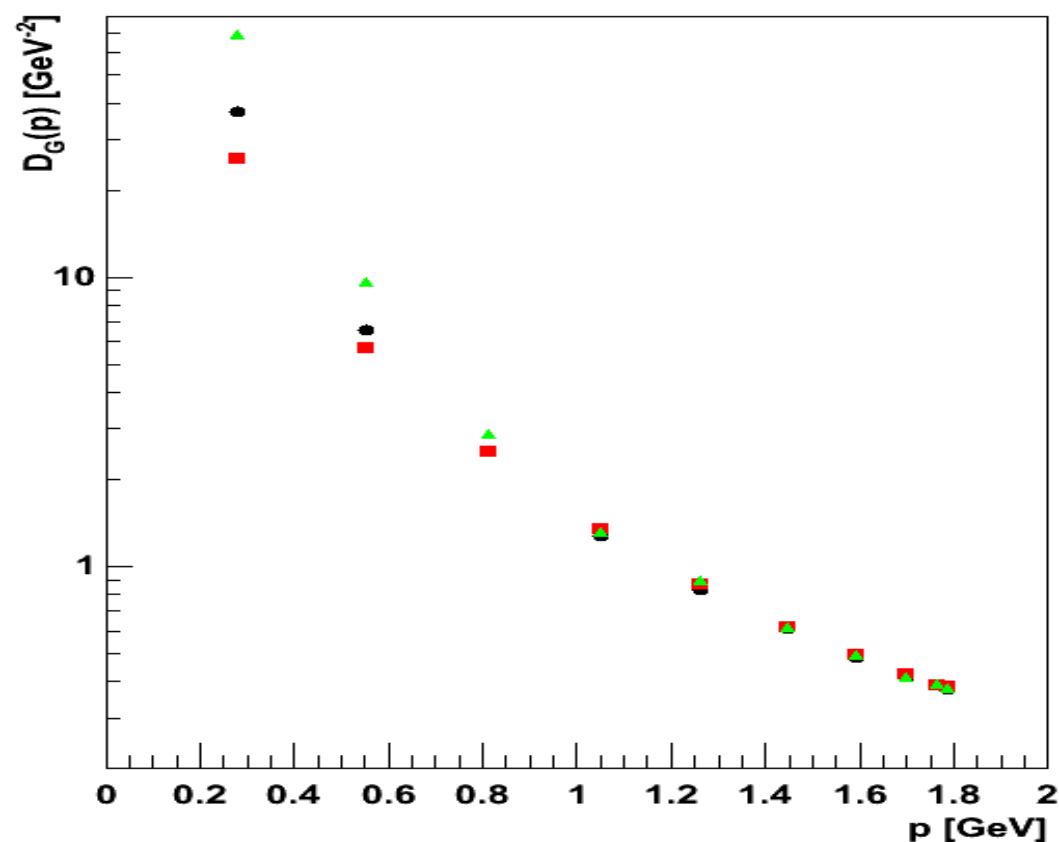
Ghost dressing function



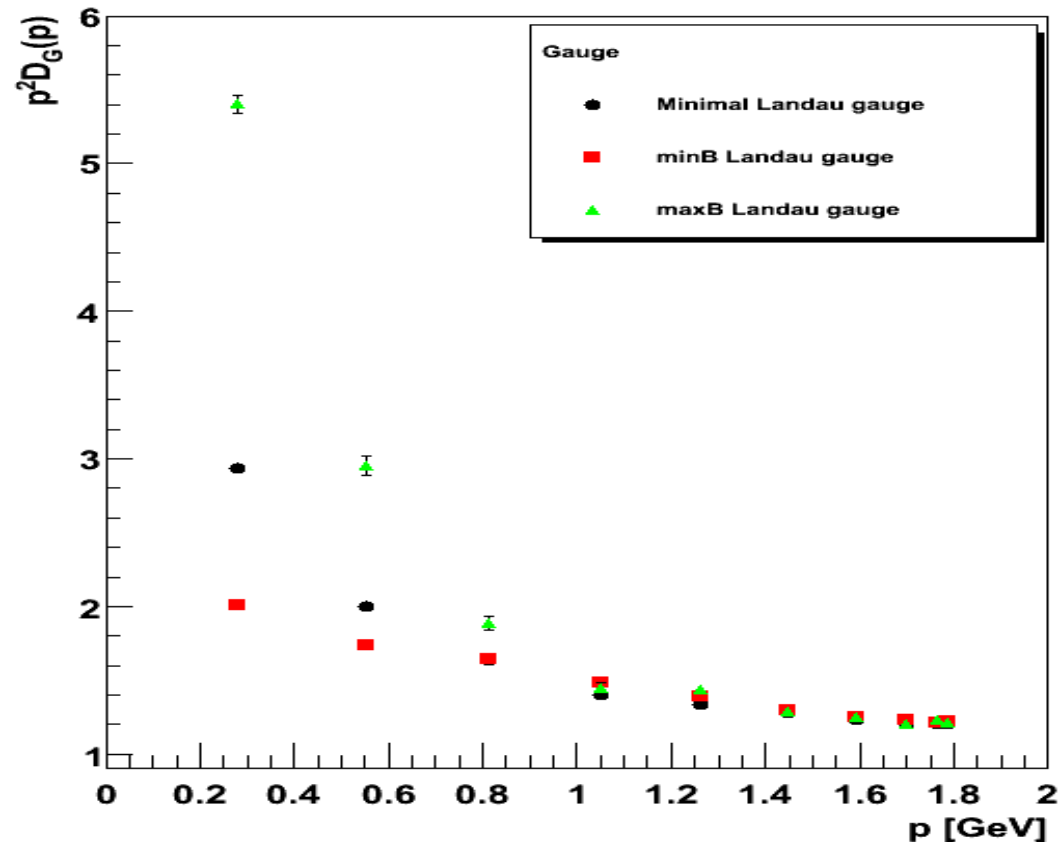
- Significantly less enhanced
 - Will have in 3d and 4d infrared finite dressing function
 - Fate in 2d yet unclear

Ghost propagator in the minimal Landau gauge [3d, Maas, unpublished]

Ghost propagator



Ghost dressing function

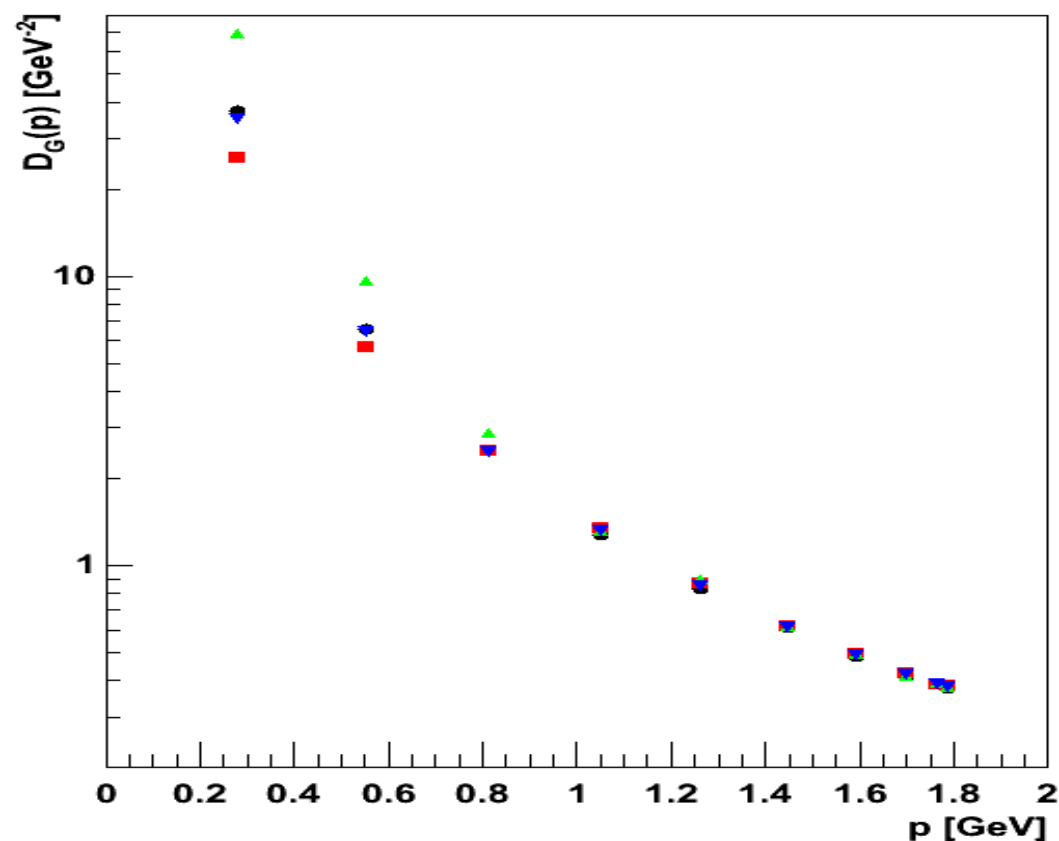


- Strongly enhanced

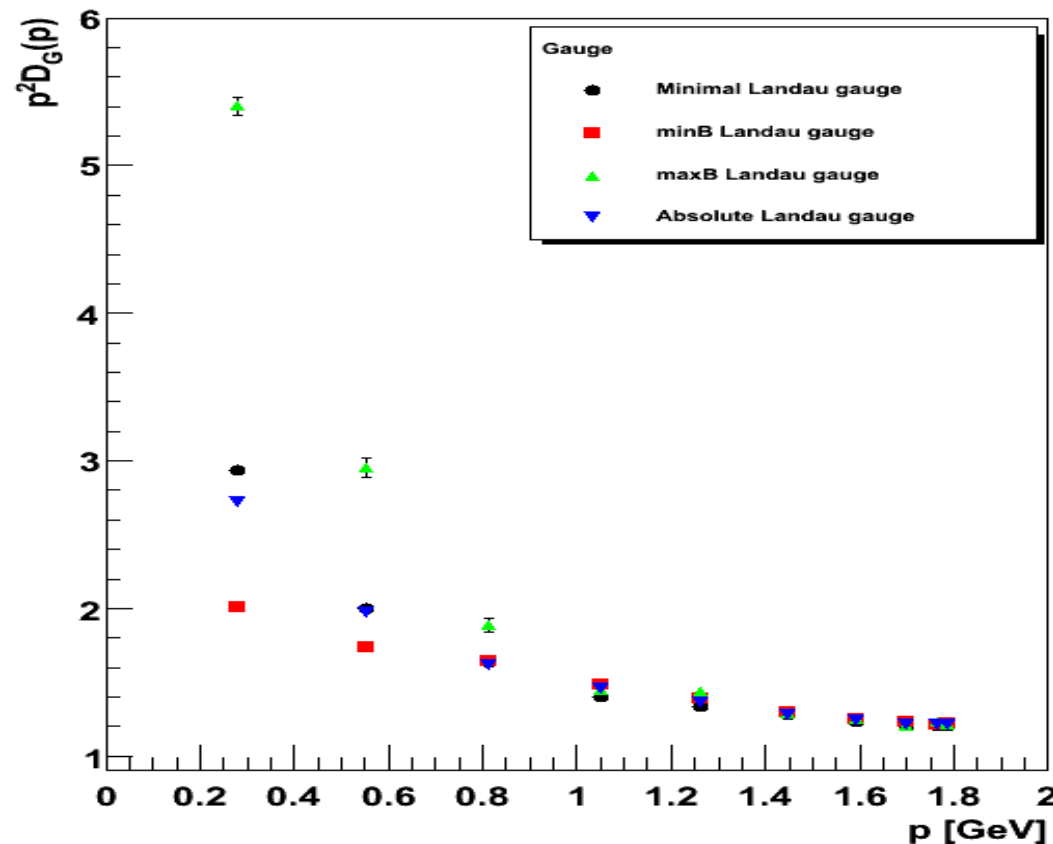
- Could have infrared divergent dressing function in all d
- Effects up to almost 1 GeV

Ghost propagator in the minimal Landau gauge [3d, Maas, unpublished]

Ghost propagator



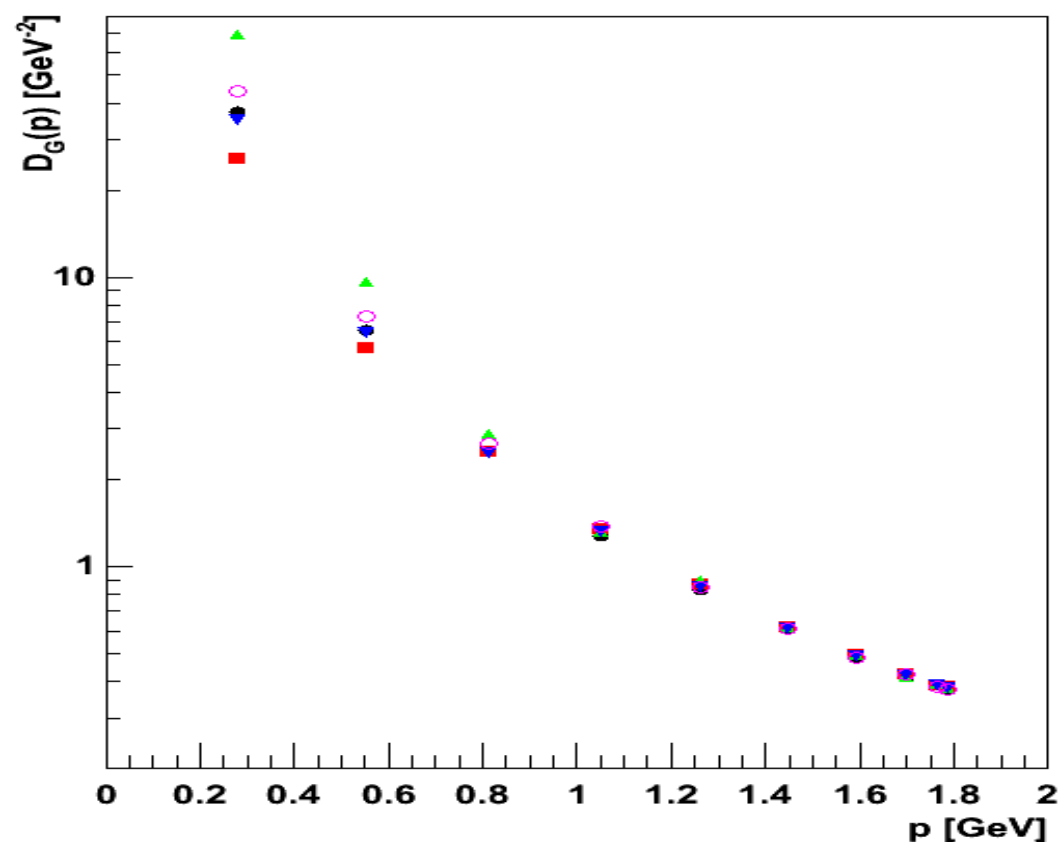
Ghost dressing function



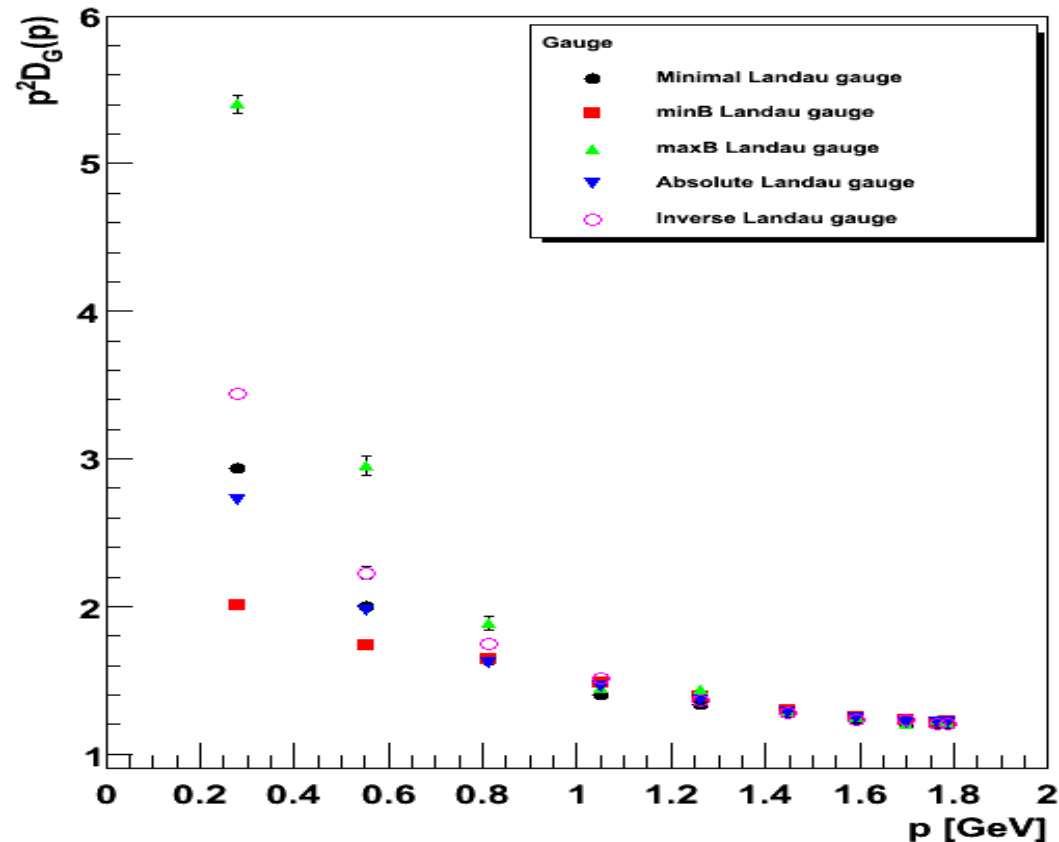
- Less strongly enhanced than minimal Landau gauge
 - But has different finite volume evolution [Maas, 2008]

Ghost propagator in the minimal Landau gauge [3d, Maas, unpublished]

Ghost propagator



Ghost dressing function



- Mildly stronger enhanced than minimal Landau gauge
 - Infinite volume fate is unclear

B gauges

- **Gluon propagator** is only mildly affected
 - Consequence of Zwanziger's entropy conjecture?

B gauges

- **Gluon propagator** is only mildly affected
 - Consequence of Zwanziger's entropy conjecture?
- **Ghost propagator** differs strongly
 - Zwanziger's entropy conjecture is not applying

B gauges

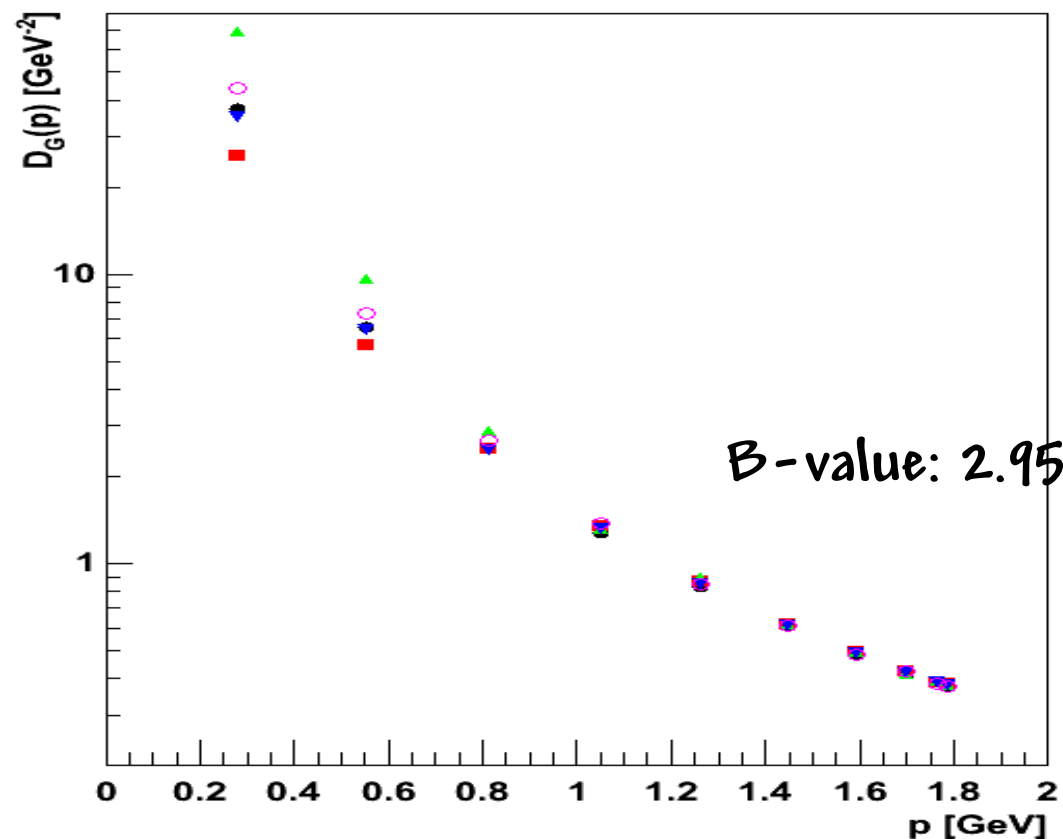
- **Gluon propagator** is only mildly affected
 - Consequence of Zwanziger's entropy conjecture?
- **Ghost propagator** differs strongly
 - Zwanziger's entropy conjecture is not applying
- **B-Landau gauges** permit unique gauge-fixing prescription

B gauges

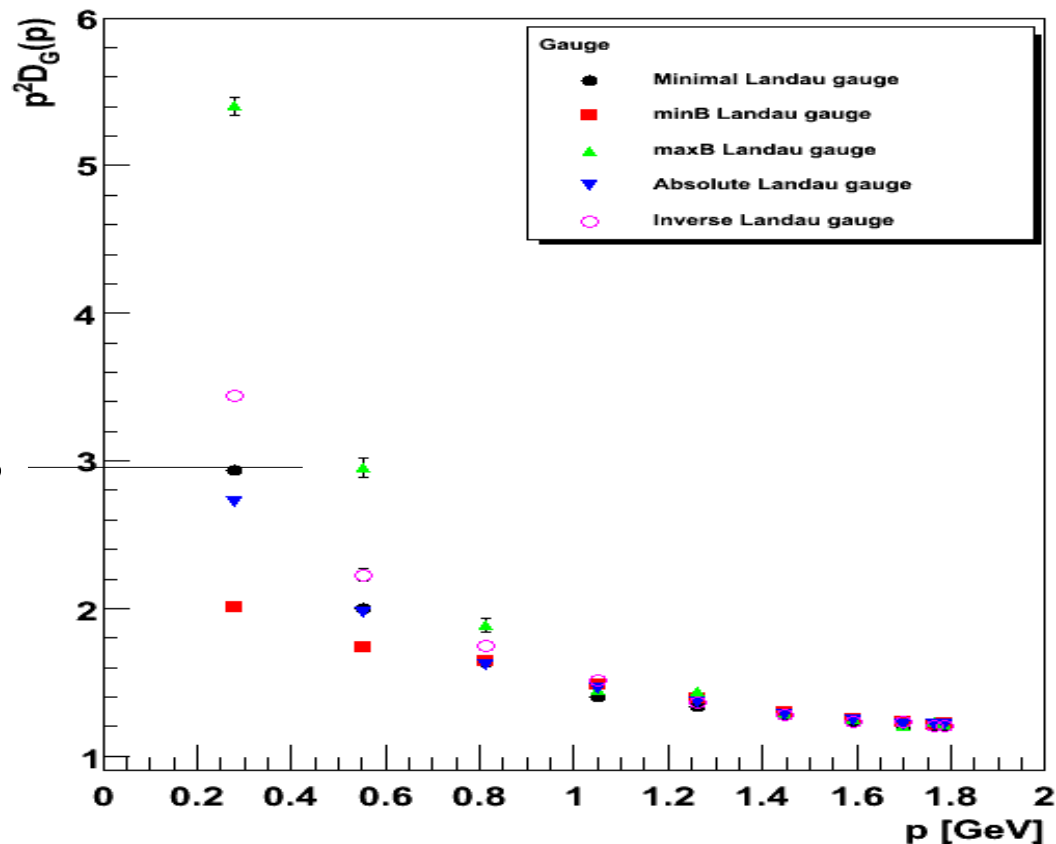
- **Gluon propagator** is only mildly affected
 - Consequence of Zwanziger's entropy conjecture?
- **Ghost propagator** differs strongly
 - Zwanziger's entropy conjecture is not applying
- **B-Landau gauges** permit unique gauge-fixing prescription
 - Are **correlation functions** in a **B gauge** identical to e.g. the **minimal Landau gauge** with its effective **b** value?
 - If it is a faithful map of the probability distribution

Mapping gauge descriptions [3d, Maas, unpublished]

Ghost propagator



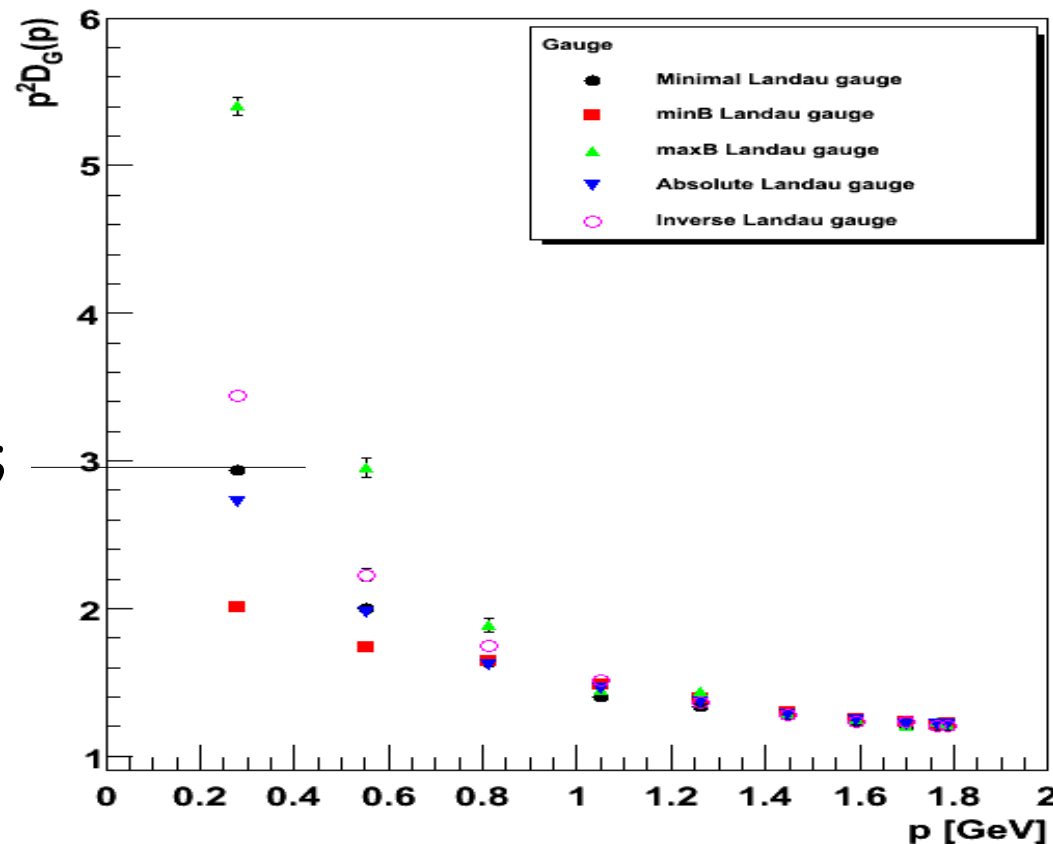
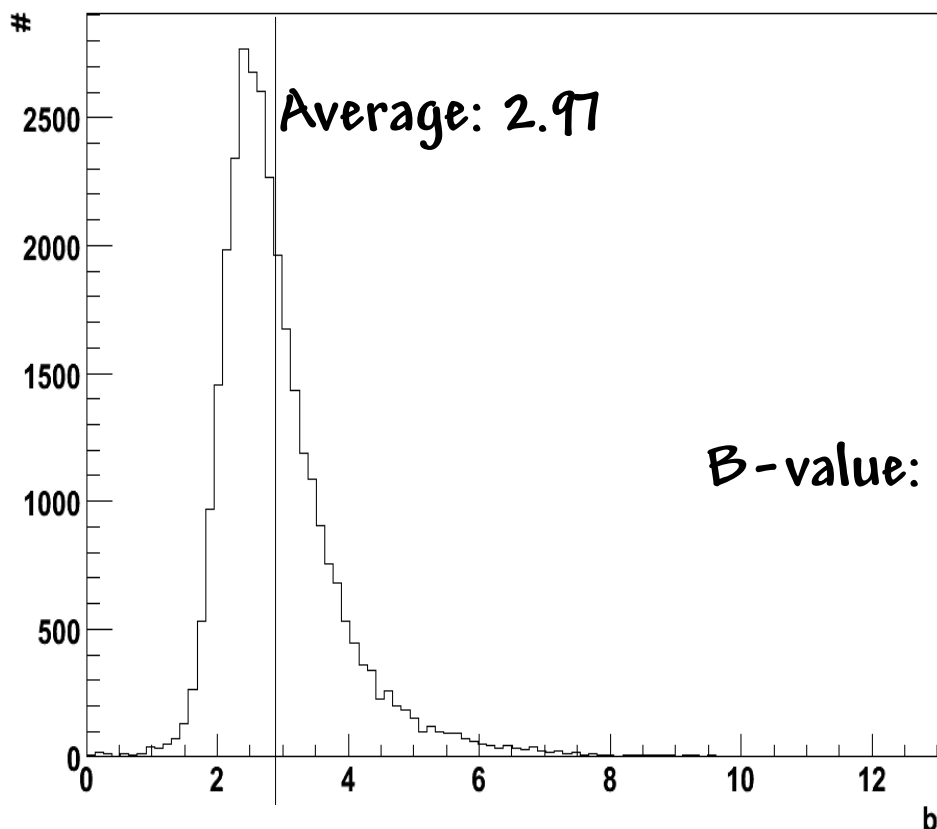
Ghost dressing function



Mapping gauge descriptions [3d, Maas, unpublished]

$b = G(0.280 \text{ GeV}) / G(\infty \text{ GeV})$ for $V = (4.4 \text{ fm})^3$

Ghost dressing function

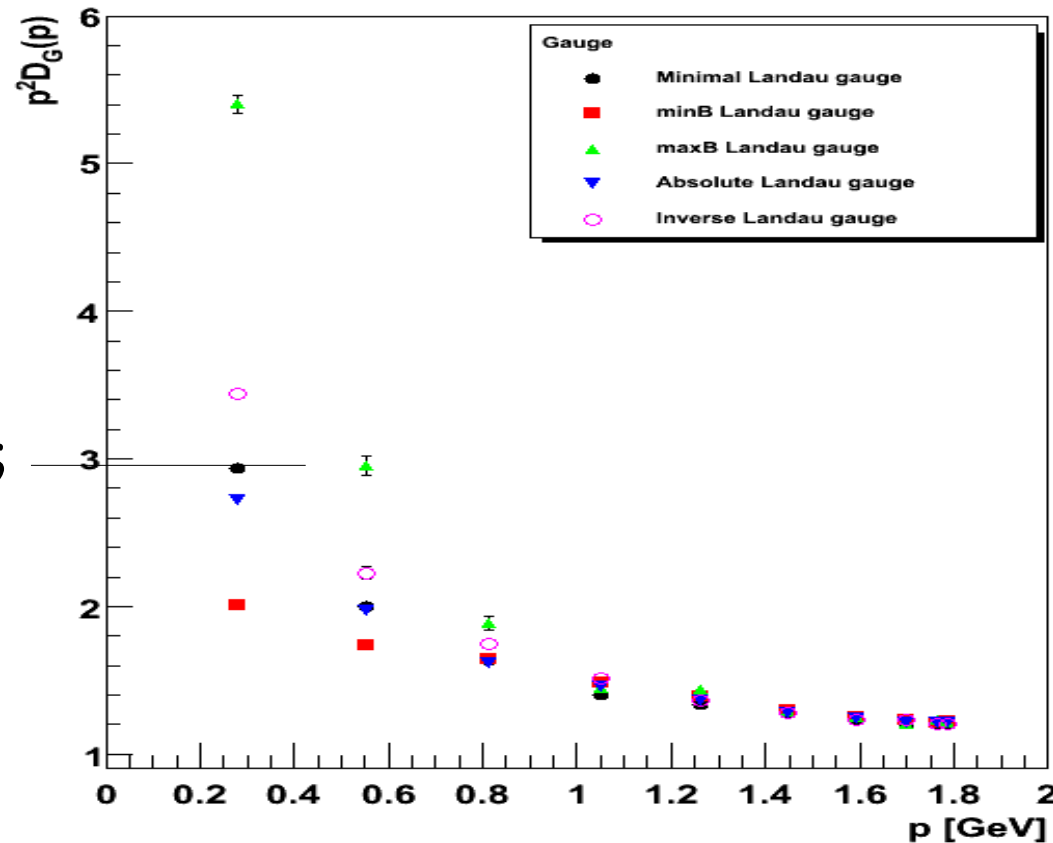
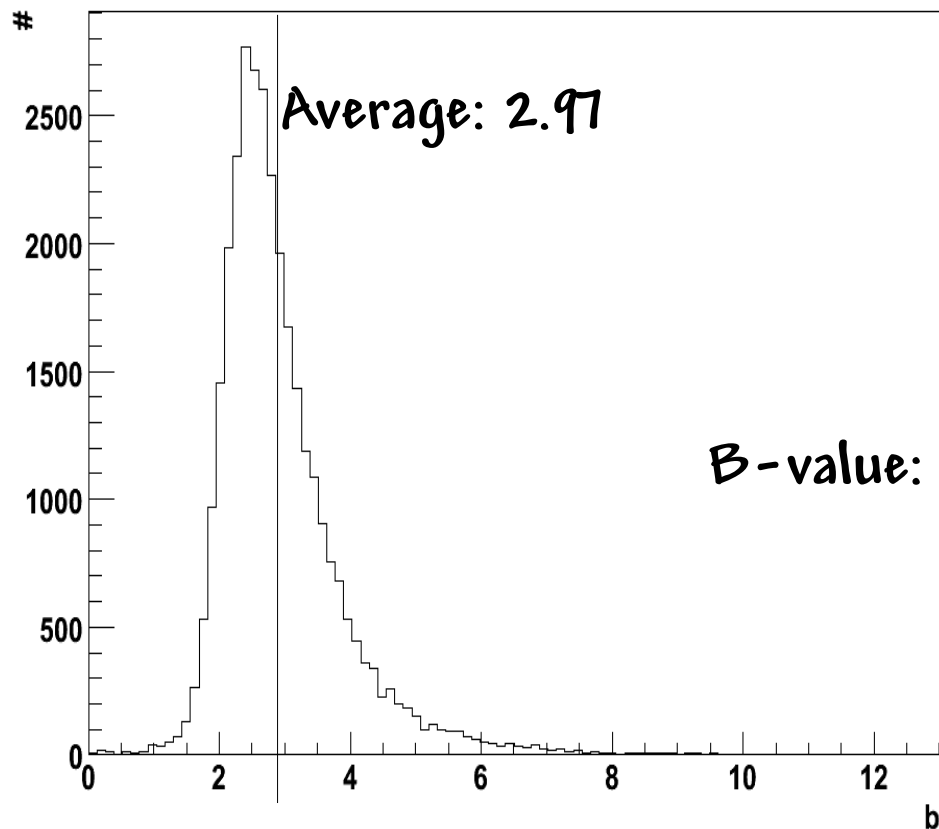


- Minimal Landau gauge is a Landau- B gauge at the average B value
- Probability map is faithful

Mapping gauge descriptions [3d, Maas, unpublished]

$$b = G(0.280 \text{ GeV}) / G(\infty \text{ GeV}) \text{ for } V = (4.4 \text{ fm})^3$$

Ghost dressing function



- Minimal Landau gauge is a Landau-B gauge at the average B value
 - Probability map is faithful
 - Also for absolute or inverse Landau gauge, but with a deformed map
 - All Landau gauges should be map*B-gauge

Landau- B gauges

- **Gluon propagator** is only mildly affected
- **Ghost propagator** differs strongly

Landau- \mathcal{B} gauges

- **Gluon propagator** is only mildly affected
- **Ghost propagator** differs strongly
- **Landau- \mathcal{B} gauges** permit unique gauge-fixing prescription
 - Are **correlation functions** in a **\mathcal{B} gauge** identical to e.g. the **minimal Landau gauge** with its effective **b** value?
 - If it is a faithful map of the probability distribution
 - Similar for other Landau gauges

Landau- \mathcal{B} gauges

- **Gluon propagator** is only mildly affected
- **Ghost propagator** differs strongly
- **Landau- \mathcal{B} gauges** permit unique gauge-fixing prescription
 - Are **correlation functions** in a **\mathcal{B} gauge** identical to e.g. the **minimal Landau gauge** with its effective **b** value?
 - If it is a faithful map of the probability distribution
 - Similar for other Landau gauges
 - Independent of the **method**

Landau- \mathcal{B} gauges

- **Gluon propagator** is only mildly affected
- **Ghost propagator** differs strongly
- **Landau- \mathcal{B} gauges** permit unique gauge-fixing prescription
 - Are **correlation functions** in a **\mathcal{B} gauge** identical to e.g. the **minimal Landau gauge** with its effective **b** value?
 - If it is a faithful map of the probability distribution
 - Similar for other Landau gauges
 - Independent of the **method**
 - But how to identify the **first Gribov horizon** outside the **lattice**?
 - Positivity and monotonicity of the **ghost propagator**? Assume

[Reinhardt et al. 2008]

Methods

- **Lattice**

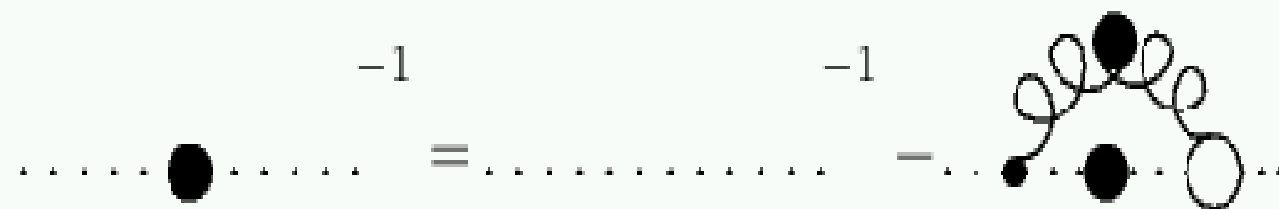
- Discretize space-time in a box and calculate the path-integral and expectation values explicitly
- Full non-perturbative dynamics correctly implemented
- Finite volume artifacts, numerical problems most severe obstacles

Methods

- **Lattice**
 - Discretize space-time in a box and calculate the path-integral and expectation values explicitly
 - Full non-perturbative dynamics correctly implemented
 - Finite volume artifacts, numerical problems most severe obstacles
- **Functional methods (DSE, RGE...)**

(Truncated) Dyson-Schwinger Equations (DSEs)

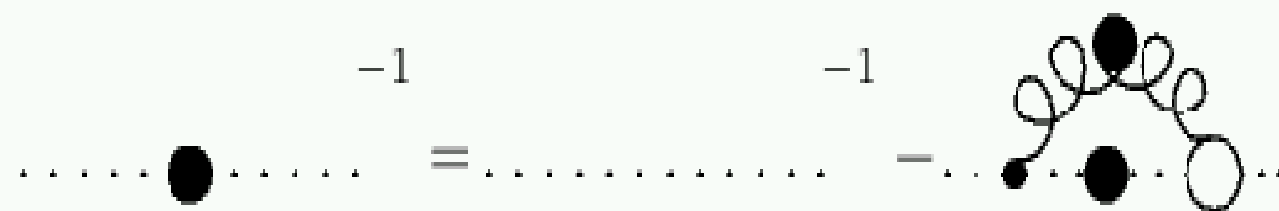
$$\frac{1}{\langle \bar{c} c \rangle(p)} = p^2 + \int dq p_\mu \langle \bar{c} c \rangle(q) \langle A_\mu A_\nu \rangle(p-q) \langle A_\nu \bar{c} c \rangle(p, q)$$



- Infinite set of coupled non-linear integral equations

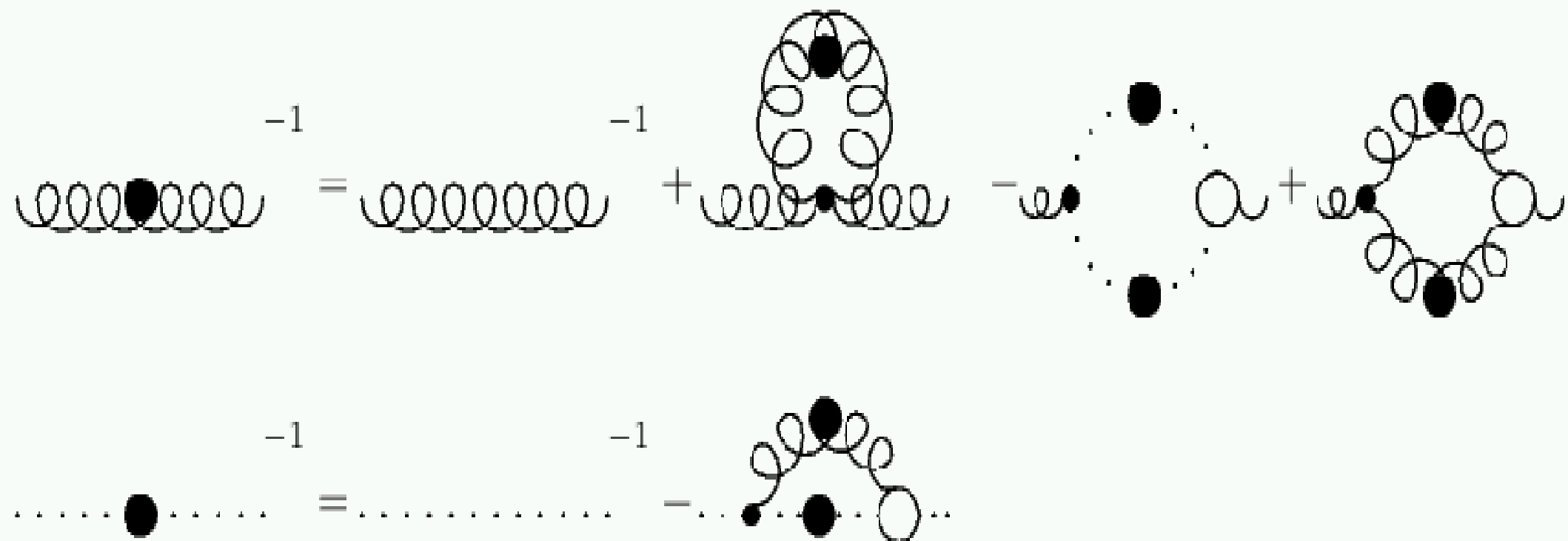
(Truncated) Dyson-Schwinger Equations (DSEs)

$$\frac{1}{\langle \bar{c} c \rangle(p)} = p^2 + \int dq p_\mu \langle \bar{c} c \rangle(q) \langle A_\mu A_\nu \rangle(p-q) \langle A_\nu \bar{c} c \rangle(p, q)$$



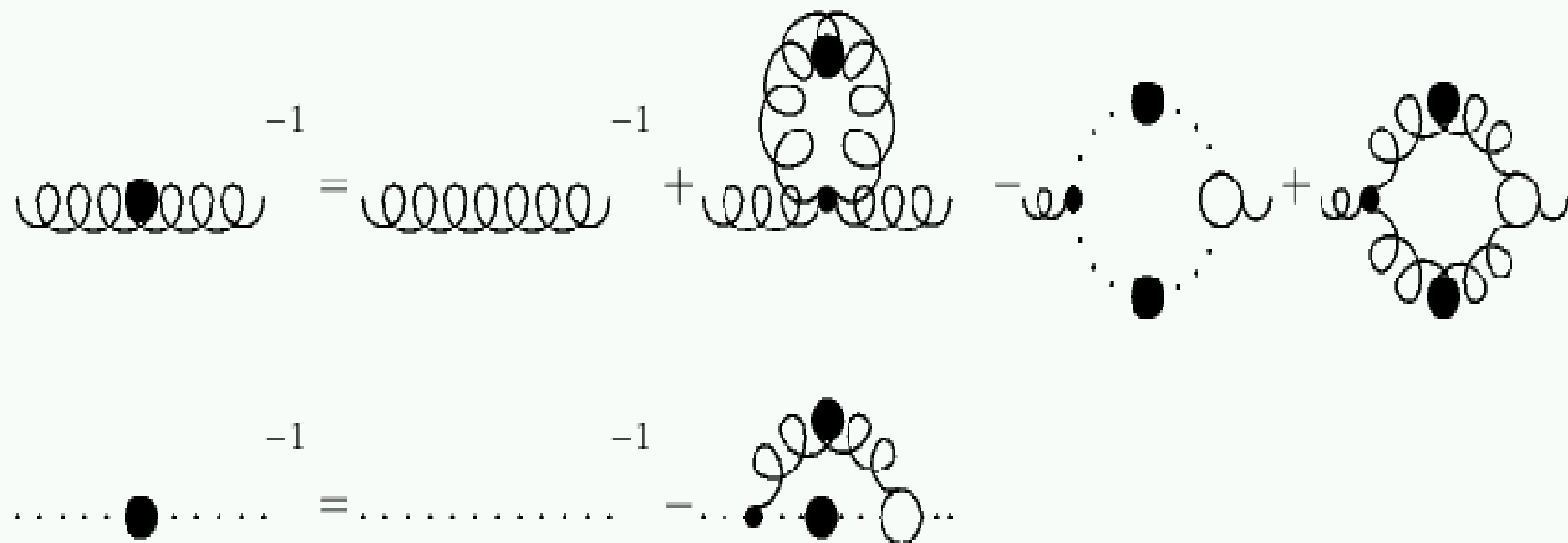
- Infinite set of coupled non-linear integral equations
- Generate also **perturbation theory**

(Truncated) Dyson-Schwinger Equations (DSEs)



- Infinite set of coupled non-linear integral equations
- Generate also **perturbation theory**
- Similar: **Renormalization group equations (RGEs)**

(Truncated) Dyson-Schwinger Equations (DSEs)



- Infinite set of coupled non-linear integral equations
- Generate also **perturbation theory**
- Similar: **Renormalization group equations (RGEs)**
- **Truncation necessary:** Ansätze for the vertices

Practical problems [Fischer et al., 2008]

- Solving the equations at all momenta requires input from the gluon equation etc

Practical problems [Fischer et al., 2008]

- Solving the equations at all momenta requires input from the gluon equation etc.
- Requires **truncations**, in particular dropping of equations for **many-legged correlation functions** [Fischer et al., 2008]

Practical problems [Fischer et al., 2008]

- Solving the equations at all momenta requires input from the gluon equation etc.
- Requires **truncations**, in particular dropping of equations for **many-legged correlation functions** [Fischer et al., 2008]
 - Should preserve renormalization and **perturbation theory**

Practical problems [Fischer et al., 2008]

- Solving the equations at all momenta requires input from the gluon equation etc.
- Requires **truncations**, in particular dropping of equations for **many-legged correlation functions** [Fischer et al., 2008]
 - Should preserve renormalization and **perturbation theory**
 - Introduces artifacts due to the violation of **Slavnov-Taylor identities**
 - Cannot be avoided, at best these can be closed self-consistently

Practical problems [Fischer et al., 2008]

- Solving the equations at all momenta requires input from the gluon equation etc.
- Requires **truncations**, in particular dropping of equations for **many-legged correlation functions** [Fischer et al., 2008]
 - Should preserve renormalization and **perturbation theory**
 - Introduces artifacts due to the violation of **Slavnov-Taylor identities**
 - Cannot be avoided, at best these can be closed self-consistently
 - **Landau gauge** is very advantageous, as such violations do not couple back
 - Only transverse contributions are coupled
 - **Slavnov-Taylor identities** only constrain longitudinal components

Practical problems [Fischer et al., 2008]

- Solving the equations at all momenta requires input from the gluon equation etc.
- Requires **truncations**, in particular dropping of equations for **many-legged correlation functions** [Fischer et al., 2008]
 - Should preserve renormalization and **perturbation theory**
 - Introduces artifacts due to the violation of **Slavnov-Taylor identities**
 - Cannot be avoided, at best these can be closed self-consistently
 - **Landau gauge** is very advantageous, as such violations do not couple back
 - Only transverse contributions are coupled
 - **Slavnov-Taylor identities** only constrain longitudinal components
 - Still other artifacts remain

Practical problems [Fischer et al., 2008]

- Solving the equations at all momenta requires input from the gluon equation etc.
- Requires **truncations**, in particular dropping of equations for **many-legged correlation functions** [Fischer et al., 2008]
 - Should preserve renormalization and **perturbation theory**
 - Introduces artifacts due to the violation of **Slavnov-Taylor identities**
 - Cannot be avoided, at best these can be closed self-consistently
 - **Landau gauge** is very advantageous, as such violations do not couple back
 - Only transverse contributions are coupled
 - **Slavnov-Taylor identities** only constrain longitudinal components
 - Still other artifacts remain
- Consistent and self-consistent truncation schemes can be developed

Methods

- **Lattice**
 - Discretize space-time in a box and calculate the path-integral and expectation values explicitly
 - Full non-perturbative dynamics correctly implemented
 - Finite volume artifacts, numerical problems most severe obstacles
- **Functional methods (DSE, RGE...)**
 - Coupled non-linear integral equations must be solved
 - Requires (often completely uncontrolled) approximations
 - Continuum, partly analytical in the far infrared

Methods

- **Lattice**
 - Discretize space-time in a box and calculate the path-integral and expectation values explicitly
 - Full non-perturbative dynamics correctly implemented
 - Finite volume artifacts, numerical problems most severe obstacles
- **Functional methods** (DSE, RGE...)
 - Coupled non-linear integral equations must be solved
 - Requires (often completely uncontrolled) approximations
 - Continuum, partly analytical in the far infrared
- **Perturbation theory** provides analytical results at large momenta

Methods

- **Lattice**
 - Discretize space-time in a box and calculate the path-integral and expectation values explicitly
 - Full non-perturbative dynamics correctly implemented
 - Finite volume artifacts, numerical problems most severe obstacles
- **Functional methods** (DSE, RGE...)
 - Coupled non-linear integral equations must be solved
 - Requires (often completely uncontrolled) approximations
 - Continuum, partly analytical in the far infrared
- **Perturbation theory** provides analytical results at large momenta
- Combination of all methods most successful!

Implementing \mathcal{B} Landau gauges in the continuum

- \mathcal{B} -Landau gauges are implemented via the **ghost dressing function equation**

- Symbolically
$$\frac{1}{G(p)} = Z + \int dq K(p, q) D(p-q) G(q)$$

Implementing \mathcal{B} Landau gauges in the continuum

- \mathcal{B} -Landau gauges are implemented via the **ghost dressing function equation**

- Symbolically
$$\frac{1}{G(p)} = Z + \int dq K(p, q) D(p-q) G(q)$$

- Kernel depends on

Implementing \mathcal{B} Landau gauges in the continuum

- \mathcal{B} -Landau gauges are implemented via the **ghost dressing function equation**

- Symbolically
$$\frac{1}{G(p)} = Z + \int dq K(p, q) D(p-q) G(q)$$

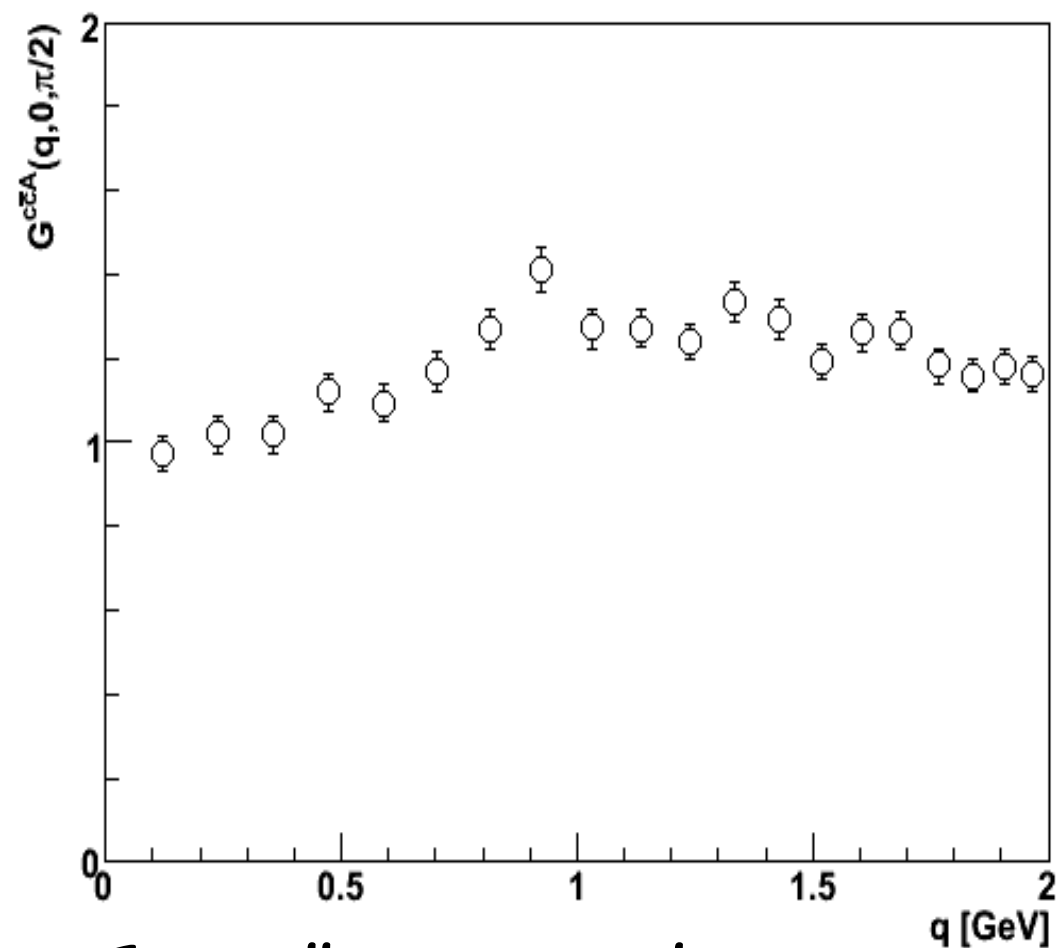
- Kernel depends on

- **Ghost-gluon vertex**: Can be taken to be almost bare

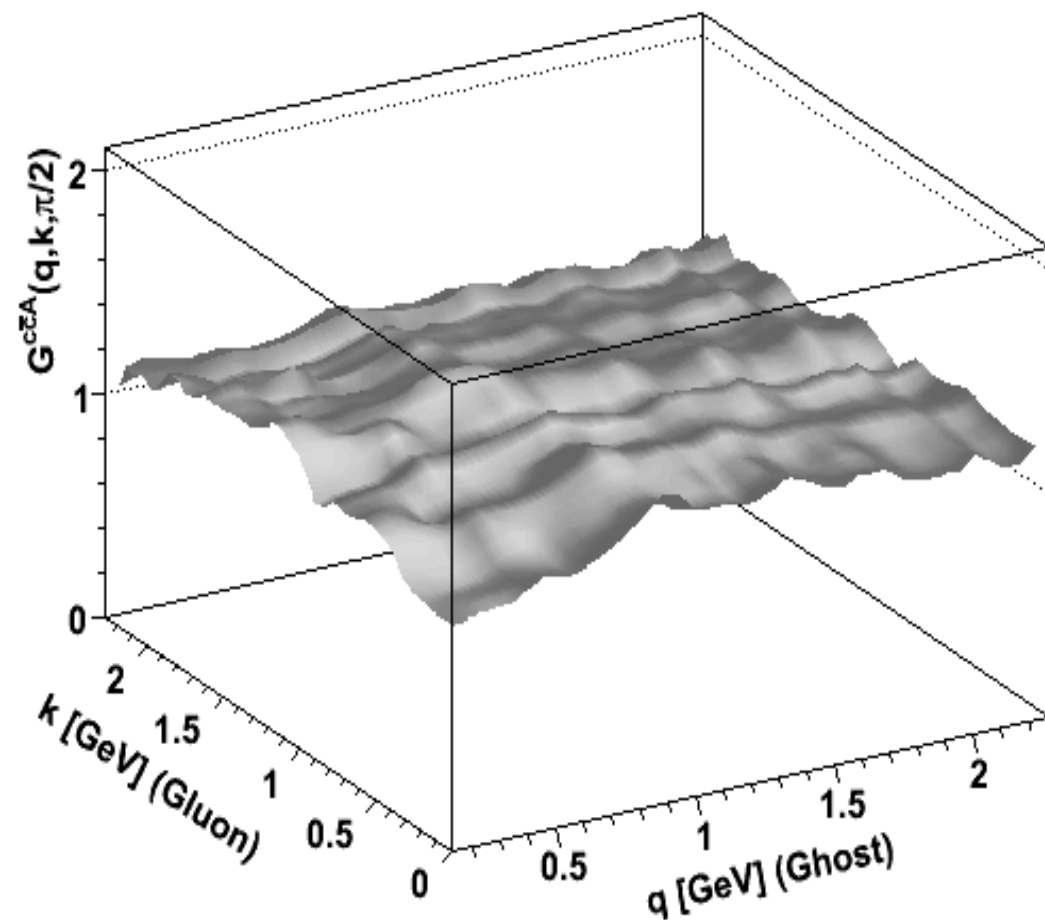
Ghost-gluon vertex in 3d

[60³, beta=4.2, Cucchieri et al., 2008]

Ghost-gluon vertex, one momentum vanishing



Ghost-gluon vertex, orthogonal momenta



- Essentially constant, only some structure at 1 GeV
- Same in 2d and 4d
- In agreement with DSE predictions [Schleifenbaum et al., PRD 2005]

Implementing \mathcal{B} Landau gauges in the continuum

- \mathcal{B} -Landau gauges are implemented via the **ghost dressing function equation**

- Symbolically
$$\frac{1}{G(p)} = Z + \int dq K(p, q) D(p-q) G(q)$$

- Kernel depends on

- **Ghost-gluon vertex**: Can be taken to be almost bare
- **Gluon propagator**: Massive or suppressed

Implementing \mathcal{B} Landau gauges in the continuum

- \mathcal{B} -Landau gauges are implemented via the **ghost dressing function equation**

- Symbolically
$$\frac{1}{G(p)} = Z + \int dq K(p, q) D(p-q) G(q)$$

- Kernel depends on
 - **Ghost-gluon vertex**: Can be taken to be almost bare
 - **Gluon propagator**: Massive or suppressed
- Then the self-energy diagram behaves as $A + \mathcal{L}(p)$ with $\mathcal{L}(0) = 0$

Implementing \mathcal{B} Landau gauges in the continuum

- \mathcal{B} -Landau gauges are implemented via the **ghost dressing function equation**

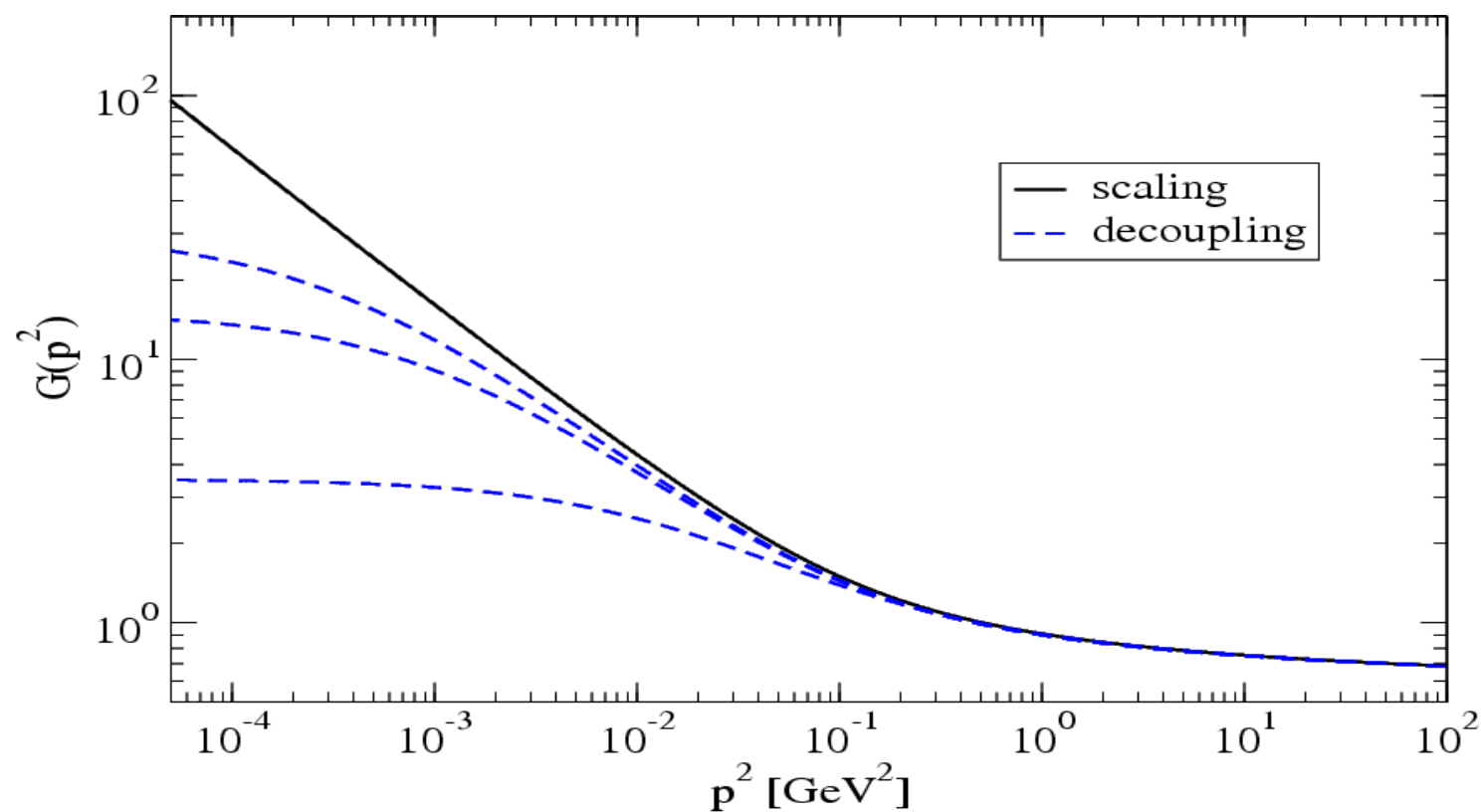
- Symbolically
$$\frac{1}{G(p)} = Z + \int dq K(p, q) D(p-q) G(q)$$

- Kernel depends on
 - **Ghost-gluon vertex**: Can be taken to be almost bare
 - **Gluon propagator**: Massive or suppressed
- Then the self-energy diagram behaves as $A + \mathcal{L}(p)$ with $\mathcal{L}(0) = 0$
- Select wave-function renormalization Z such as to implement a particular \mathcal{B} -value $\mathcal{B} = 1/(Z+A)$ and thus **\mathcal{B} -Landau gauge**

Ghost dressing function in the continuum [Fischer et al., 2008]

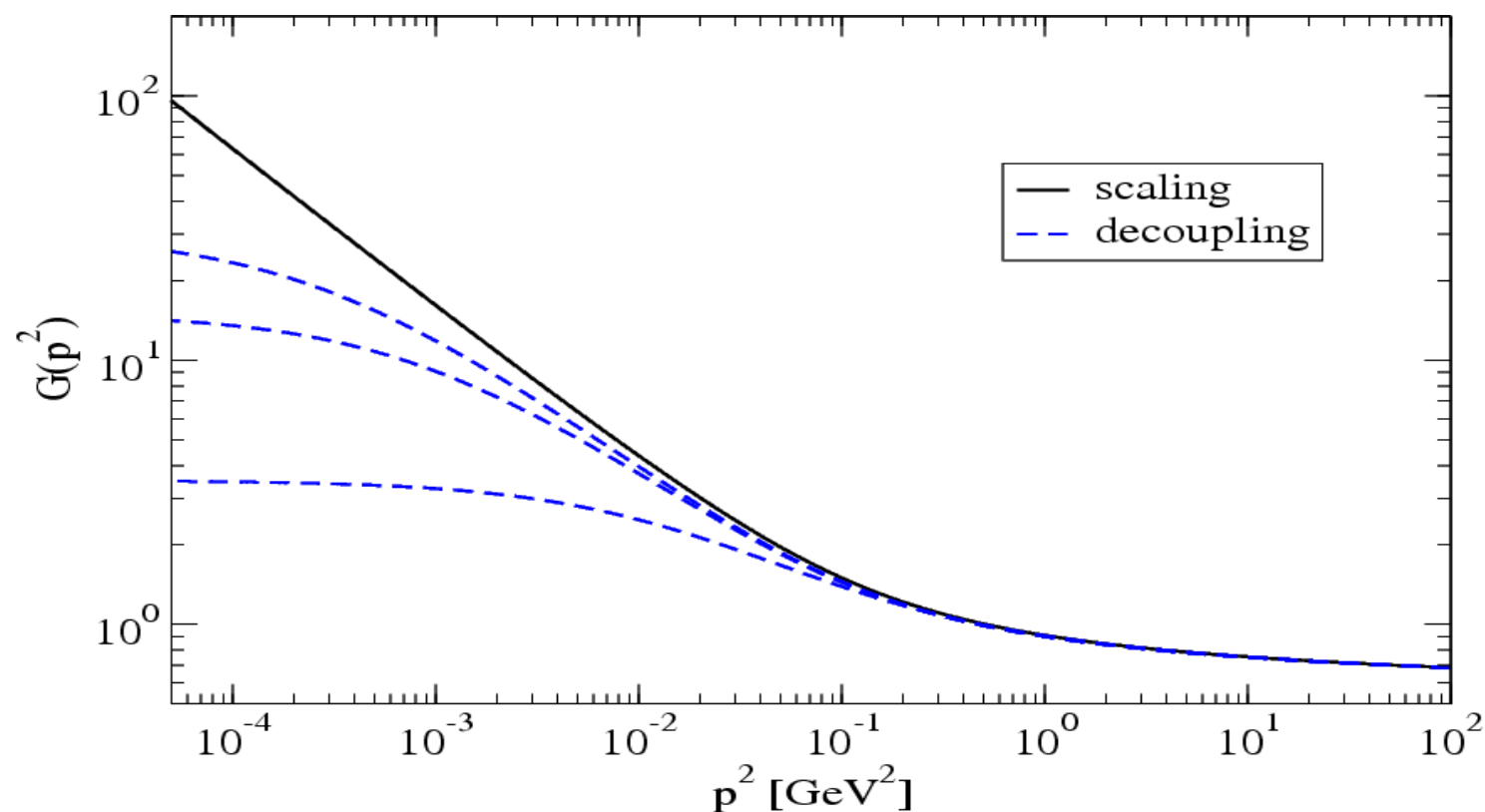
- Continuous 1-parameter family of solutions

Ghost dressing function in the continuum [Fischer et al., 2008]



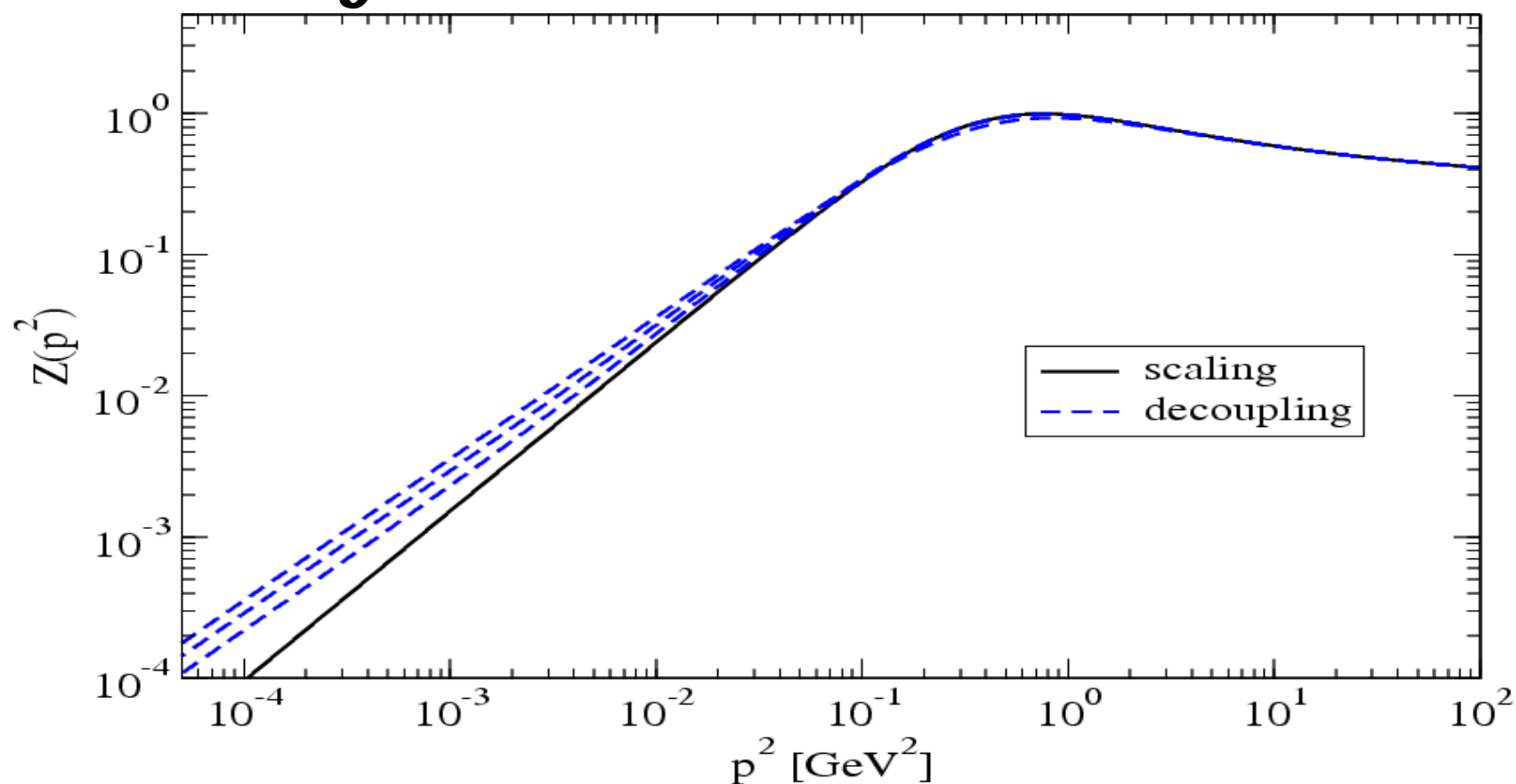
- Continuous 1-parameter family of solutions
- Starting at finite B up to positive infinite B (divergent G)
 - Classify as decoupling and scaling type-solution

Ghost dressing function in the continuum [Fischer et al., 2008]



- Continuous 1-parameter family of solutions
- Starting at finite B up to positive infinite B (divergent G)
 - Classify as decoupling and scaling type-solution
 - Continuum implementation of the B -Landau gauges

Gluon dressing function in the continuum [Fischer et al., 2008]



- **Decoupling gauges** yield a decoupling (infrared massive) **gluon propagator** - consistent with the lattice result
- The **scaling gauge** yields an infrared vanishing **gluon propagator**

Gauge-dependence of the propagators

- The Landau-B gauge 1-parameter family can be constructed with lattice, DSE and FRG methods

Gauge-dependence of the propagators

- The Landau- β gauge 1-parameter family can be constructed with lattice, DSE and FRG methods
- The endpoint appears to be a scaling solution, while all other Landau- β gauges show a decoupling-type behavior

Gauge-dependence of the propagators

- The **Landau-B gauge** 1-parameter family can be constructed with **lattice, DSE and FRG methods**
- The endpoint appears to be a scaling solution, while all other **Landau-B gauges** show a decoupling-type behavior
- The qualitative infrared behavior is thus a matter of the selected non-perturbative gauge!
 - Possible physical information is mixed with pure gauge artifacts
 - It is possible to select thus the infrared behavior fitting the needs of a particular problem, similar to perturbation theory

Gauge-dependence of the propagators

- The **Landau-B gauge** 1-parameter family can be constructed with **lattice, DSE and FRG methods**
- The endpoint appears to be a scaling solution, while all other **Landau-B gauges** show a decoupling-type behavior
- The qualitative infrared behavior is thus a matter of the selected non-perturbative gauge!
 - Possible physical information is mixed with pure gauge artifacts
 - It is possible to select thus the infrared behavior fitting the needs of a particular problem, similar to perturbation theory
 - This needs a more formal formulation, volume studies...etc.

Example: Scaling gauge and the physical Hilbert space

- Use the **scaling gauge** to determine the structure of the physical Hilbert space

Example: Scaling gauge and the physical Hilbert space

- Use the **scaling gauge** to determine the structure of the physical Hilbert space
- Assume: No explicit transition functions

Example: Scaling gauge and the physical Hilbert space

- Use the **scaling gauge** to determine the structure of the physical Hilbert space
- Assume: No explicit transition functions
 - Manifold of a gauge theory is non-trivial, no unique coordinate system
 - Requires transition function from one coordinate system to another
 - Assume that this can be included into the definition of the field variables
 - Then equations are form equivalent to the standard form
 - All non-triviality contained in the **correlation functions** (**B value** etc.)
 - Globally well-defined in the sense of differential geometry

Example: Scaling gauge and the physical Hilbert space

- Use the **scaling gauge** to determine the structure of the physical Hilbert space
- Assume: No explicit transition functions
 - Manifold of a gauge theory is non-trivial, no unique coordinate system
 - Requires transition function from one coordinate system to another
 - Assume that this can be included into the definition of the field variables
 - Then equations are form equivalent to the standard form
 - All non-triviality contained in the **correlation functions** (**B value** etc.)
 - Globally well-defined in the sense of differential geometry
 - If not correct: All of this is void

Infrared properties of the scaling solution

- **DSEs and RGs** deliver a qualitative asymptotic infrared solution for all **Green's functions** and all dimensions [Alkofer et al., PLB 2005, Fischer et al. PRD 2007]

K

Infrared properties of the scaling solution

- **DSEs and RGs** deliver a qualitative asymptotic infrared solution for all **Green's functions** and all dimensions [Alkofer et al., PLB 2005, Fischer et al. PRD 2007]
 - Power-laws. Simplest case: All momenta equal and m gluon and n ghost legs in d dimensions:
 - $p^{(\frac{n_{ghost}}{2} - m_{gluon})\kappa + (1 - \frac{n_{ghost}}{2})(\frac{d}{2} - 2)}$ [Alkofer et al. 2006, Huber et al. 2008]
 - Analytic solution without truncation
 - It is the only power-law-like solution [Fischer et al. PRD 2007]

K

Infrared properties of the scaling solution

- **DSEs and RGs** deliver a qualitative asymptotic infrared solution for all **Green's functions** and all dimensions [Alkofer et al., PLB 2005, Fischer et al. PRD 2007]
 - Power-laws. Simplest case: All momenta equal and m gluon and n ghost legs in d dimensions:
 - $p^{(\frac{n_{ghost}}{2} - m_{gluon})\kappa + (1 - \frac{n_{ghost}}{2})(\frac{d}{2} - 2)}$ [Alkofer et al. 2006, Huber et al. 2008]
 - Analytic solution without truncation
 - It is the only power-law-like solution [Fischer et al. PRD 2007]
- **Truncated DSEs** give quantitative predictions for the exponent κ [Zwanziger, PRD 2002, Lerche et al. PRD 2002, Maas et al., EPJC 2004]

Infrared properties of the scaling solution

- **DSEs and RGs** deliver a qualitative asymptotic infrared solution for all **Green's functions** and all dimensions [Alkofer et al., PLB 2005, Fischer et al. PRD 2007]
 - Power-laws. Simplest case: All momenta equal and m gluon and n ghost legs in d dimensions:
 - $p^{(\frac{n_{ghost}}{2} - m_{gluon})\kappa + (1 - \frac{n_{ghost}}{2})(\frac{d}{2} - 2)}$ [Alkofer et al. 2006, Huber et al. 2008]
 - Analytic solution without truncation
 - It is the only power-law-like solution [Fischer et al. PRD 2007]
- **Truncated DSEs** give quantitative predictions for the exponent κ [Zwanziger, PRD 2002, Lerche et al. PRD 2002, Maas et al., EPJC 2004]
- In particular: **IR-vanishing gluon propagator, IR-enhanced ghost propagator**

Infrared properties of the scaling solution

- **DSEs and RGs** deliver a qualitative asymptotic infrared solution for all **Green's functions** and all dimensions [Alkofer et al., PLB 2005, Fischer et al. PRD 2007]
 - Power-laws. Simplest case: All momenta equal and m gluon and n ghost legs in d dimensions:
 - $p^{(\frac{n_{ghost}}{2} - m_{gluon})\kappa + (1 - \frac{n_{ghost}}{2})(\frac{d}{2} - 2)}$ [Alkofer et al. 2006, Huber et al. 2008]
 - Analytic solution without truncation
 - It is the only power-law-like solution [Fischer et al. PRD 2007]
- **Truncated DSEs** give quantitative predictions for the exponent κ [Zwanziger, PRD 2002, Lerche et al. PRD 2002, Maas et al., EPJC 2004]
- In particular: **IR-vanishing gluon propagator, IR-enhanced ghost propagator**
- Can be expanded to the case with matter fields [Alkofer et al., 2007/8]

Absence of the gluon from the physical spectrum [Zwanziger, 1994-2009]

- **Gluon propagator** violates positivity

Absence of the gluon from the physical spectrum [Zwanziger, 1994-2009]

- **Gluon propagator** violates positivity
 - **Propagator** determines spectral function

$$\text{Propagator} = \text{One particle part} + \int dq^2 \frac{\text{spectral function}(q^2)}{p^2 + q^2}$$

- One particle part by definition positive

Absence of the gluon from the physical spectrum [Zwanziger, 1994-2009]

- **Gluon propagator** violates positivity

- **Propagator** determines spectral function

$$\text{Propagator} = \text{One particle part} + \int dq^2 \frac{\text{spectral function}(q^2)}{p^2 + q^2}$$

- One particle part by definition positive
- **Gluon propagator** vanishes at zero momentum
- Gluon not part of the physical sub-space!

Absence of the gluon from the physical spectrum [Zwanziger, 1994-2009]

- **Gluon propagator** violates positivity

- **Propagator** determines spectral function

$$\text{Propagator} = \text{One particle part} + \int dq^2 \frac{\text{spectral function}(q^2)}{p^2 + q^2}$$

- One particle part by definition positive
- **Gluon propagator** vanishes at zero momentum
- Gluon not part of the physical sub-space!
- In decoupling-type **Landau-B gauges** this is not obvious

Absence of the gluon from the physical spectrum [Zwanziger, 1994-2009]

- **Gluon propagator** violates positivity

- **Propagator** determines spectral function

$$\text{Propagator} = \text{One particle part} + \int dq^2 \frac{\text{spectral function}(q^2)}{p^2 + q^2}$$

- One particle part by definition positive
 - **Gluon propagator** vanishes at zero momentum
 - Gluon not part of the physical sub-space!
 - In decoupling-type **Landau-B gauges** this is not obvious
 - But explicit calculations show this to be the case
- [Fischer et al., 2008, Cucchieri et al. 2004]
- Gluon in general **Landau-B gauges** not part of the spectrum

Absence of the gluon from the physical spectrum [Zwanziger, 1994-2009]

- **Gluon propagator** violates positivity

- **Propagator** determines spectral function

$$\text{Propagator} = \text{One particle part} + \int dq^2 \frac{\text{spectral function}(q^2)}{p^2 + q^2}$$

- One particle part by definition positive
- **Gluon propagator** vanishes at zero momentum
- Gluon not part of the physical sub-space!
- In decoupling-type **Landau-B gauges** this is not obvious
- But explicit calculations show this to be the case
[Fischer et al., 2008, Cucchieri et al. 2004]
- Gluon in general **Landau-B gauges** not part of the spectrum
- Can be extended to Wilson criteria for quarks [Pawlowski et al., 2008]

What is with general colored states?

- **Perturbatively**, physical states are selected by the cohomology of the **BRST charge**

What is with general colored states?

- **Perturbatively**, physical states are selected by the cohomology of the **BRST charge**
- Can **BRST** be constructed non-perturbatively?

What is with general colored states?

- **Perturbatively**, physical states are selected by the cohomology of the **BRST charge**
- Can **BRST** be constructed non-perturbatively?
- Yes: Neuberger-von Smekal construction [Neuberger, 1987, von Smekal et al. 2007/8]
 - Gives an explicit lattice realization of a globally well-defined BRST charge, which differs from the perturbative definitions on the lattice
 - Requires at finite discretization a modified definition of the gluon field
 - In the continuum limit, the standard construction is recovered

What is with general colored states?

- **Perturbatively**, physical states are selected by the cohomology of the **BRST charge**
- Can **BRST** be constructed non-perturbatively?
- Yes: Neuberger-von Smekal construction [Neuberger, 1987, von Smekal et al. 2007/8]
 - Gives an explicit lattice realization of a globally well-defined BRST charge, which differs from the perturbative definitions on the lattice
 - Requires at finite discretization a modified definition of the gluon field
 - In the continuum limit, the standard construction is recovered
- However: Yet unknown, to which **Landau-B gauge** this corresponds, or if it applies generally to all **Landau gauges**
 - Will require always large (not infinitesimal) transformations

What is with general colored states?

- **Perturbatively**, physical states are selected by the cohomology of the **BRST charge**
- Can **BRST** be constructed non-perturbatively?
- Yes: Neuberger-von Smekal construction [Neuberger, 1987, von Smekal et al. 2007/8]
 - Gives an explicit lattice realization of a globally well-defined BRST charge, which differs from the perturbative definitions on the lattice
 - Requires at finite discretization a modified definition of the gluon field
 - In the continuum limit, the standard construction is recovered
- However: Yet unknown, to which **Landau-B gauge** this corresponds, or if it applies generally to all **Landau gauges**
 - Will require always large (not infinitesimal) transformations
- Assume it to be valid in the **scaling gauge**

The Kugo-Ojima construction [Kugo & Ojima 1979]

- If a globally well-defined BRST exists, the Kugo-Ojima construction can be performed

The Kugo-Ojima construction [Kugo & Ojima 1979]

- If a globally well-defined BRST exists, the Kugo-Ojima construction can be performed
- If the **ghost dressing function** is infrared singular, this construction implies the absence of colored states from the spectrum

The Kugo-Ojima construction [Kugo & Ojima 1979]

- If a globally well-defined BRST exists, the Kugo-Ojima construction can be performed
- If the **ghost dressing function** is infrared singular, this construction implies the absence of colored states from the spectrum
- Given the assumptions, this is realized in the **scaling Landau-B gauge**

The Kugo-Ojima construction [Kugo & Ojima 1979]

- If a globally well-defined BRST exists, the Kugo-Ojima construction can be performed
- If the **ghost dressing function** is infrared singular, this construction implies the absence of colored states from the spectrum
- Given the assumptions, this is realized in the **scaling Landau-B gauge**
 - Provides an explicit realization of confinement and of the physical Hilbert space in this gauge

The Kugo-Ojima construction [Kugo & Ojima 1979]

- If a globally well-defined BRST exists, the Kugo-Ojima construction can be performed
- If the **ghost dressing function** is infrared singular, this construction implies the absence of colored states from the spectrum
- Given the assumptions, this is realized in the **scaling Landau-B gauge**
 - Provides an explicit realization of confinement and of the physical Hilbert space in this gauge
 - Very attractive, if correct

Summary

- *Gauge-fixing* in the non-perturbative domain requires additional non-local conditions

Summary

- *Gauge-fixing* in the non-perturbative domain requires additional non-local conditions
- Perturbative gauges become a family of non-perturbative gauges

Summary

- **Gauge-fixing** in the non-perturbative domain requires additional non-local conditions
- Perturbative gauges become a family of non-perturbative gauges
- An explicit construction in the **Landau gauge** case is the one-parameter family of **Landau- β** gauges

Summary

- **Gauge-fixing** in the non-perturbative domain requires additional non-local conditions
- Perturbative gauges become a family of non-perturbative gauges
- An explicit construction in the **Landau gauge** case is the one-parameter family of **Landau-B** gauges
 - Can be formulated using only **correlation functions**, independent of the **method**

Summary

- **Gauge-fixing** in the non-perturbative domain requires additional non-local conditions
- Perturbative gauges become a family of non-perturbative gauges
- An explicit construction in the **Landau gauge** case is the one-parameter family of **Landau- β** gauges
 - Can be formulated using only **correlation functions**, independent of the **method**
- Additional freedom: Select a **gauge** which suits the needs

Summary

- **Gauge-fixing** in the non-perturbative domain requires additional non-local conditions
- Perturbative gauges become a family of non-perturbative gauges
- An explicit construction in the **Landau gauge** case is the one-parameter family of **Landau- \mathcal{B}** gauges
 - Can be formulated using only **correlation functions**, **independent of the method**
- Additional freedom: Select a **gauge** which suits the needs
 - Example: **Scaling Landau- \mathcal{B}** gauge to determine the physical Hilbert space

Summary

- **Gauge-fixing** in the non-perturbative domain requires additional non-local conditions
- Perturbative gauges become a family of non-perturbative gauges
- An explicit construction in the **Landau gauge** case is the one-parameter family of **Landau- \mathcal{B}** gauges
 - Can be formulated using only **correlation functions**, independent of the **method**
- Additional freedom: Select a **gauge** which suits the needs
 - Example: **Scaling Landau- \mathcal{B}** gauge to determine the physical Hilbert space
- But a formal formulation and more details are still required...