

Describing Gluons

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Overview

- *Gauge freedom* in Yang-Mills theory

Supported by the FWF

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- Gauge freedom in Yang-Mills theory
- Non-perturbative gauges

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- Comparing lattice and continuum

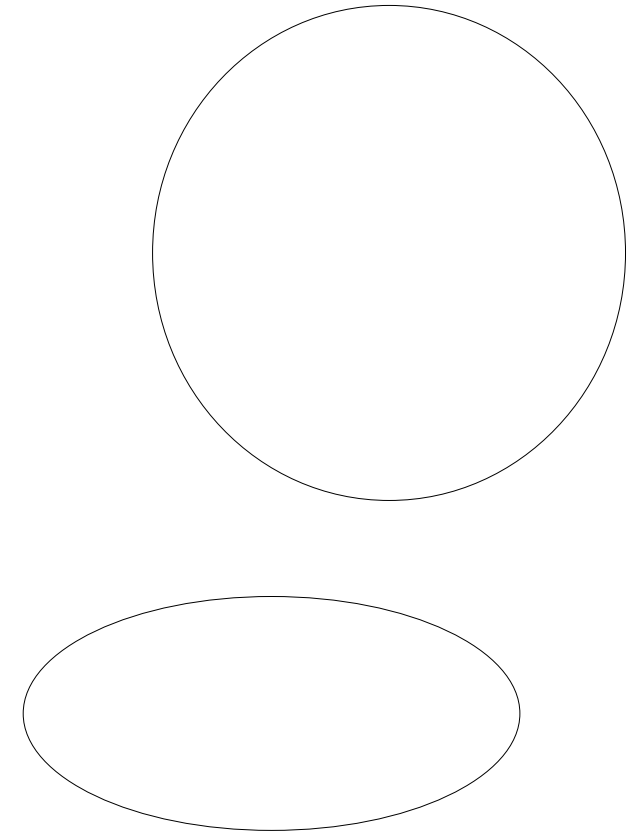
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Overview

- Gauge freedom in Yang-Mills theory
- Non-perturbative gauges
- Comparing lattice and continuum
- Correlation functions and physics

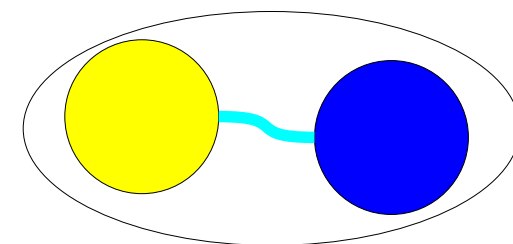
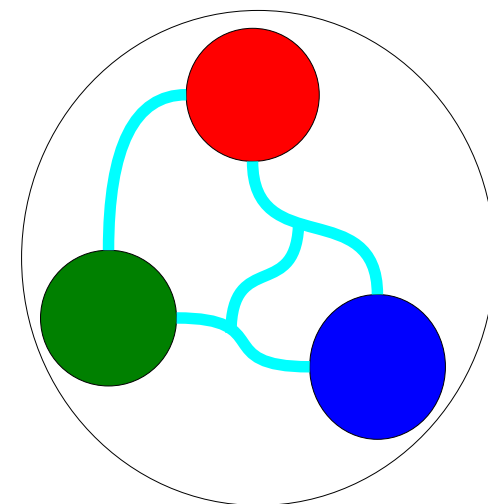
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Strong interactions



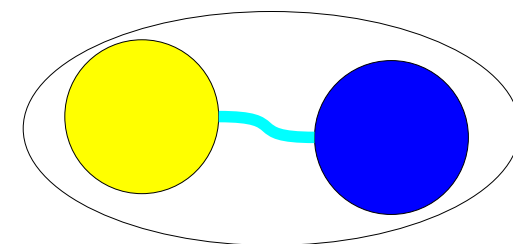
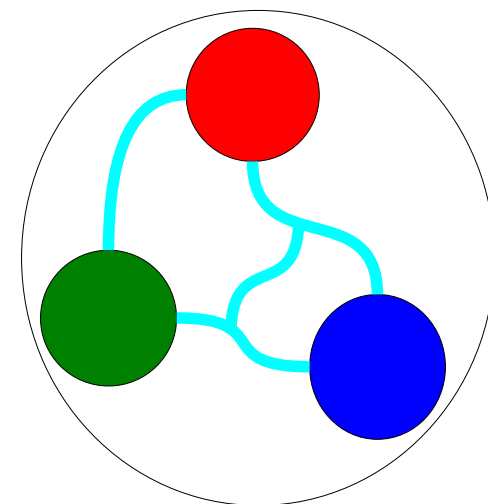
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- Substructure can be described by QCD
 - Degrees of freedom are quarks and gluons



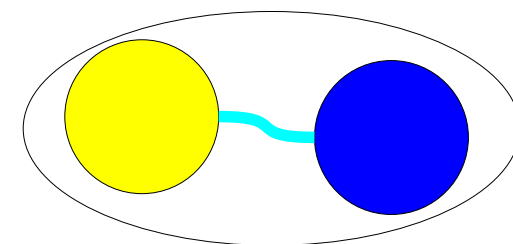
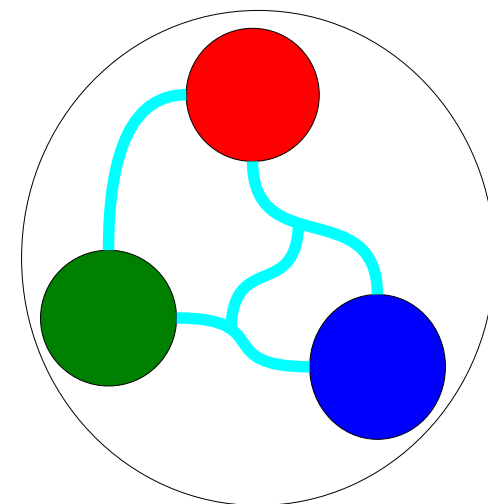
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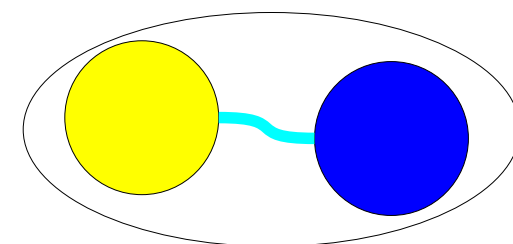
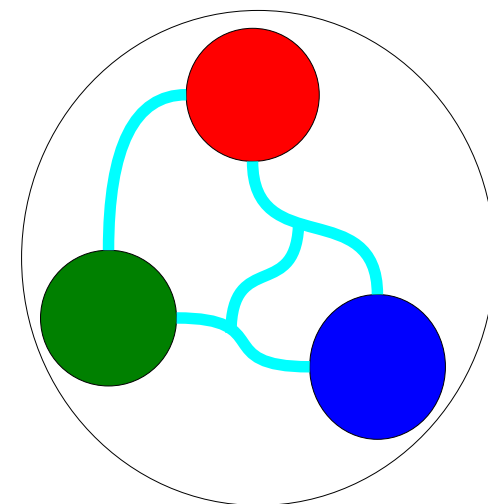
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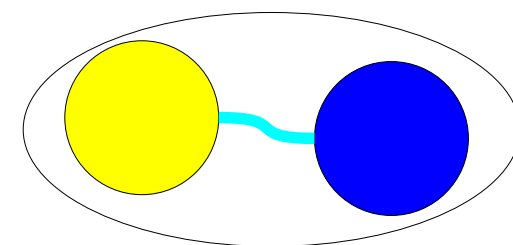
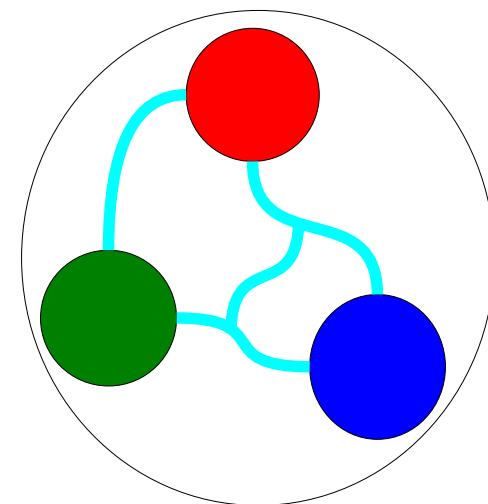
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 - Different from QED
 - No concept of a local gauge-invariant color charge distribution
 - Similar to energy density in general relativity



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 - Different from QED
 - No concept of a local gauge-invariant color charge distribution
 - Similar to energy density in general relativity
- Reduce complexity and ignore Quarks: Yang-Mills theory
 - Affects gluon properties only quantitatively



Yang-Mills Theory

- Lagrangian (after Wick rotation to Euclidean space-time):

$$L = -\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - gf^{abc} A_\mu^b A_\nu^c$$

- Degrees of freedom:

Gluons: A_μ^a

- g is the coupling constant, giving the strength of coupling
- f^{abc} are numbers, depending on the **gauge group**, $SU(3)$ for QCD:
quarks & gluon are organized in multiplets, just as with spin

Gauge-fixing

- Yang-Mills theory is a **gauge theory**

- **Gauge transformations** $A_\mu^a \rightarrow A_\mu^a + (\delta^{ab} \partial_\mu - g f^{abc} A_\mu^c) \phi^b(x)$

with arbitrary $\phi^a(x)$ change the gauge fields, but leave physics invariant

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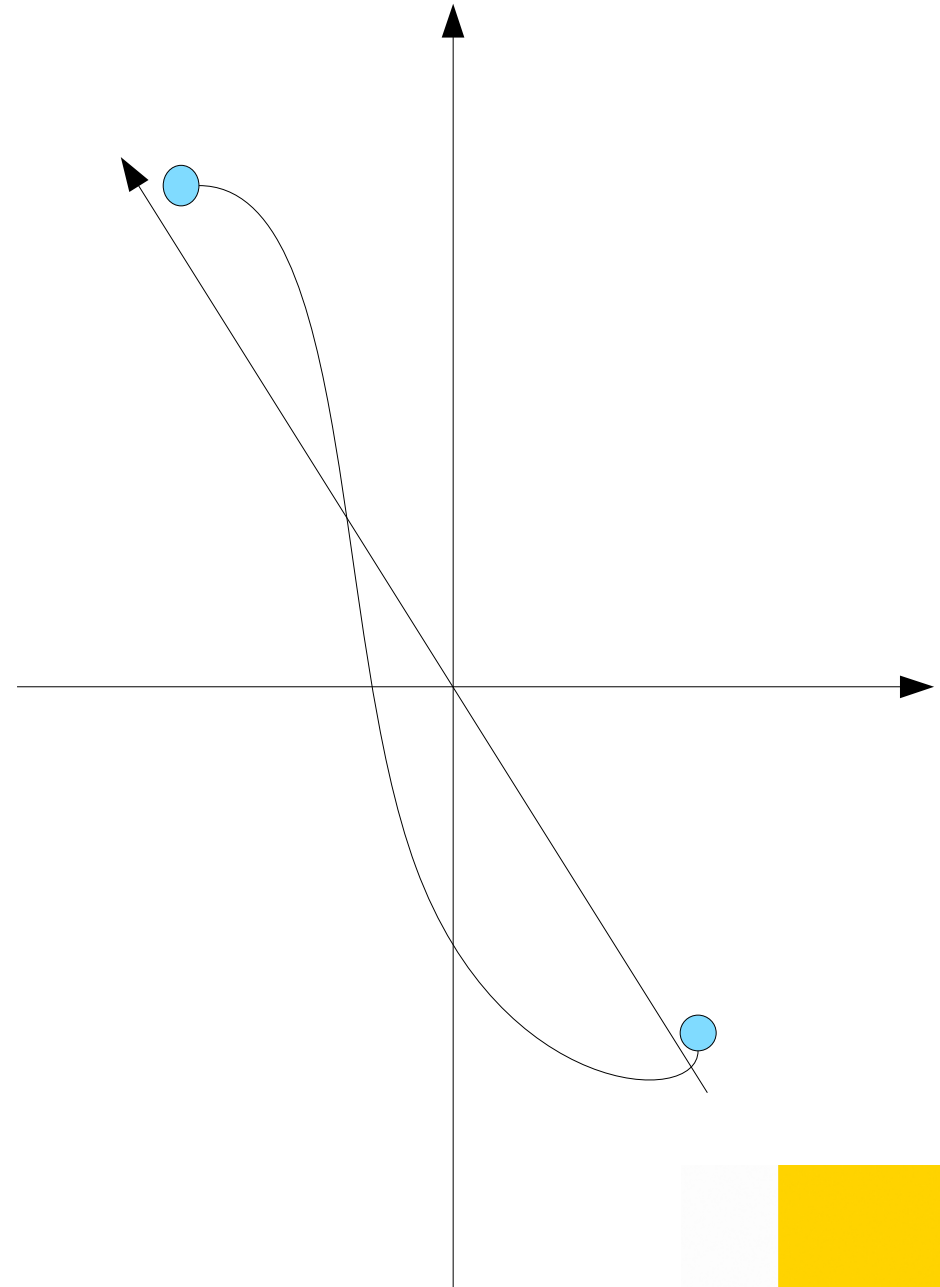
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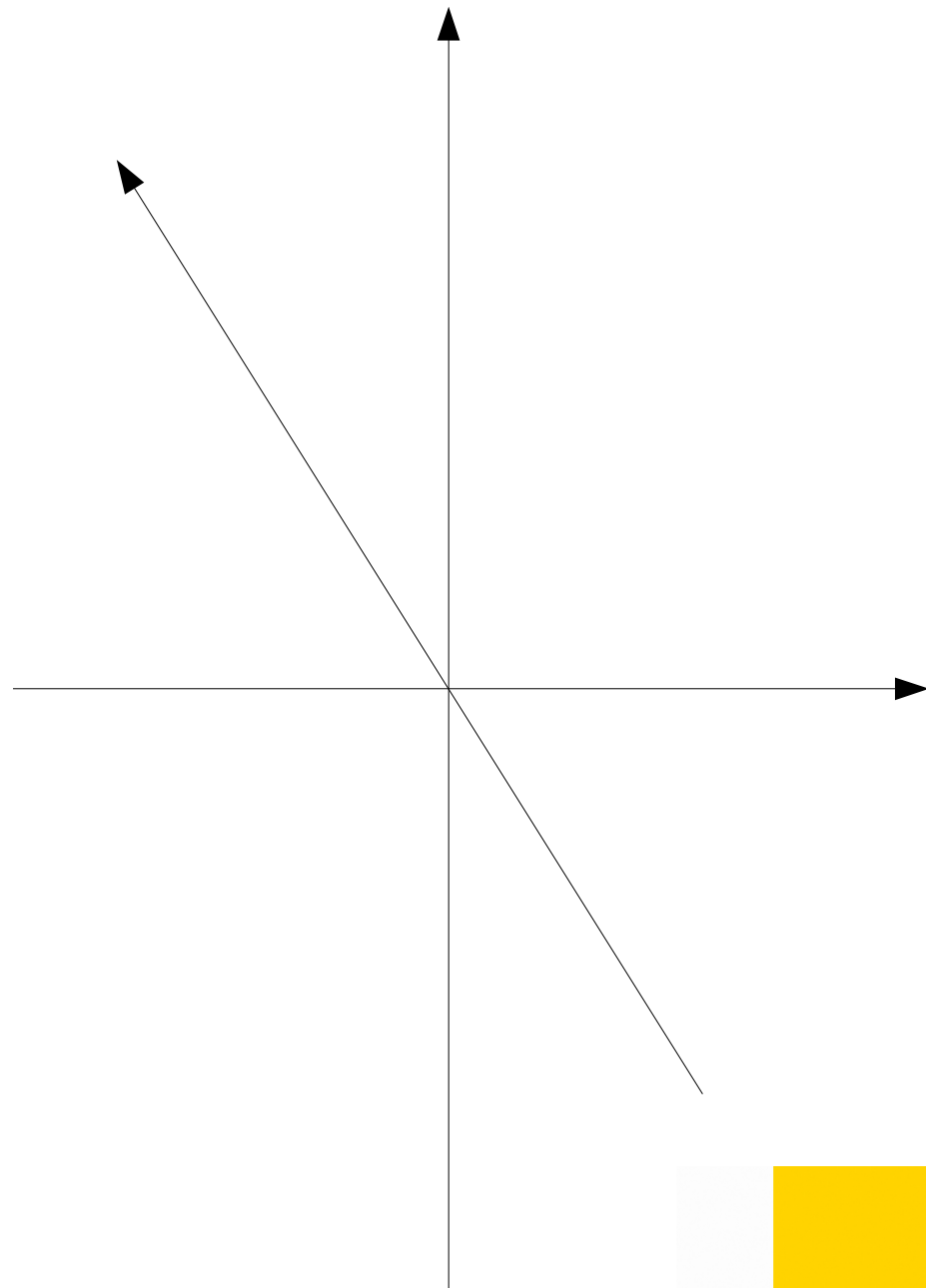
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- **Correlation functions** are in general **gauge-dependent**
 - **Gauge-fixing is required**
- **Example: Landau gauge condition** $\partial_\mu A_\mu^a = 0$
 - Here only Landau gauge results will be considered
 - Many other gauges have been studied

Configuration space (artist's view)

- Gauge fields not unique
 - Gauge transformation does not change physics

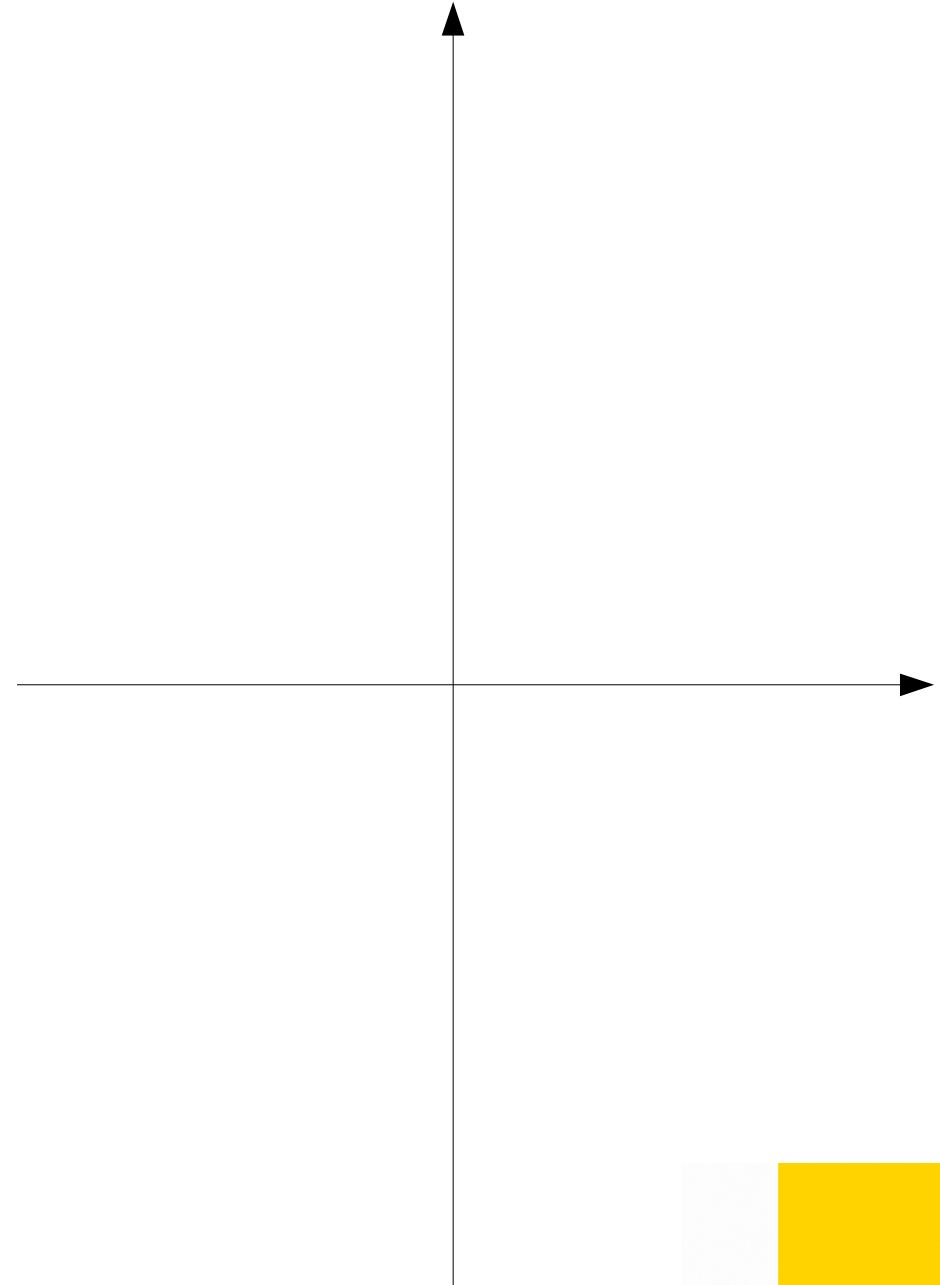


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- Impose **Landau gauge** condition
 - Reduces configuration space to a hypersurface



Unique gauge-fixing [For an introduction: Sobreiro & Sorella, 2005]

- Local gauge condition
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- Local gauge condition
 - Landau gauge: $\partial_{\mu} A_{\mu}^a = 0$
- This condition can be implemented using auxiliary fields, the so-called ghost fields
 - No physical objects: Pure mathematical convenience

(Perturbative) Landau gauge

- Lagrangian:

$$L = -\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a - \bar{c}^a \partial_\mu D_\mu^{ab} c^b$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - gf^{abc} A_\mu^b A_\nu^c$$

$$D_\mu^{ab} = \delta^{ab} \partial_\mu - gf^{abc} A_\mu^c$$
- Degrees of freedom:
 - Gluons: A_μ^a
 - Ghosts: \bar{c}^a, c^a
- Ghosts interact with gluons: They have to be included

Unique gauge-fixing [For an introduction: Sobreiro & Sorella, 2005]

- Local gauge condition
 - Landau gauge: $\partial_\mu A_\mu^a = 0$
- Sufficient for **perturbation theory**

Proceeding

- Once the gauge is fixed, all kind of (perturbative) calculations can be done
- Use **Green's or correlation functions** as basic entities

Green's Functions or correlation functions [Alkofer & von Smekal PR 2001]

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- Expectation values of a product of field operators
 - Build from the fields, here gluons and ghost
 - E.g.: $\langle \bar{c} c \rangle$

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- Full **Green's functions** contain all information
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- If having a non-vanishing color charge they change under **gauge transformation**
- Simplest non-zero **Green's functions: 2-point functions** or **propagators**
 - Expectation values of products of two field operators
 - **1-point functions** vanish

Propagators

- In Landau gauge: Gluon and one auxiliary field: Ghost
- Gluon:

$$D_{\mu\nu}^{ab}(x-y) = \langle A_{\mu}^a(x) A_{\nu}^b(y) \rangle$$

$$D_{\mu\nu}(p) = \left(\delta_{\mu\nu} - \frac{p_{\mu} p_{\nu}}{p^2} \right) \frac{Z(p)}{p^2}$$

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- Ghost:

$$D_G^{ab}(x-y) = \langle \bar{c}^a(x) c^b(y) \rangle$$

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- **Ghost propagator** can be expressed as a gluon operator, the inverse Faddeev-Popov operator

$$D_G^{ab}(x-y) \sim \langle (\partial_\mu D_\mu^{ab})^{-1} \rangle = \langle (\partial_\mu (\delta^{ab} \partial_\mu - g f^{abc} A_\mu^c))^{-1} \rangle$$

Proceeding

- Once the gauge is fixed, all kind of (perturbative) calculations can be done
- Use **Green's or correlation functions** as basic entities
- Combination of **gauge-variant Green's functions** yield gauge-invariant results
 - E.g. scattering cross-sections

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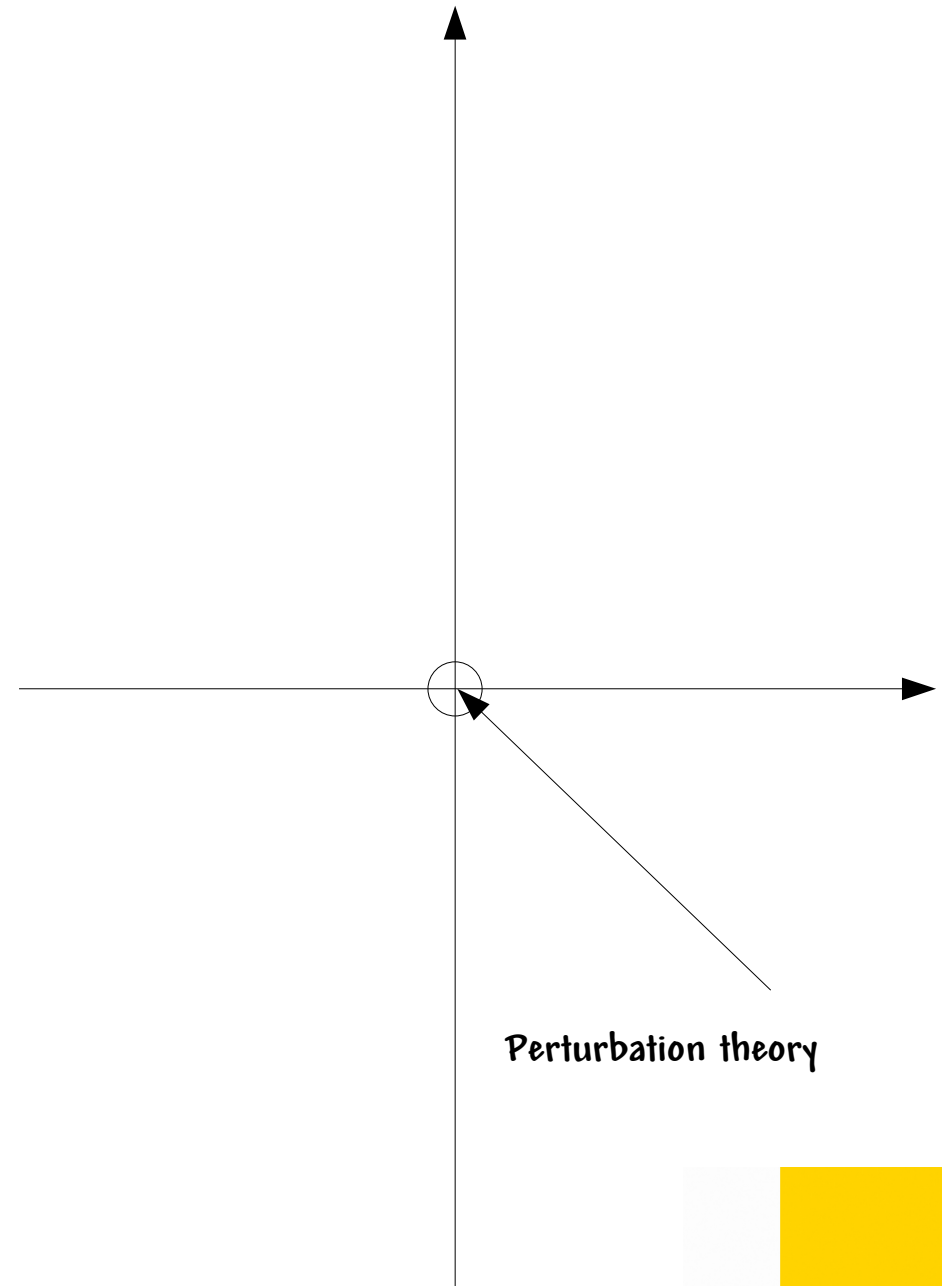
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- Almost all **perturbative calculations** proceed via **gauge-variant correlation functions**
- Also for the non-perturbative physics?

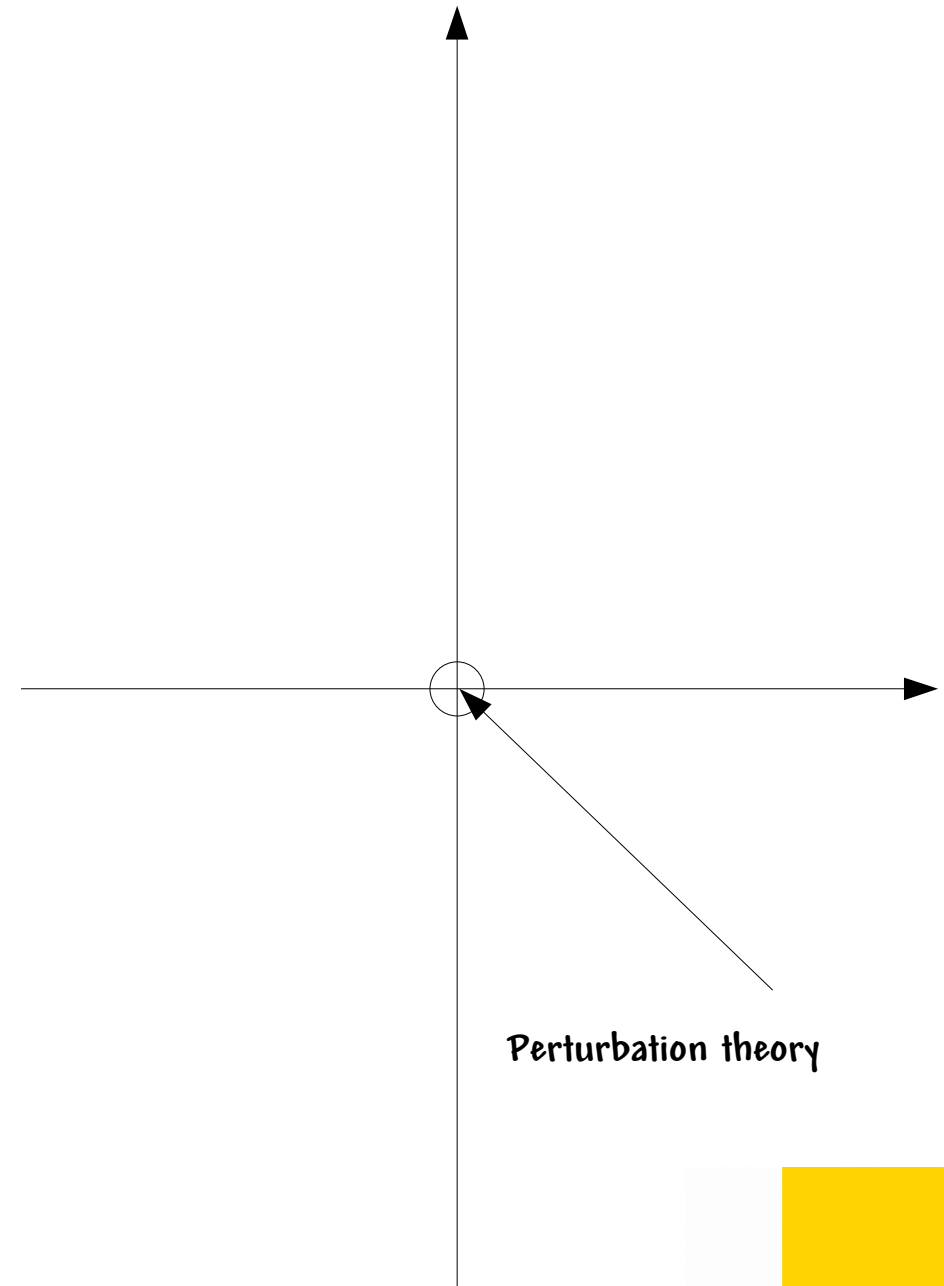
Configuration space (artist's view)

- **Perturbation theory** is applicable close to the origin



Configuration space (artist's view)

- **Perturbation theory** is applicable close to the origin
- Non-perturbative physics probes the complete hypersurface



Unique gauge-fixing [For an introduction: Sobreiro & Sorella, 2005]

- Local gauge condition
 - Landau gauge: $\partial_\mu A_\mu^a = 0$
- Sufficient for perturbation theory
- Insufficient beyond perturbation theory
 - There are gauge-equivalent configurations which obey the same local gauge-condition: Gribov copies [Gribov 1978]

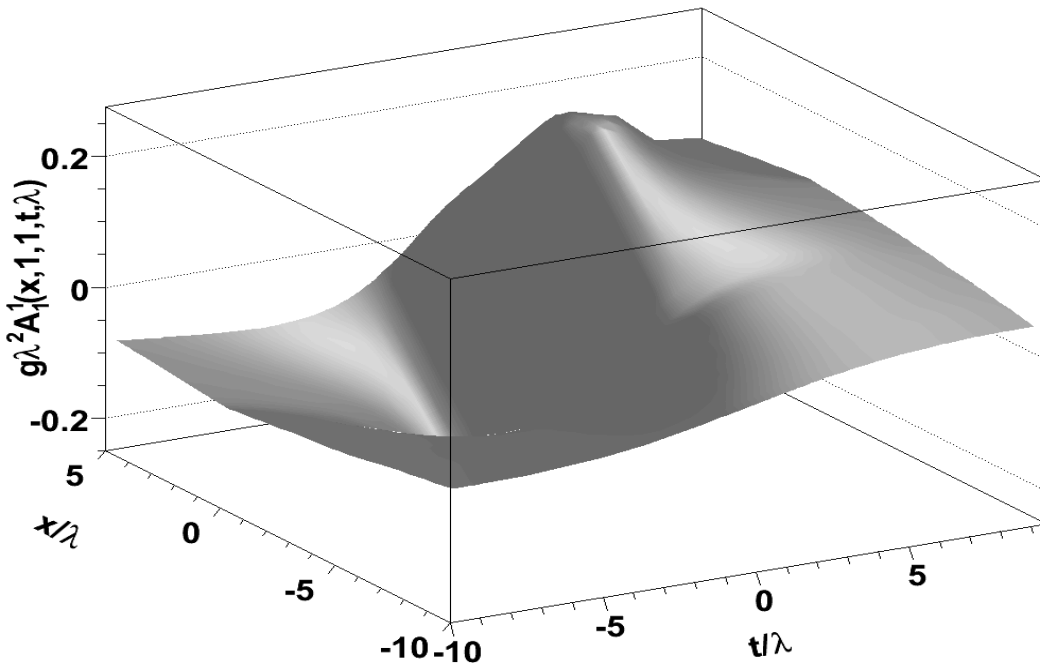
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- There are no local gauge conditions known, which select a unique gauge field configuration [Singer 1978]
 - Non-local conditions possible

Example: Instanton

[Maas, EPJC 2006]

Instanton field

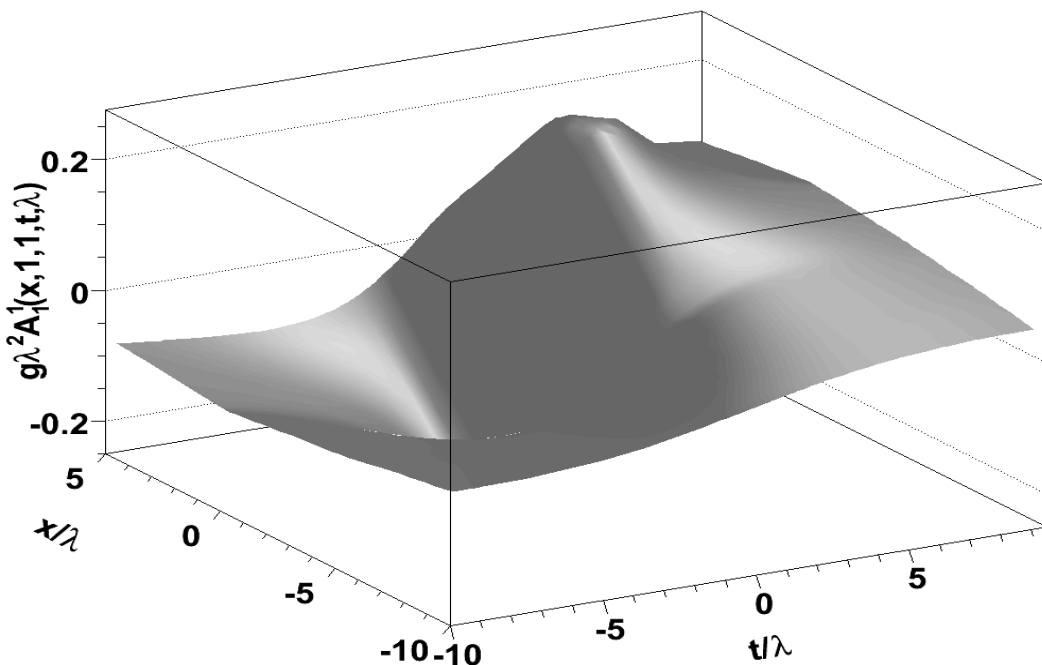


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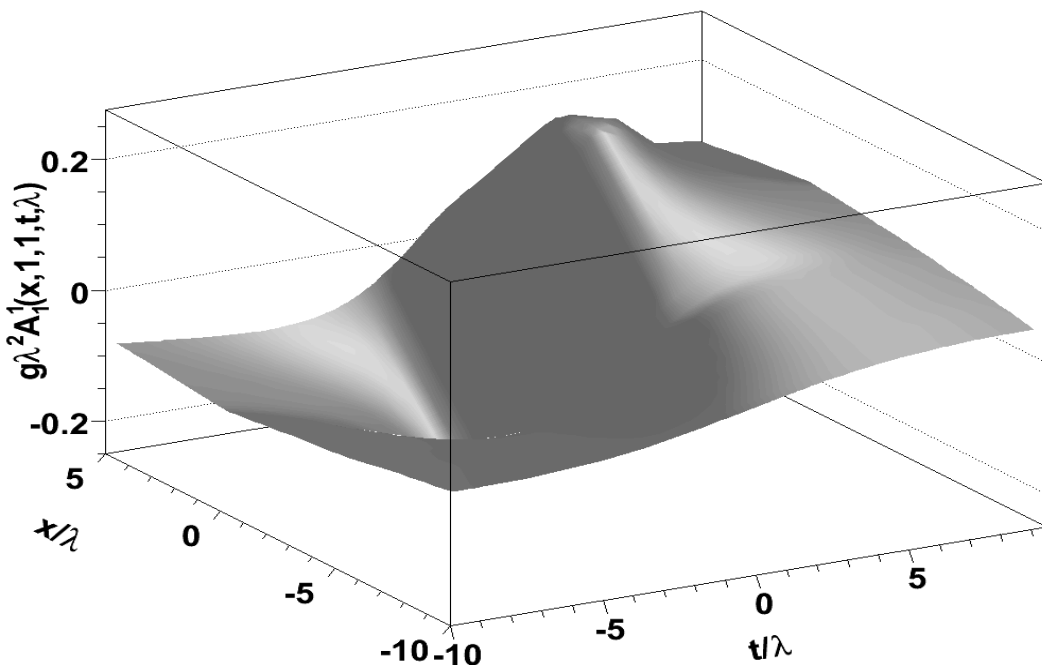


- Instanton field configuration is $A_{\mu}^a(r, \lambda) = 2r_{\nu} \eta_{\nu\mu}^a / (g(r^2 + \lambda^2))$
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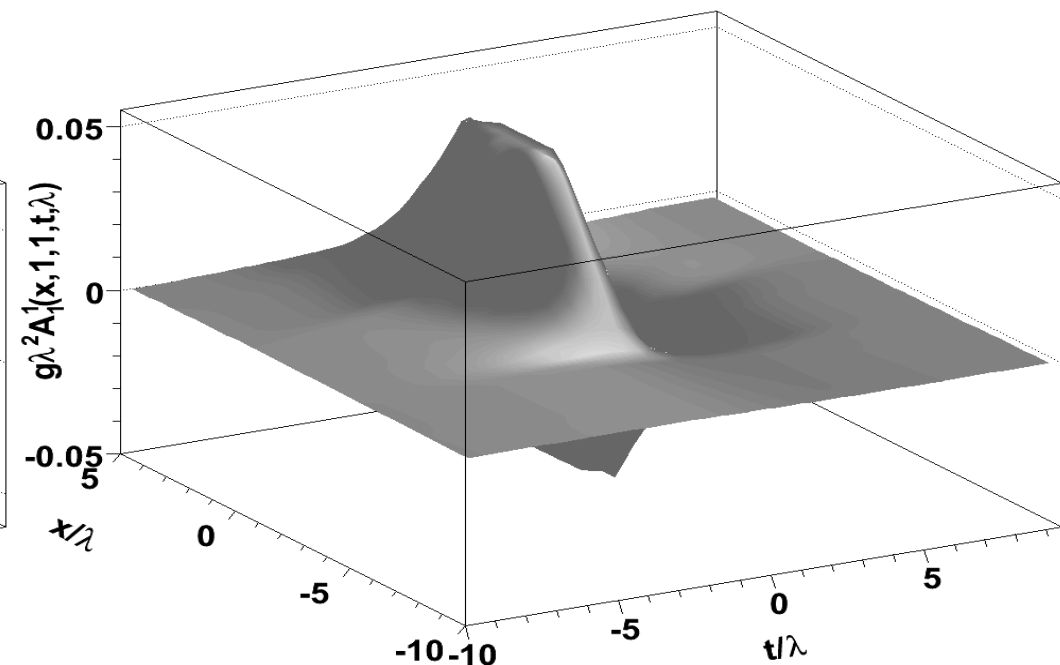
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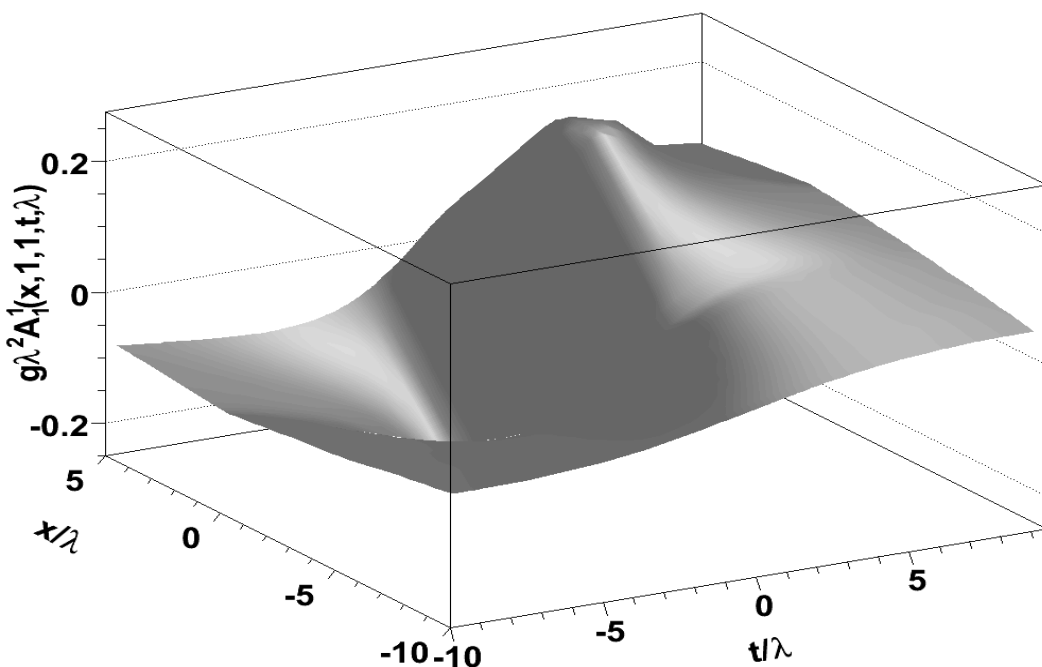


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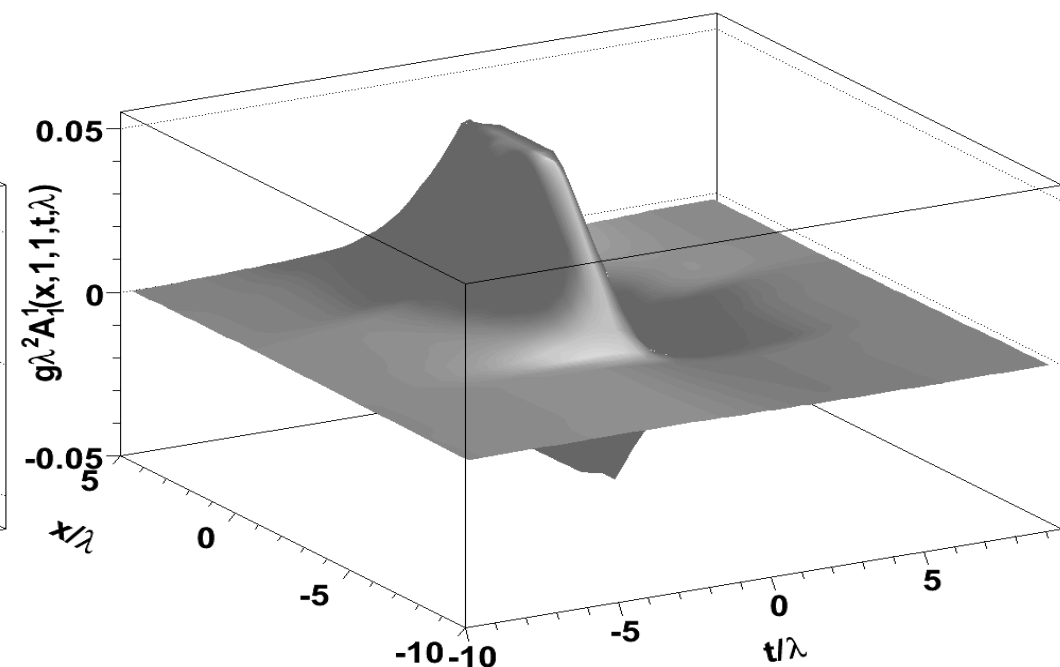
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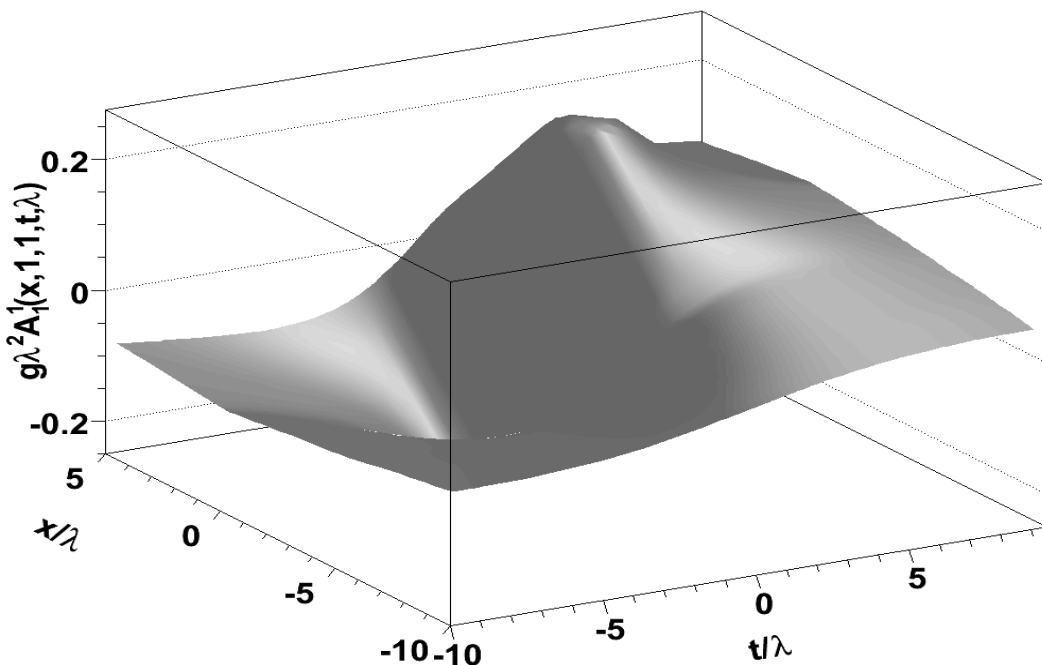


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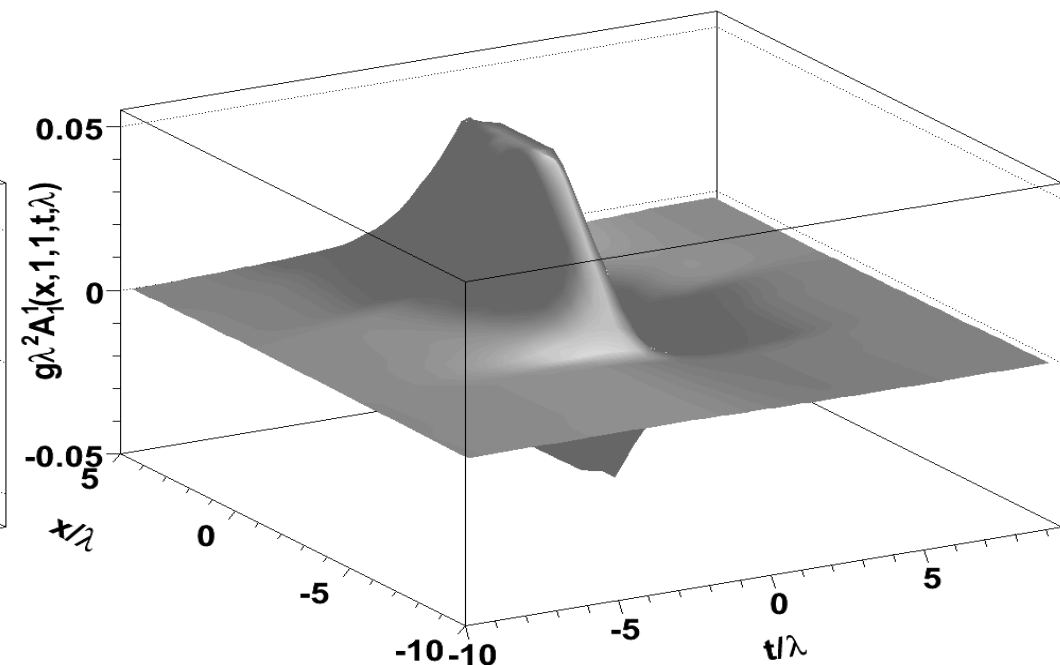
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 - **Gribov copy**
 - Non-perturbative: Depends on $1/g$

Residual freedom

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 - Reduces configuration space to a hypersurface

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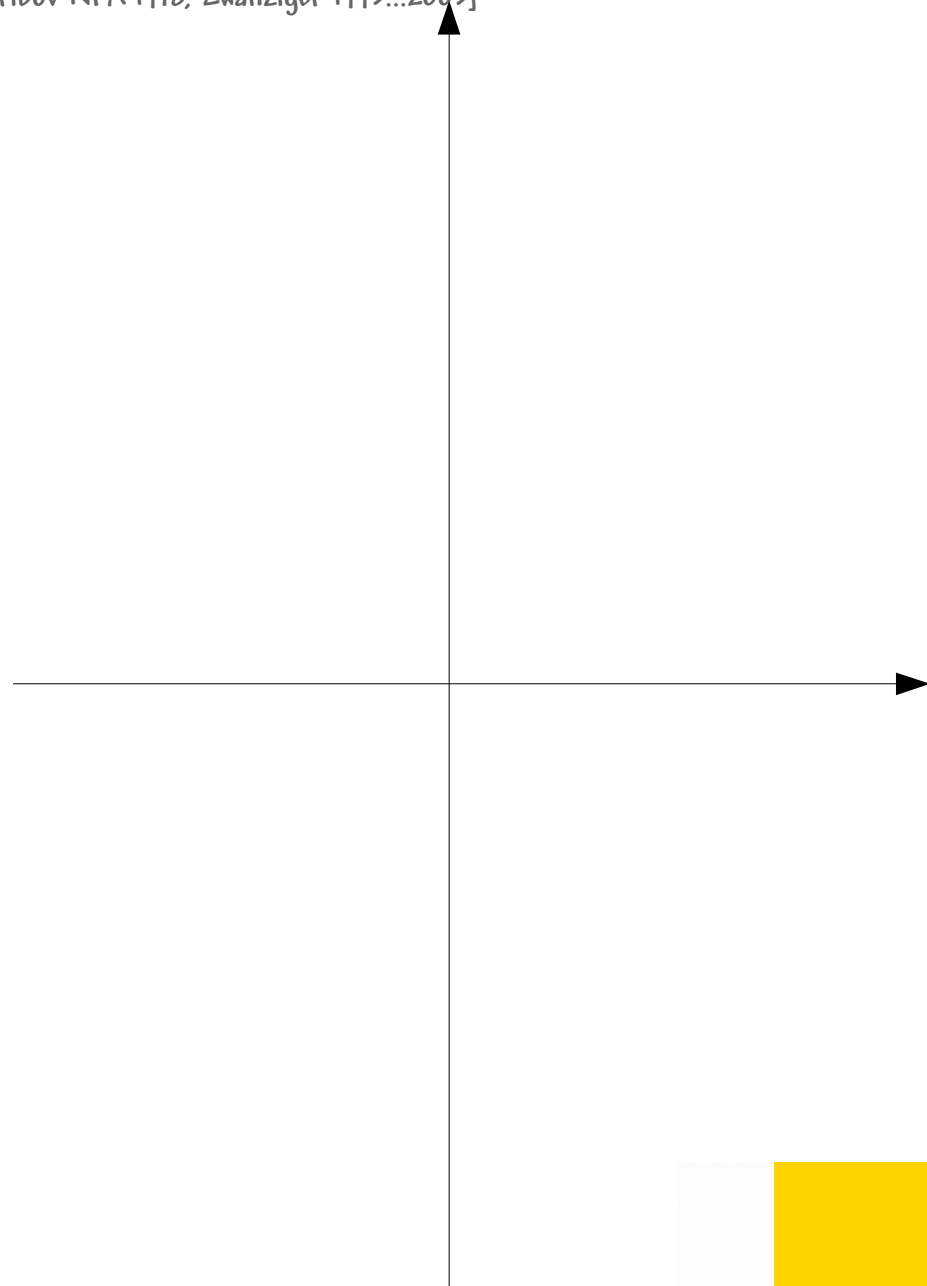
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- Construct a non-local condition instead to solve the problem
- Choice: Leave the global color symmetry unfixed

Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

- Absolute Landau gauge



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- Absolute Landau gauge

- A global minimum of

$$-\int d^d x A_\mu^a(x) A_\mu^a(x) \sim -\int d^d p D_{\mu\mu}^{aa}(p)$$

defines the **fundamental modular region**

- This region is bounded and convex



Fundamental modular region

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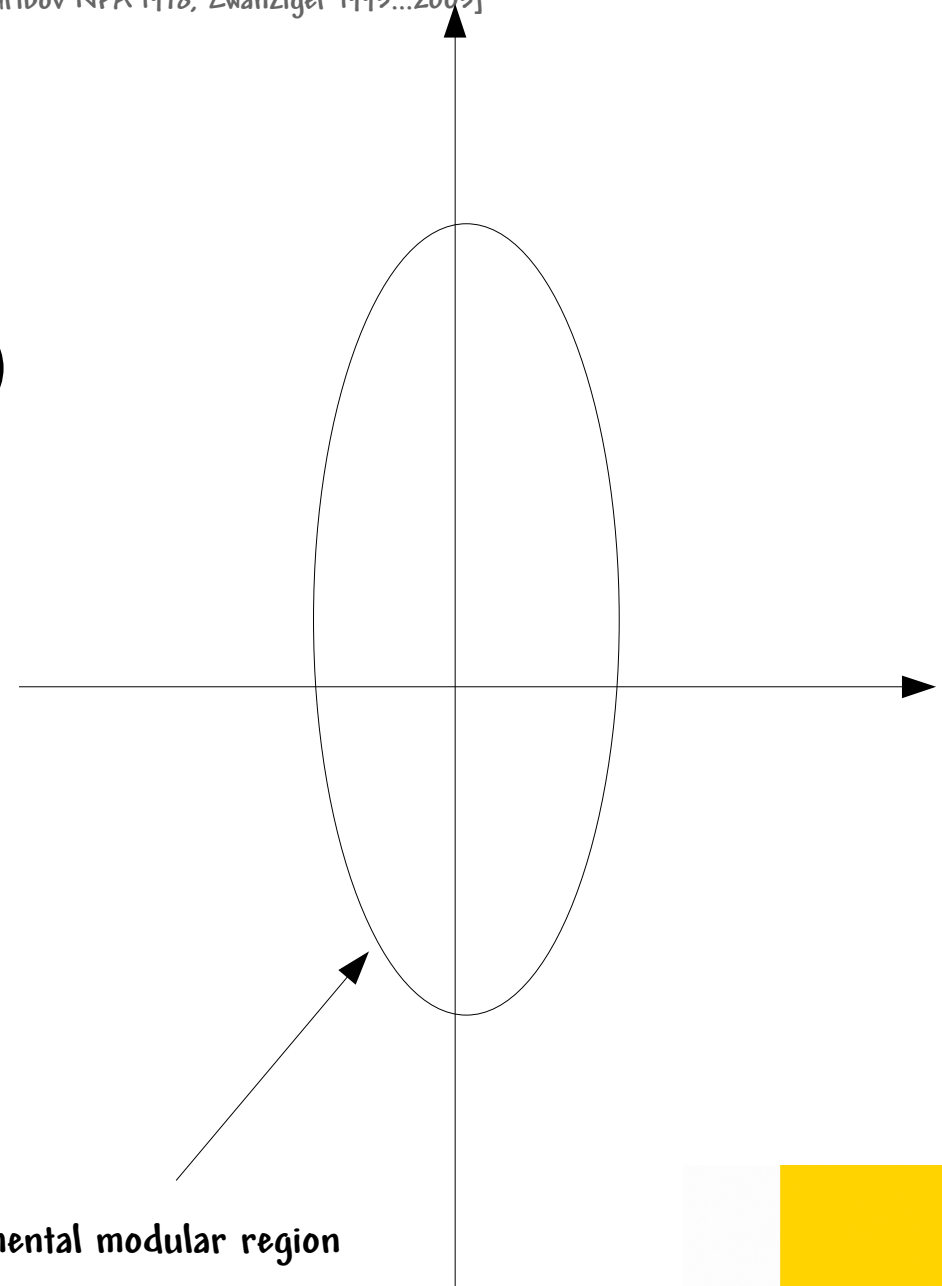
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- Singles out exactly one gauge copy
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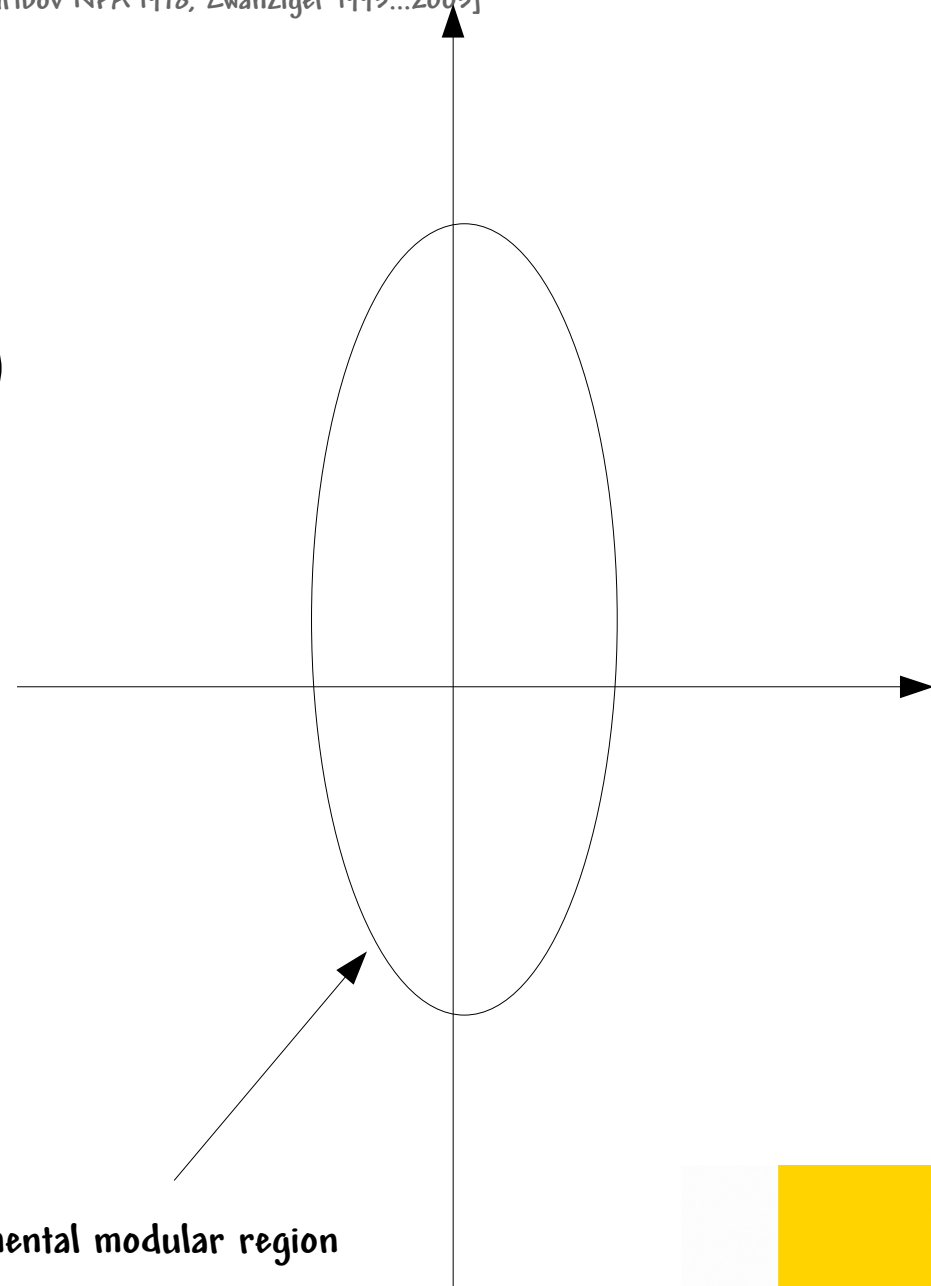
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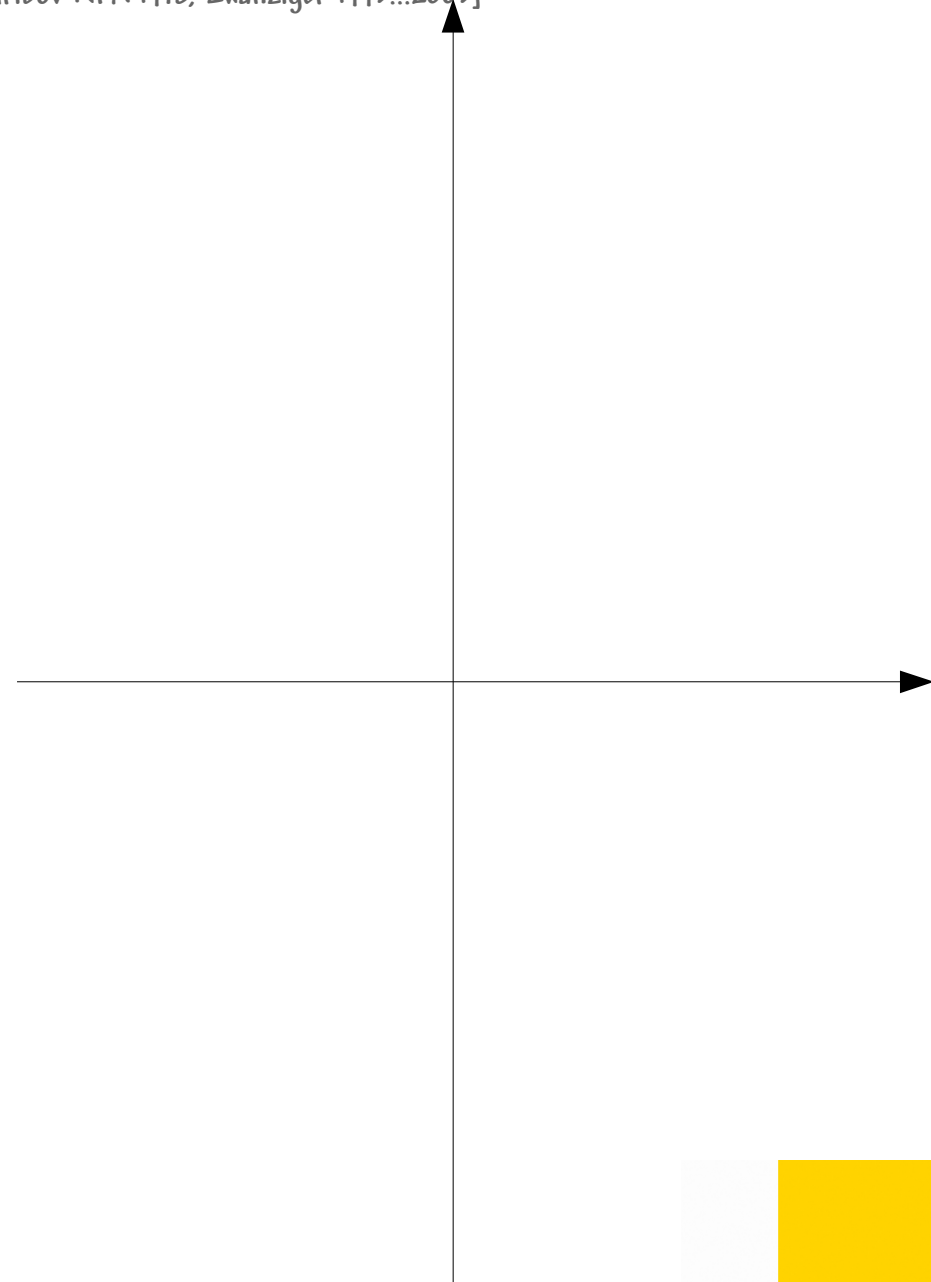
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- Drawback: Badly divergent and regularization may introduce degeneracies



Fundamental modular region

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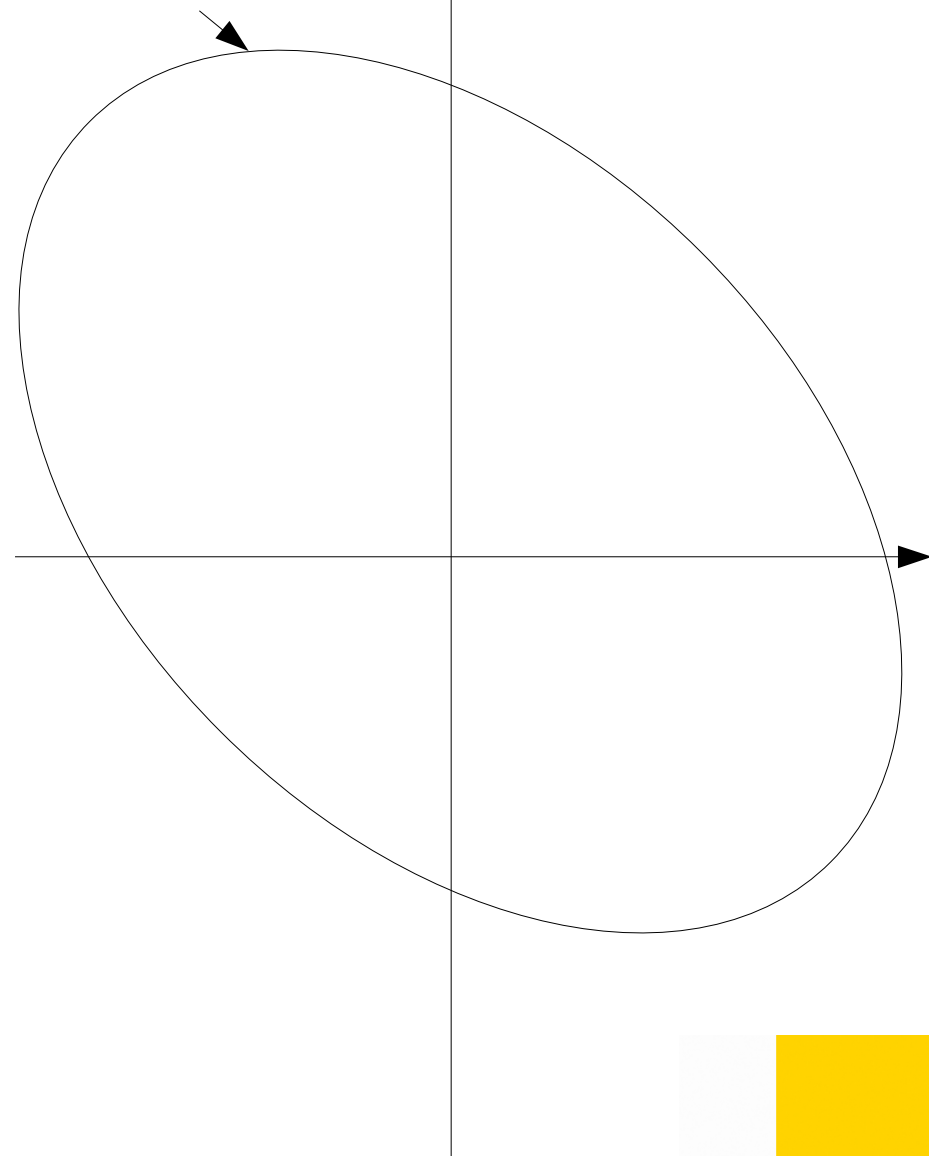
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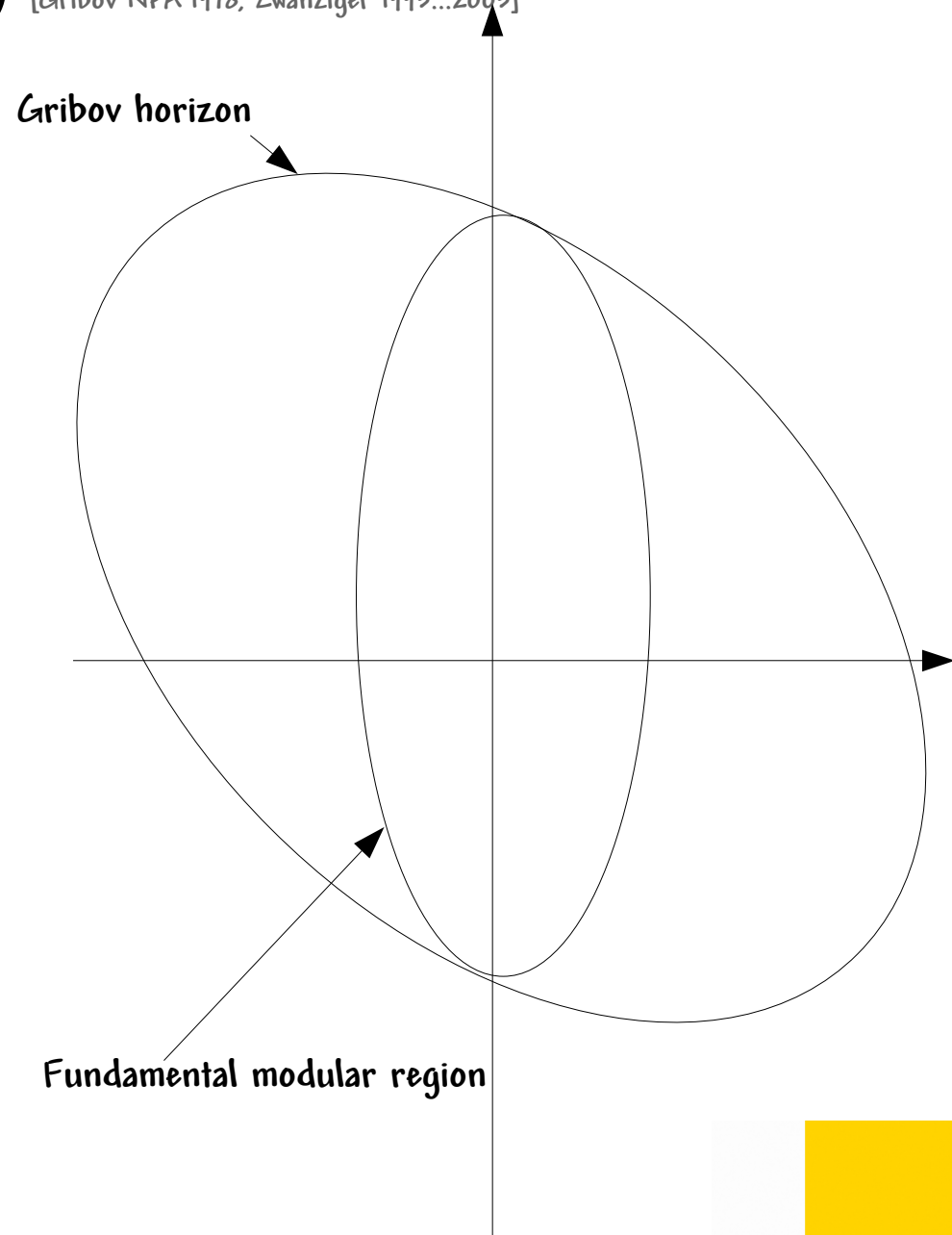
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Gribov horizon



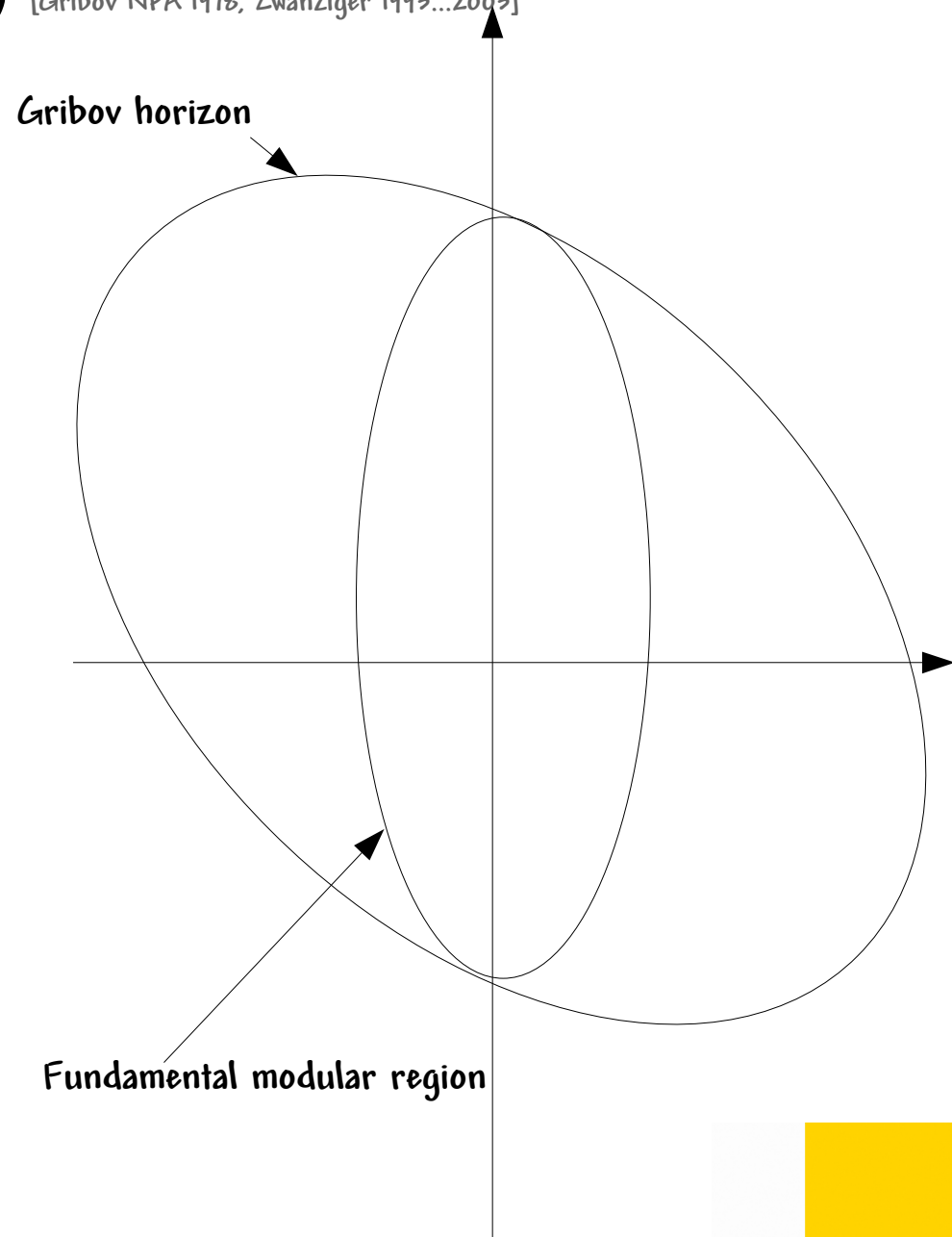
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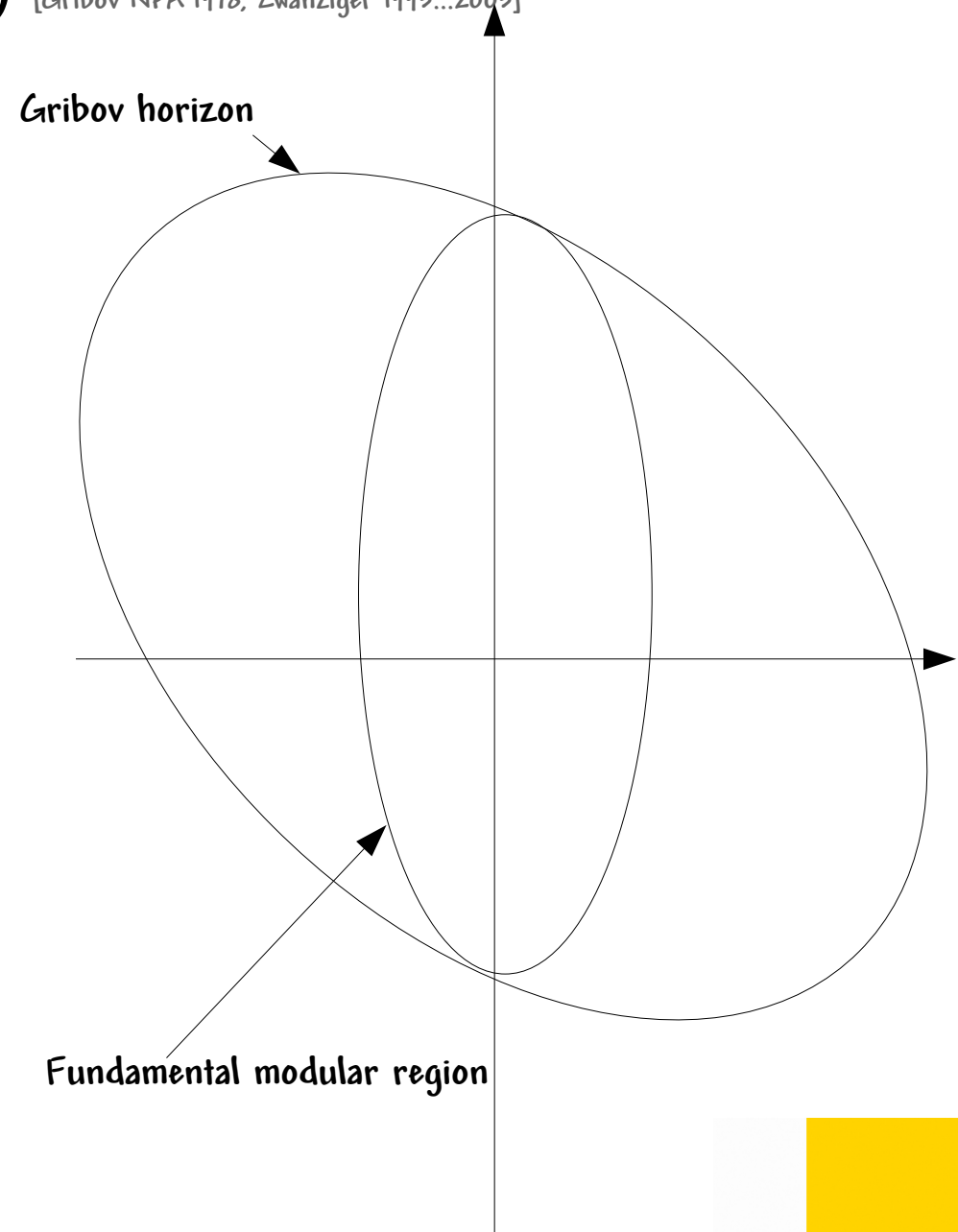
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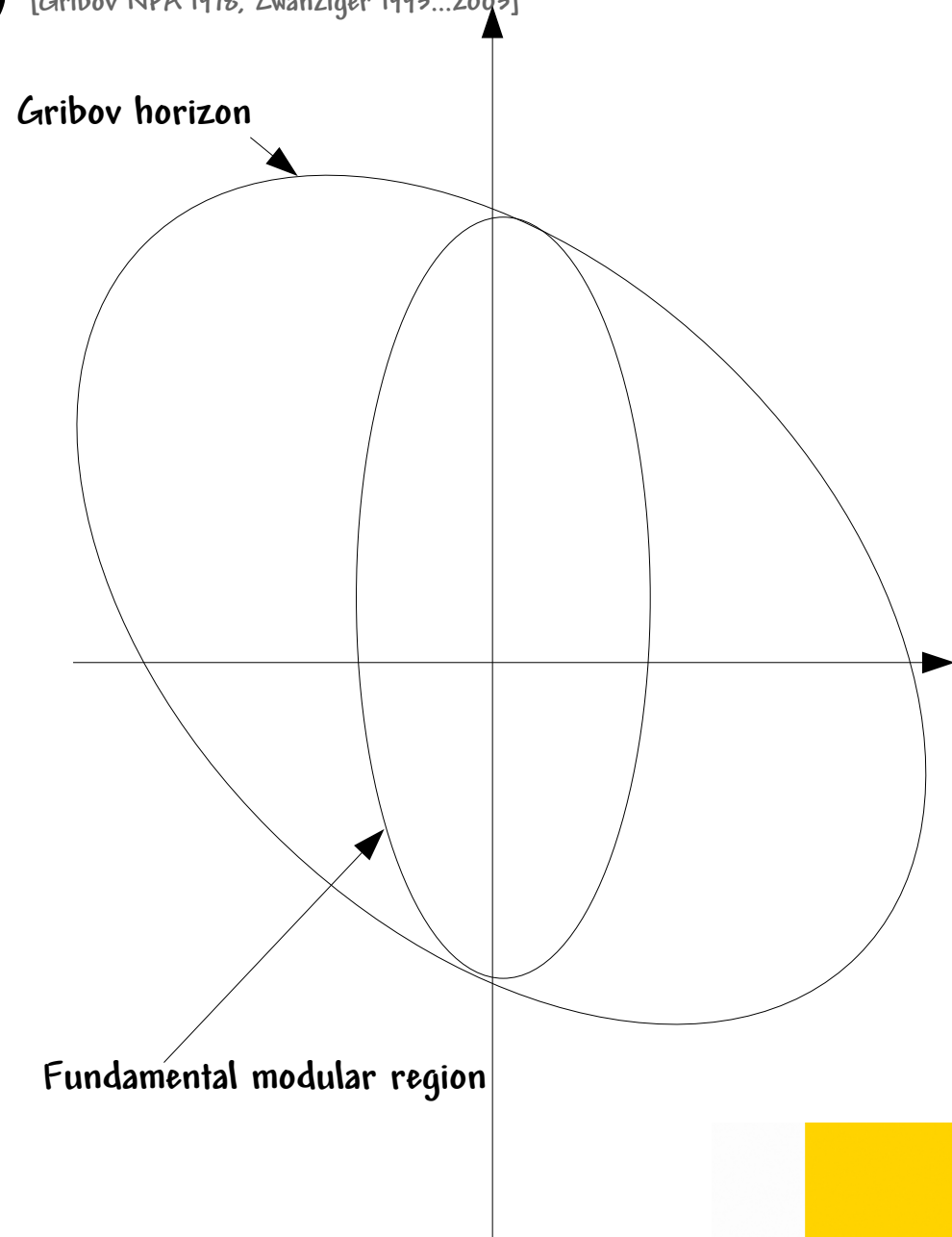
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 - How many is many?



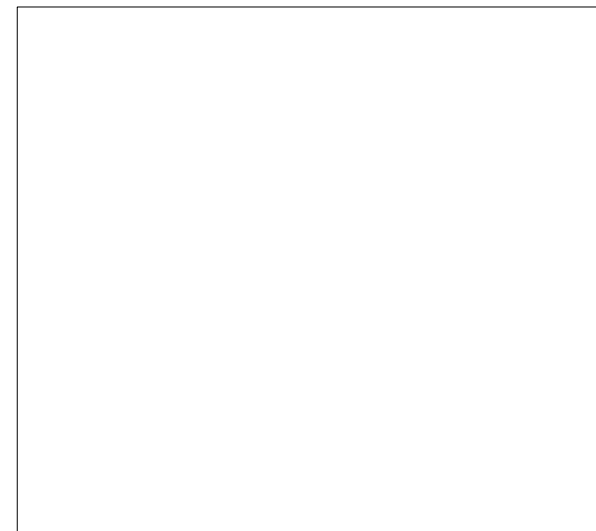
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 - How many is many?
 - Requires a **tool** for investigations



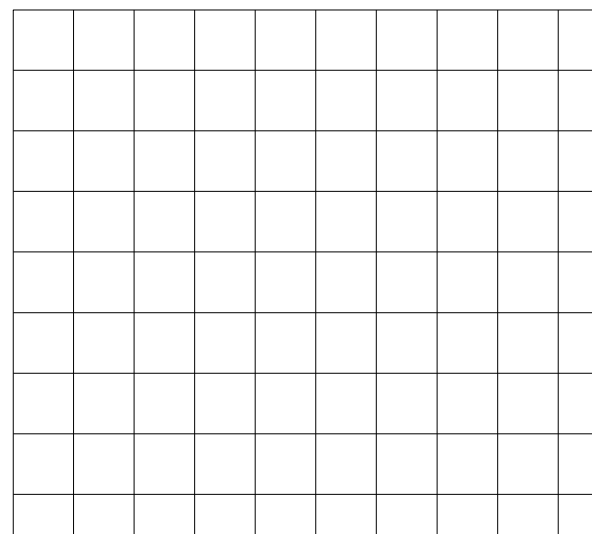
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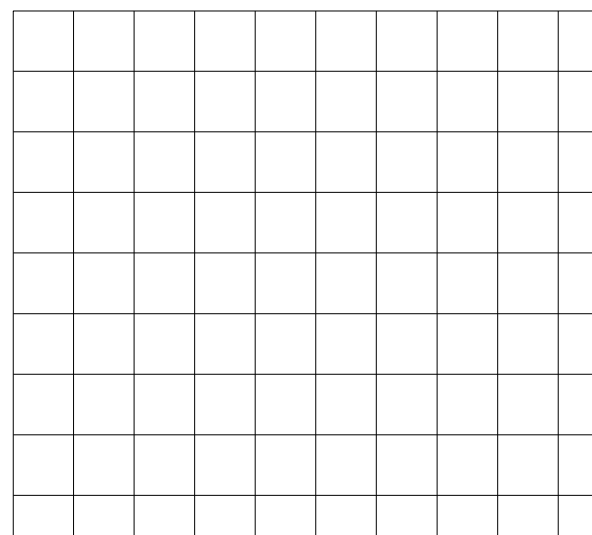
Lattice calculations

- Take a **finite volume** – usually a hypercube
- Discretize it, and get a **finite, hypercubic lattice**



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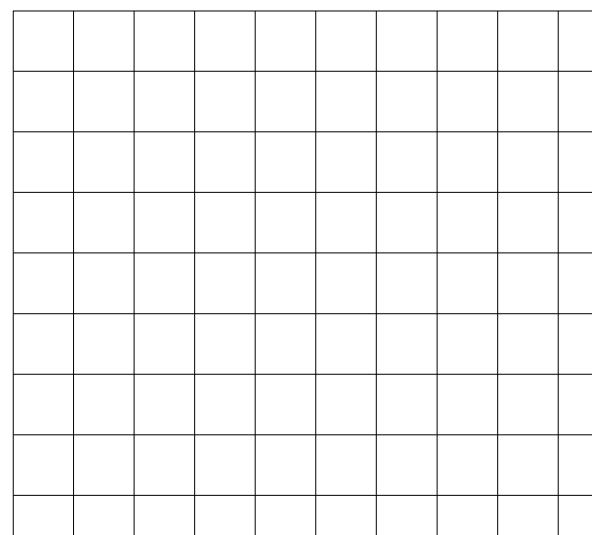
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- **Artifacts**
 - Finite volume/discretization
 - Zero momentum problematic



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Identifying Gribov copies

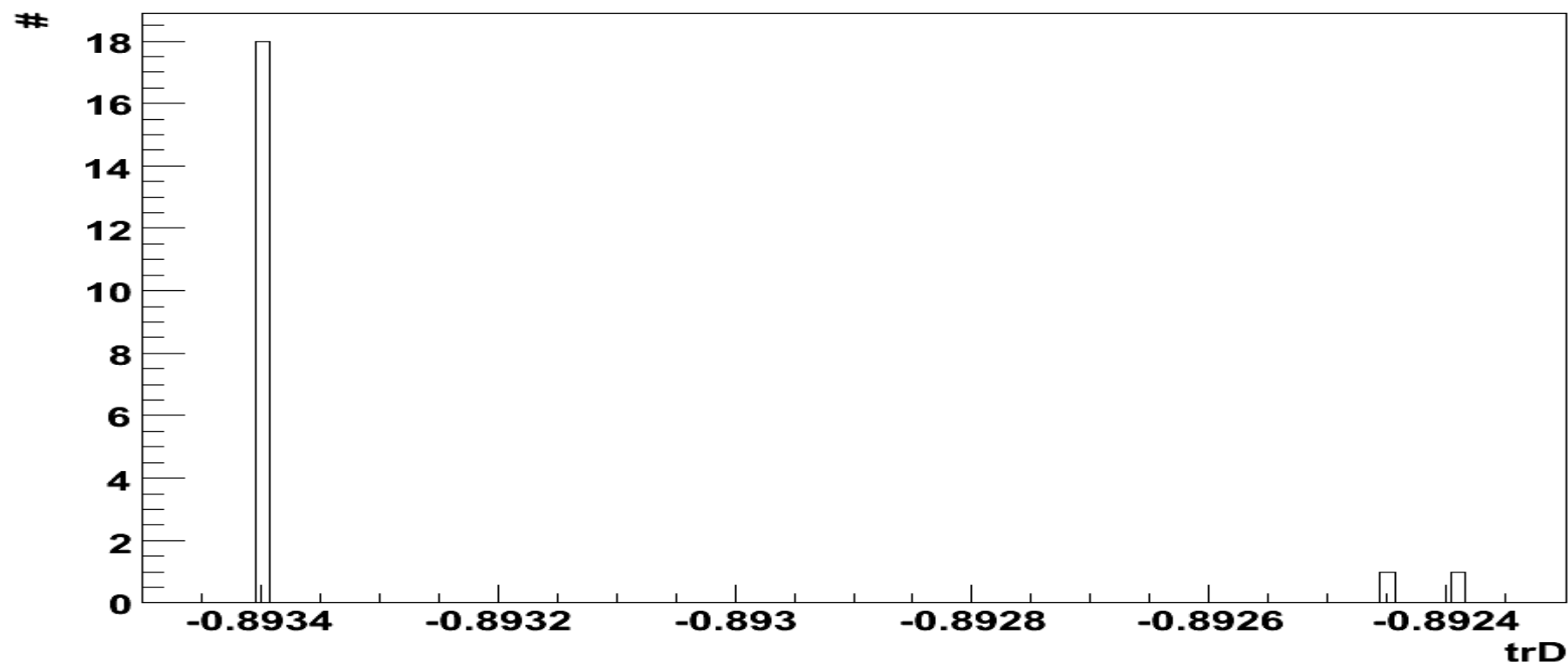
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 - How to decide whether two copies are different?
 - Non-trivial problem. Requires in principle to show a connection by a non-trivial gauge transformation
 - Practical solution: Require sufficiently different value of $\text{Tr}D$
 - Other are dismissed as **numerical/lattice artifacts**
 - Is this good? Will possibly work only on a **lattice!**

TrD on the residual gauge orbit [3d, beta=3.46, Maas, unpublished]

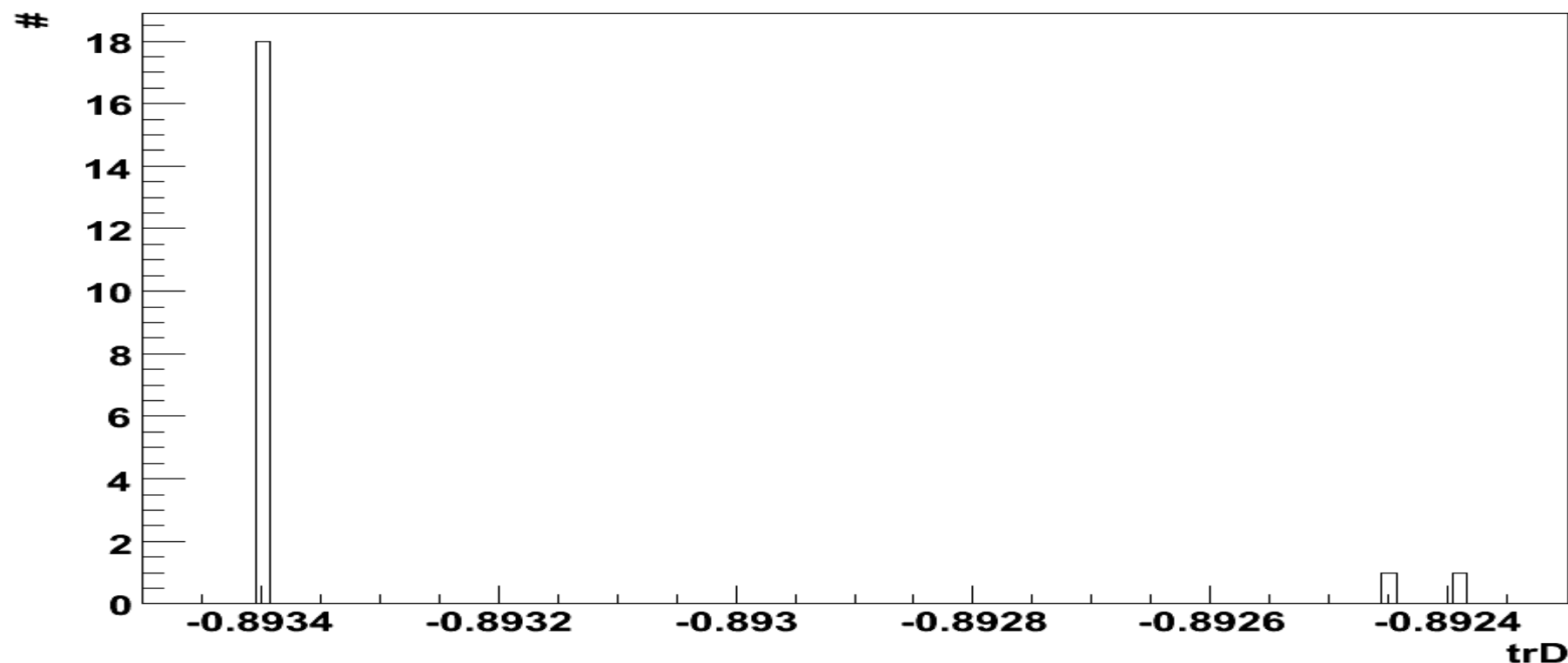
TrD for $V=(3.1 \text{ fm})^3$



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- Many artificial copies of one particular copy: Attractive basin in TrD

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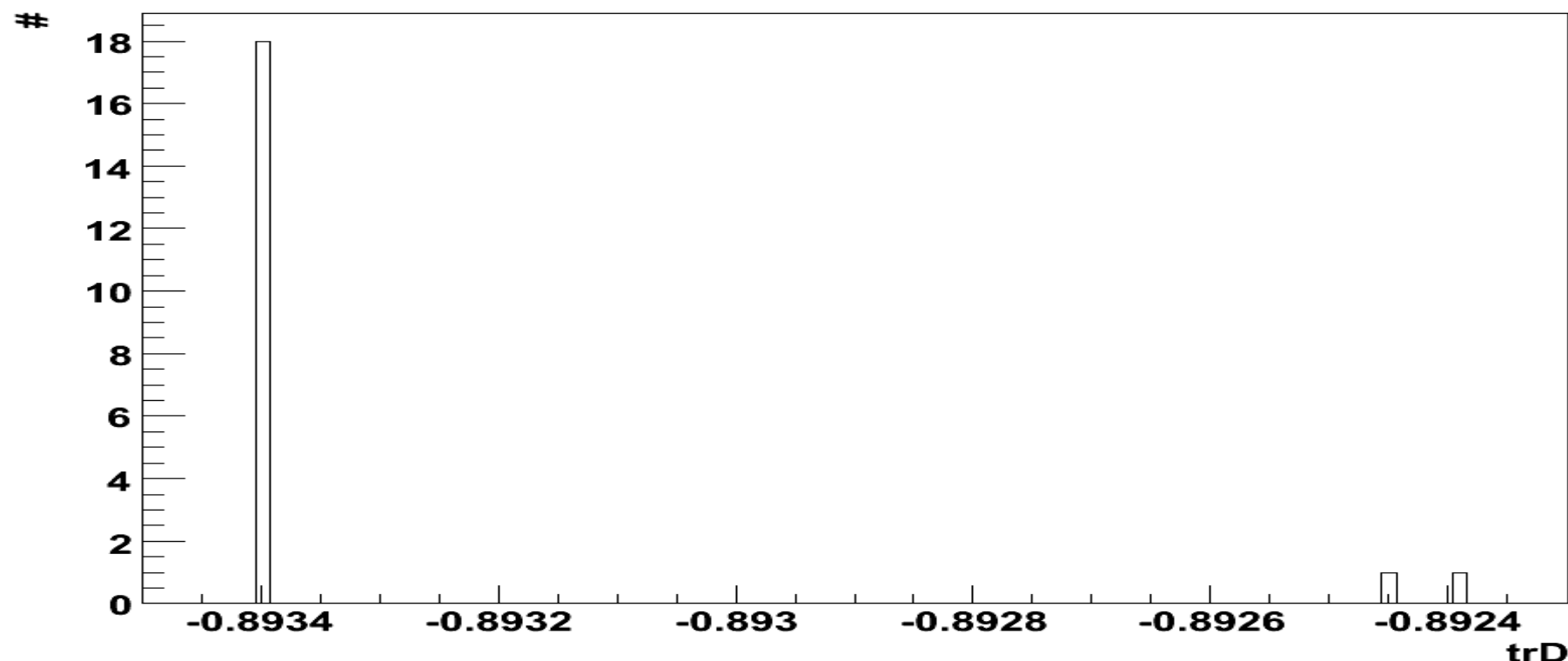
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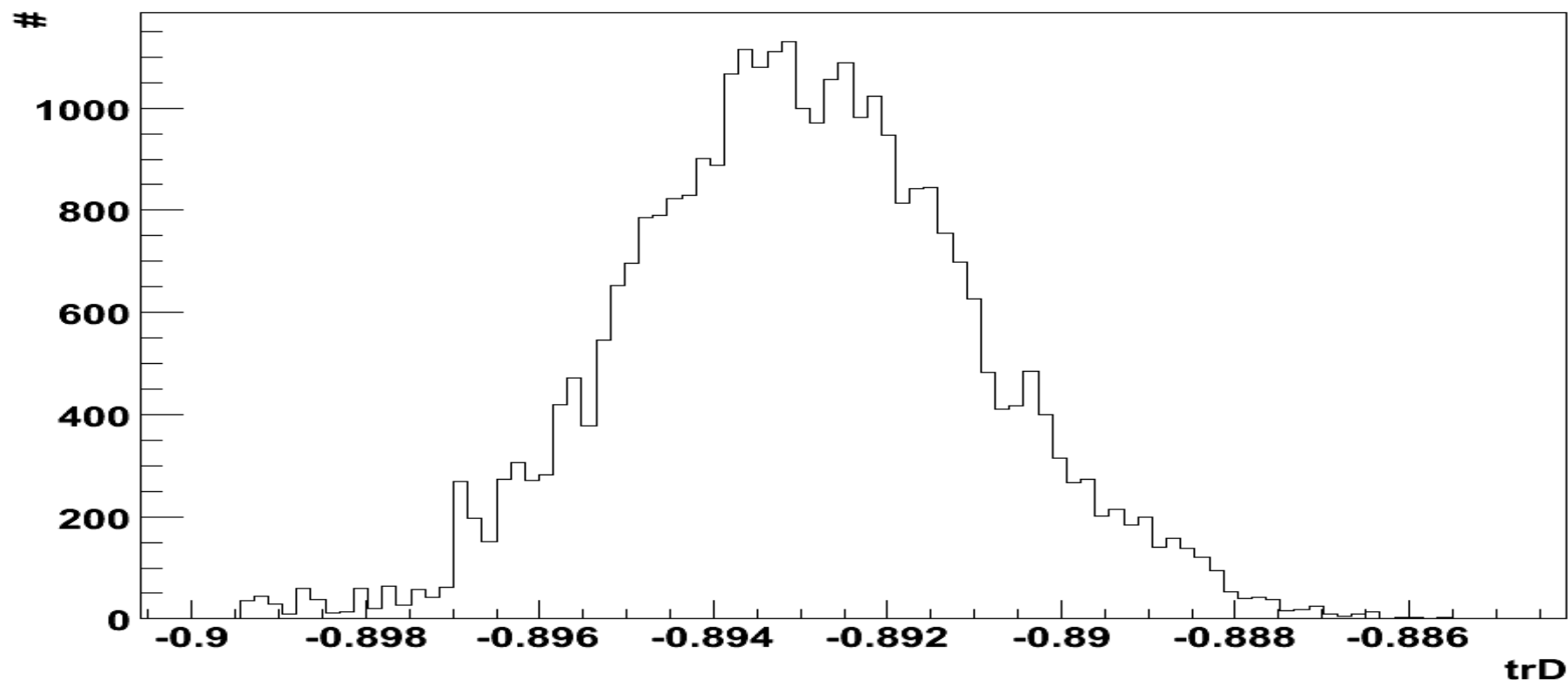
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- Difference in peak is below 10^{-10} : Numerical/lattice artifacts
- Pragmatic definition: Copy requires a separation of 10^{-5} in TrD

Distribution of TrD [3d, beta=3.46, Maas, unpublished]

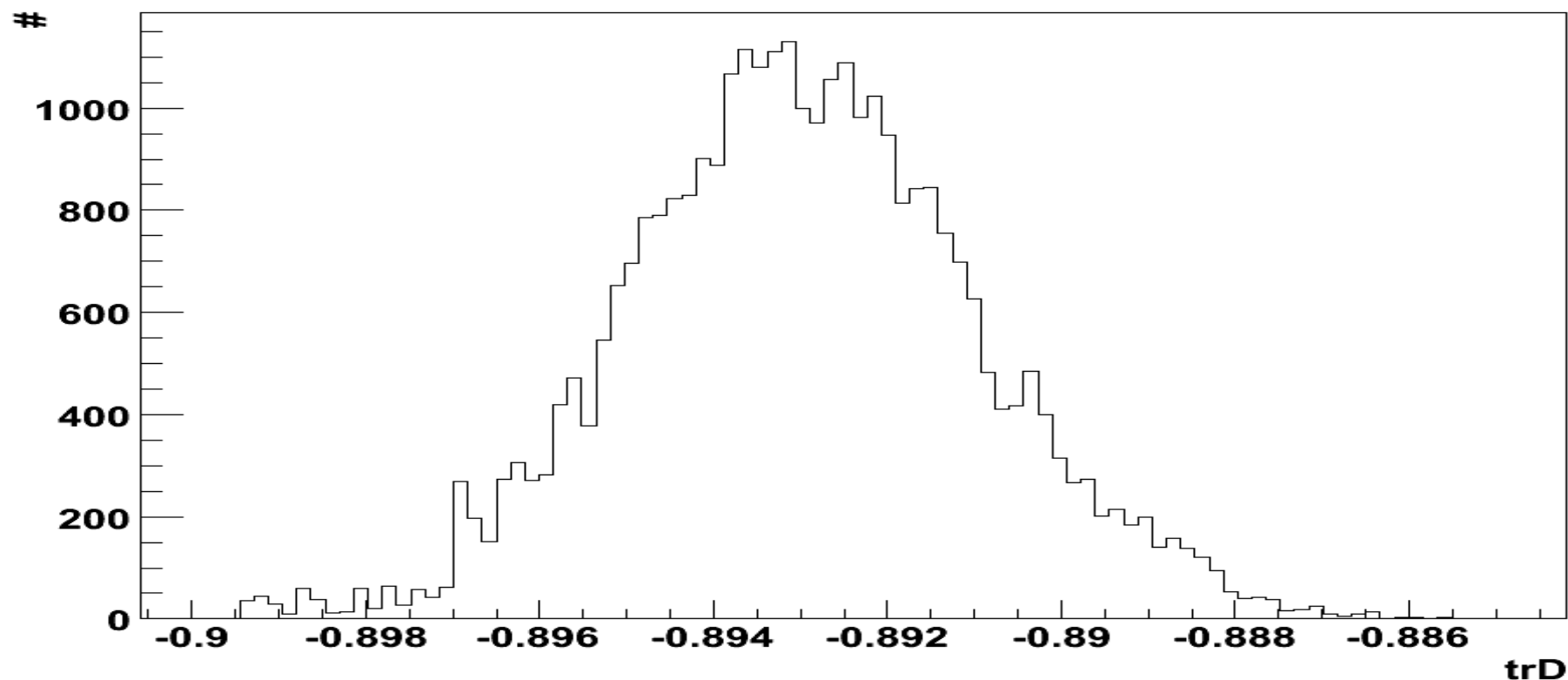
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 - Never a real Gaussian

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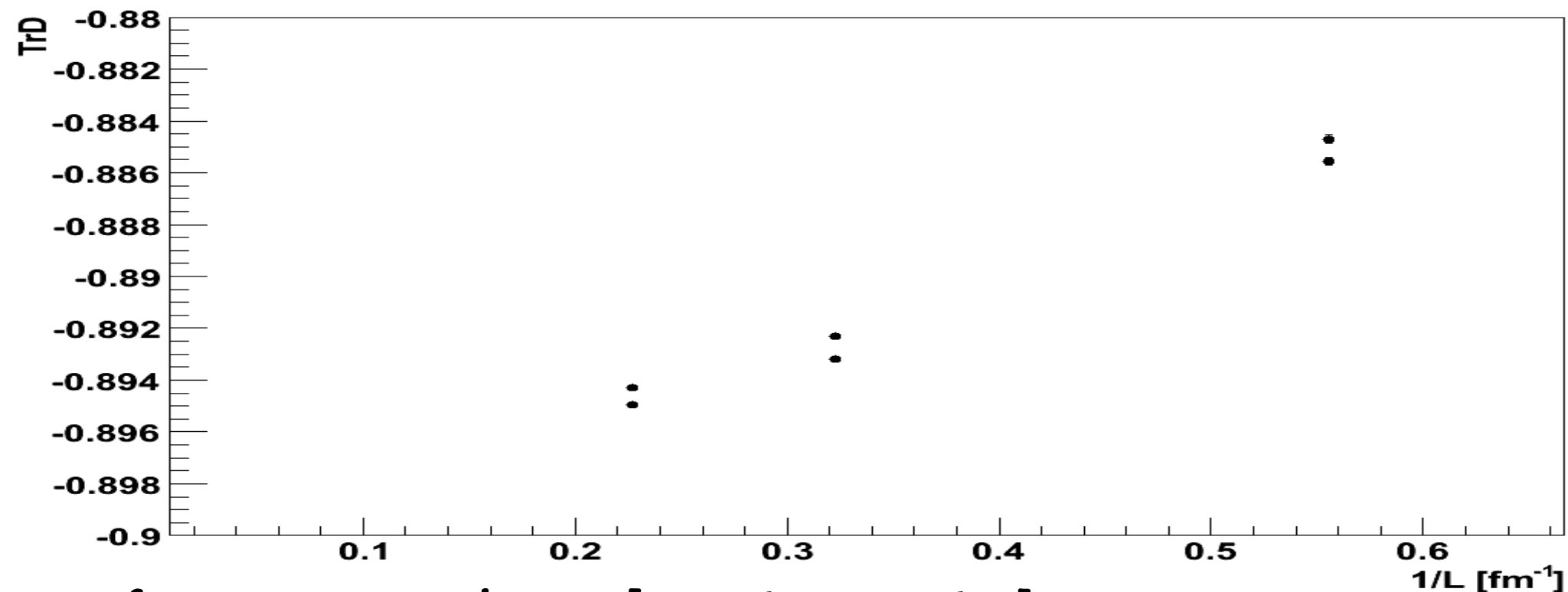
- Distribution is almost Gaussian
 - Never a real Gaussian
- Variation of TrD corresponds to varying gluon propagator

TrD as a gauge parameter [3d, beta=3.45, Maas, unpublished]

- Define gauge corridor as $[\max \text{TrD}, \min \text{TrD}]$

TrD as a gauge parameter [3d, beta=3.45, Maas, unpublished]

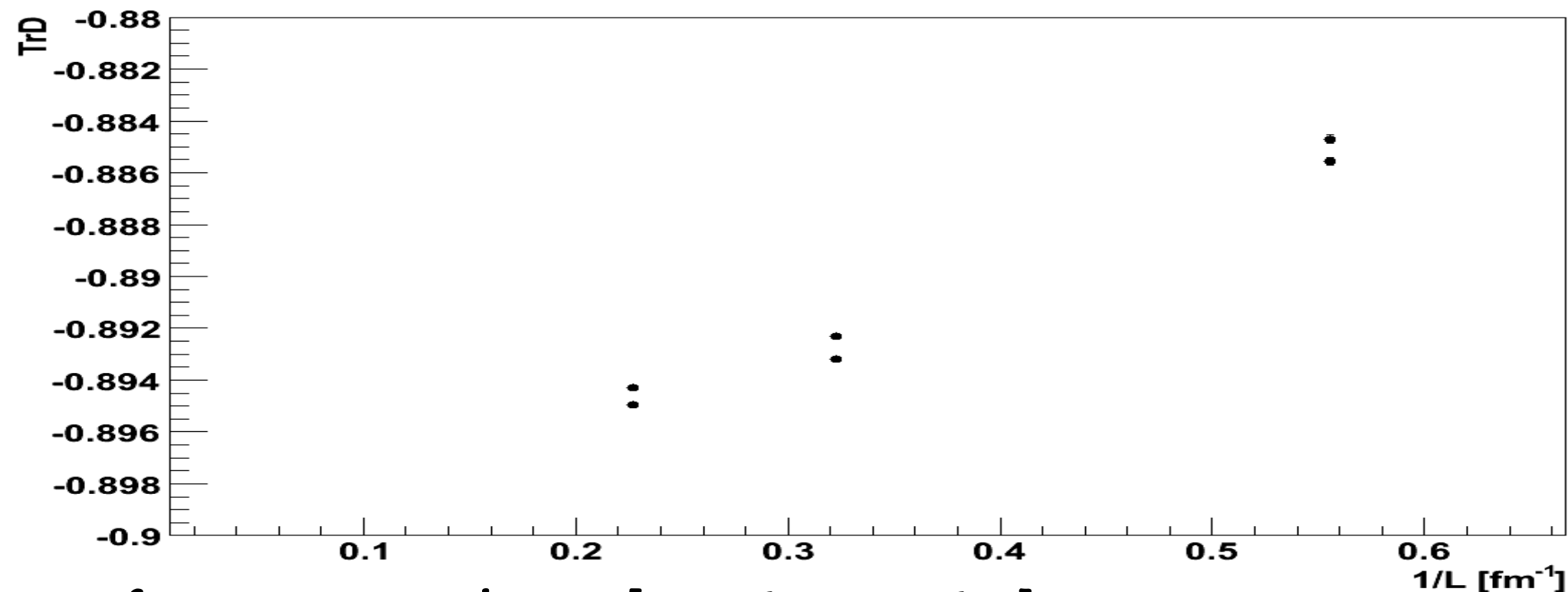
TrD corridor



- Define gauge corridor as [max TrD, min TrD] [Zwanziger, 2003]

TrD as a gauge parameter [3d, beta=3.45, Maas, unpublished]

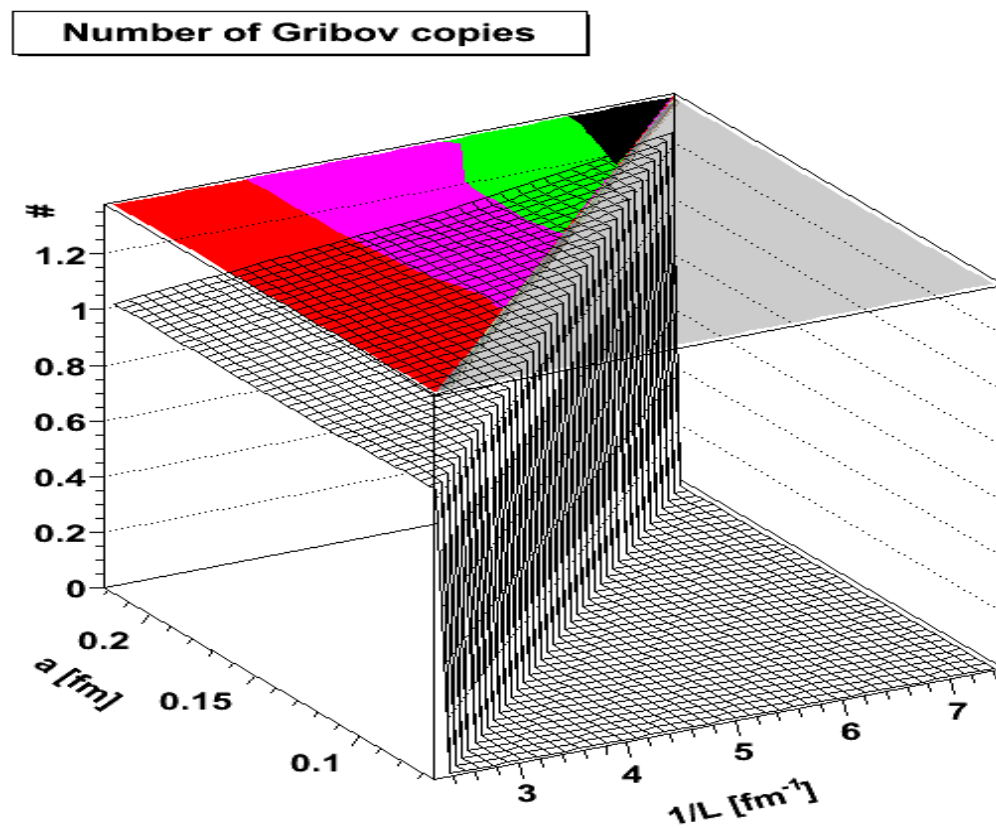
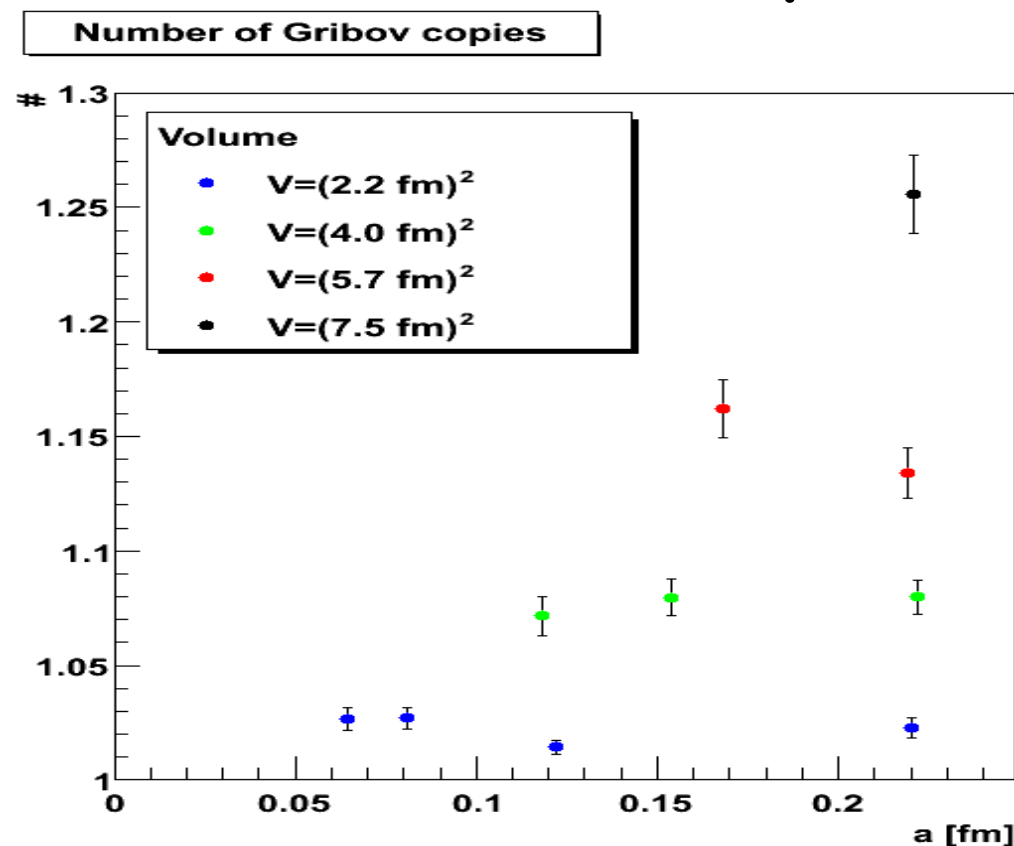
TrD corridor



- Define gauge corridor as [max TrD, min TrD]
- TrD is not showing a strong volume dependence
 - For some theories known to degenerate [de Forcrand, 1994]
 - Degeneracies would support the Zwanziger entropy conjecture

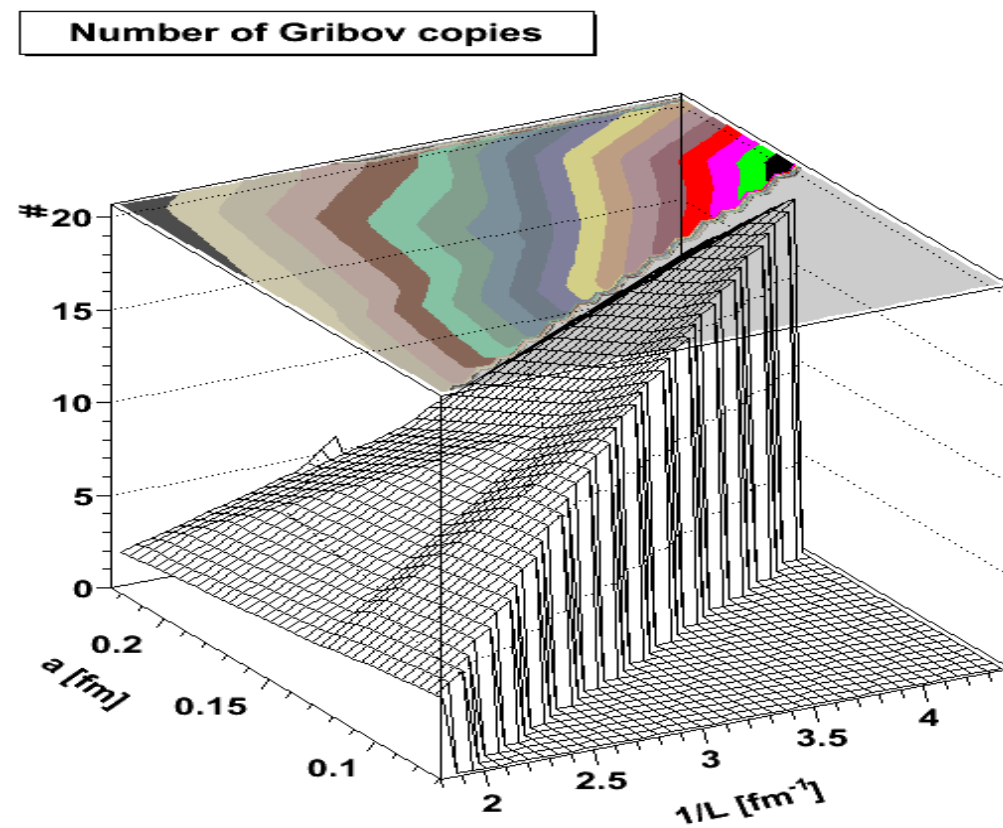
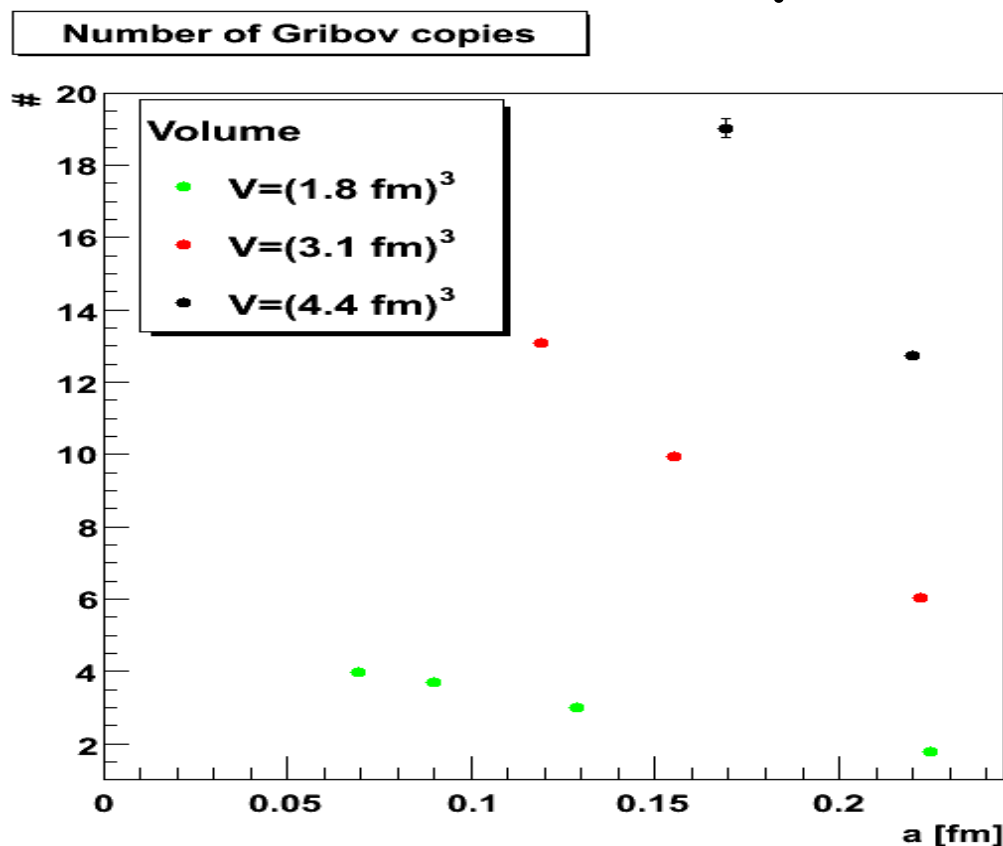
[Zwanziger, 2003]

Number of Gribov copies in two dimensions [Maas, unpublished]



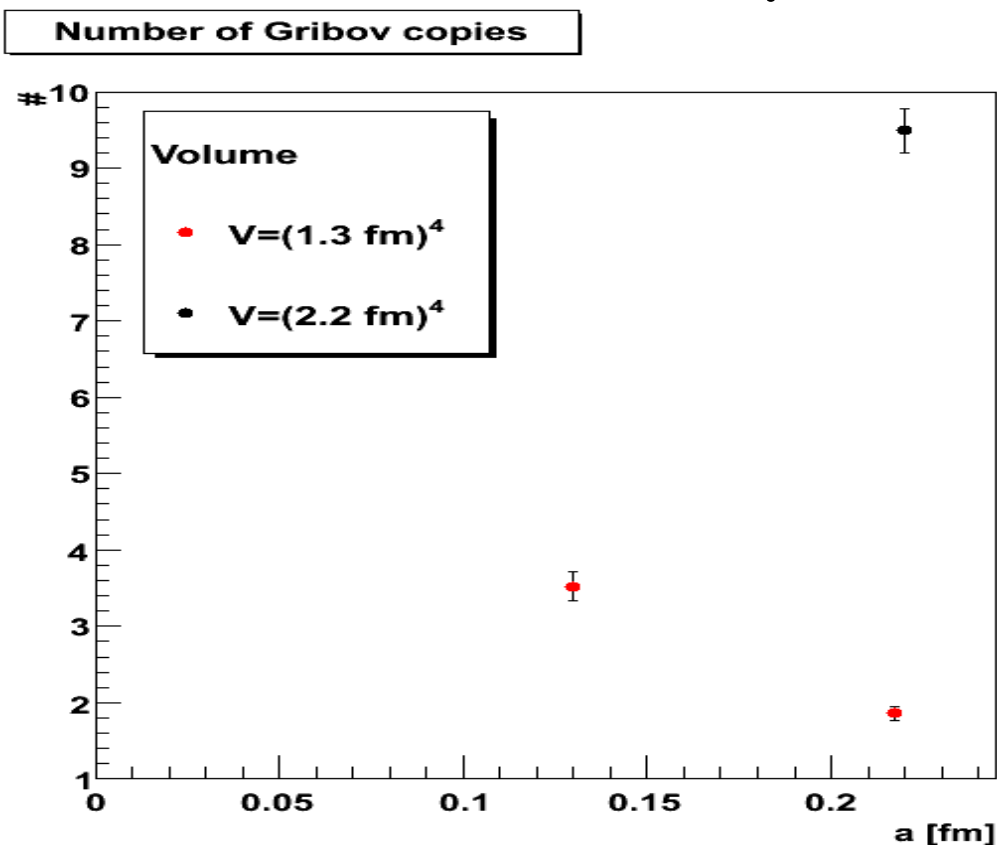
- Very slow evolution with volume and discretization
- Very large volumes will be needed to see an effect

Number of Gribov copies in three dimensions [Maas, unpublished]



- Evolution with volume and discretization
- Already rather small volumes have a large number of copies

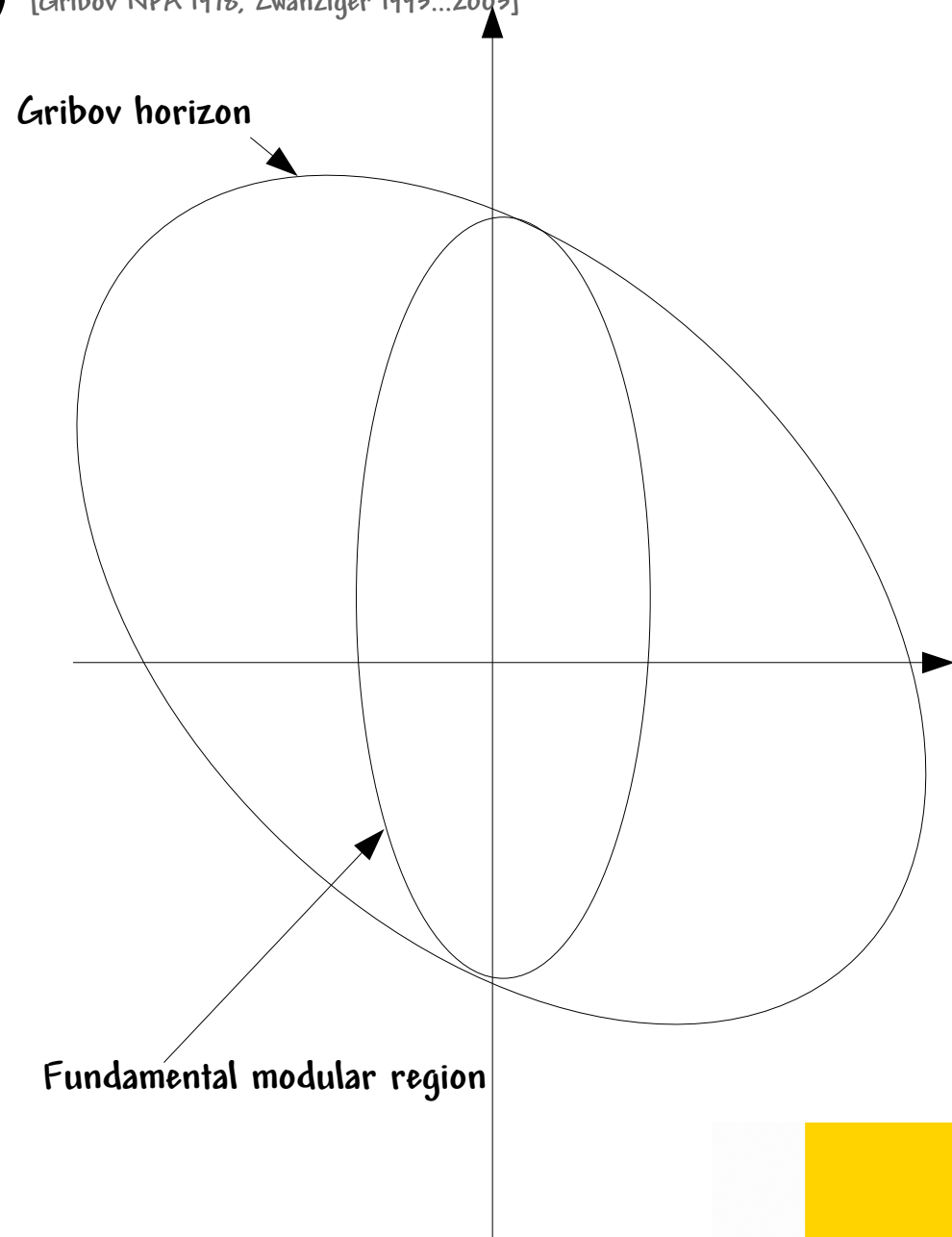
Number of Gribov copies in four dimensions [Maas, unpublished]



- Large number of copies even on small volumes
- Big effects can be expected
- Number of Gribov copies rises quickly with the number of dimensions

Configuration space (artist's view) [Gribov NPA 1978, Zwanziger 1993...2003]

- Different approach: Enlarge the search space
- **Gribov horizon** encloses all field configurations with positive Faddeev-Popov operator $(-\partial_\mu D_\mu)$
- Includes the **fundamental modular region**
- All gauge orbits pass through this region
 - Many **Gribov copies** for each
- How to select a unique representative?



Constructing a global gauge condition

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- Not directly. But the **ghost propagator** is influenced by the eigenspectrum:

$$\frac{-G(p)}{p^2} = D_G^{ab}(p) \sim \langle (\partial_\mu D_\mu^{ab})^{-1} \rangle \sim \sum_i \frac{1}{\lambda_i} |\psi_i|^2$$

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- **$G(p)$** candidate for a characterization of a copy

Constructing a global gauge condition [Maas, unpublished, Maas, 2008, Fischer et al. 2008]

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- Is the renormalization-group invariant $b = G(0)/G(P)$ with $P > 0$ fixed a possible second gauge parameter?
 - Fourier-transform at low momentum: Highly non-local
 - Does it vary along the (residual) gauge orbit inside the first Gribov region?

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- Investigate using the lattice
 - Not possible to access $G(0)$ on the lattice
 - Take lowest momentum instead: Finite-volume Artifact

Investigating the residual gauge orbit

- There are many copies along the residual gauge orbit

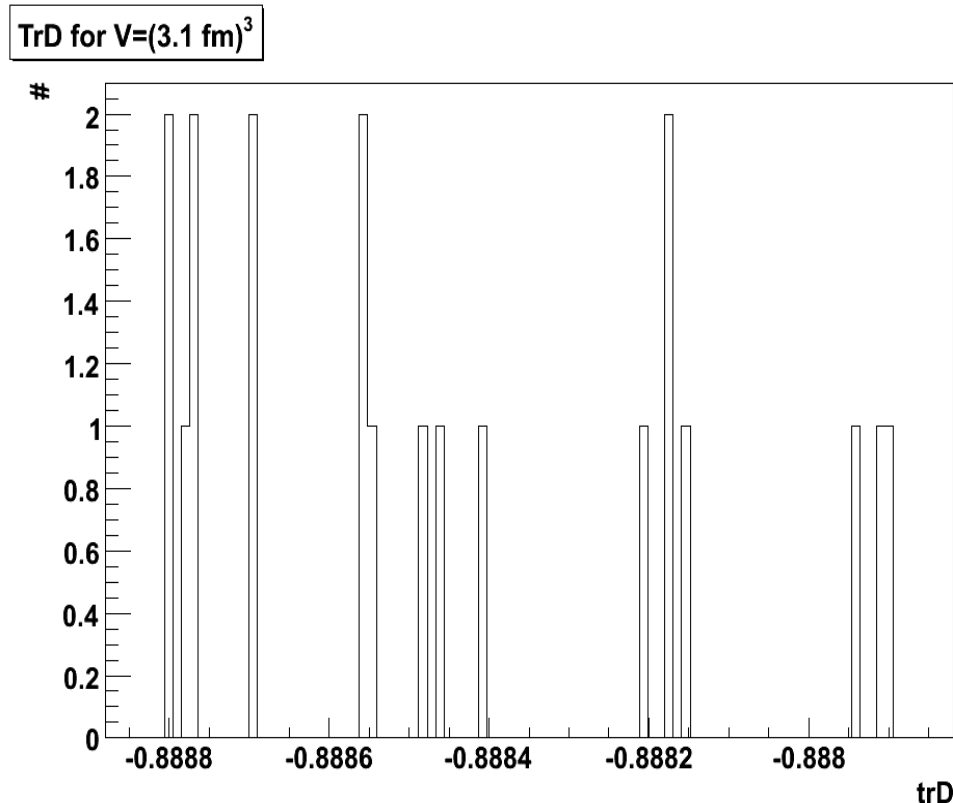
Investigating the residual gauge orbit

- There are many copies along the residual gauge orbit
- Is b different for each copy?
- If, how is b distributed?
- How is b correlated with the quantity $\text{Tr}D$?
- How evolves b with volume and discretization?
- Is there a possibility to use b as a gauge parameter?

Distribution of b [3d, beta=3.46, Maas, unpublished]

- At small volumes: Small number of *Gribov copies*

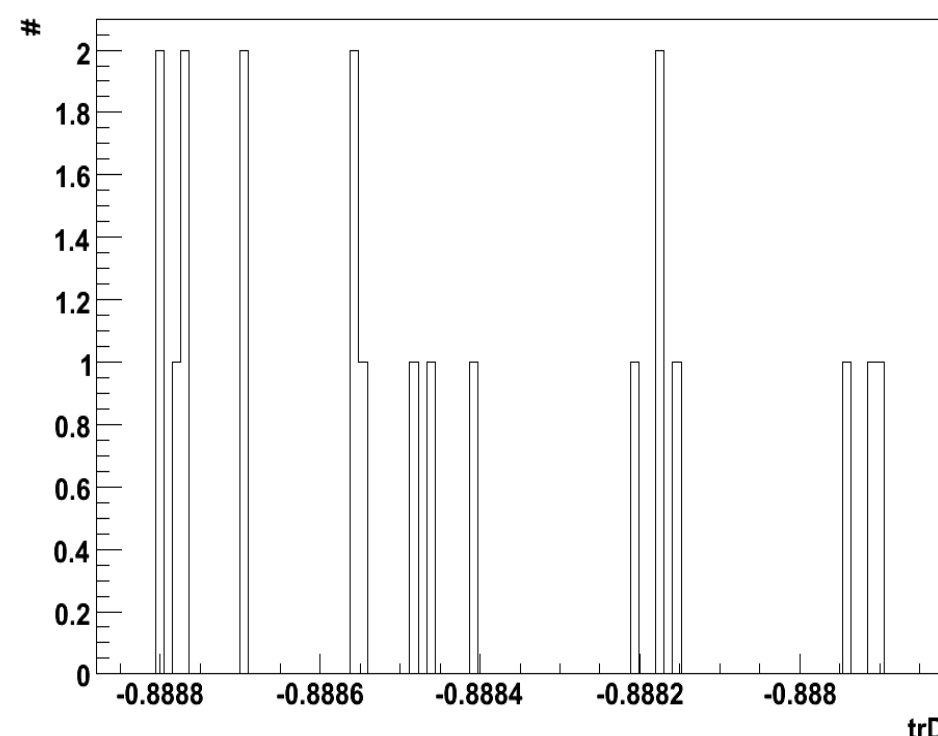
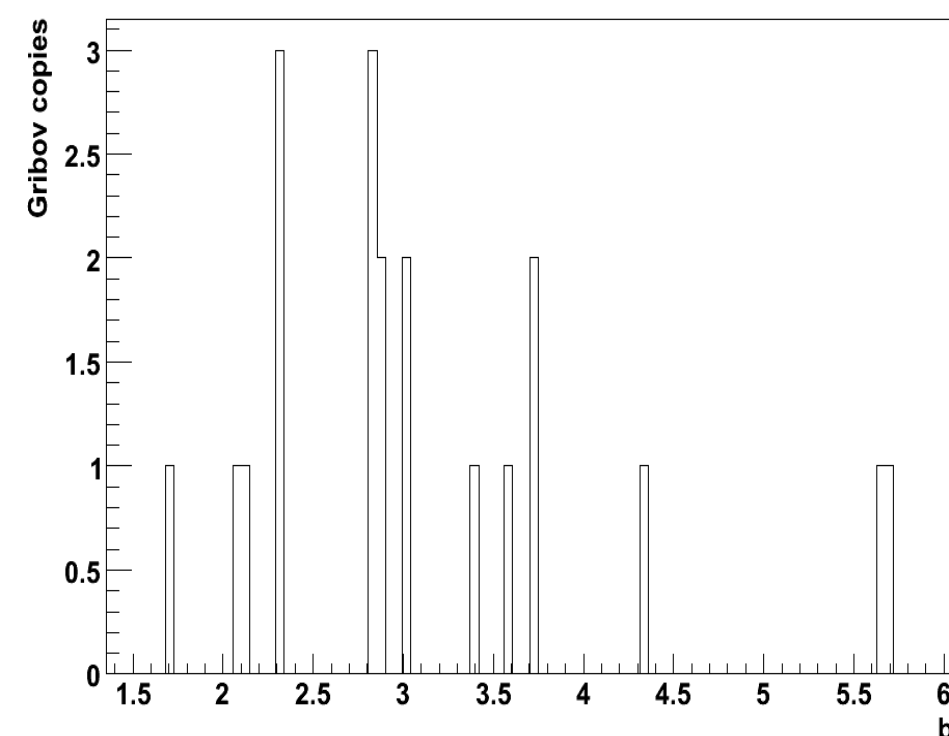
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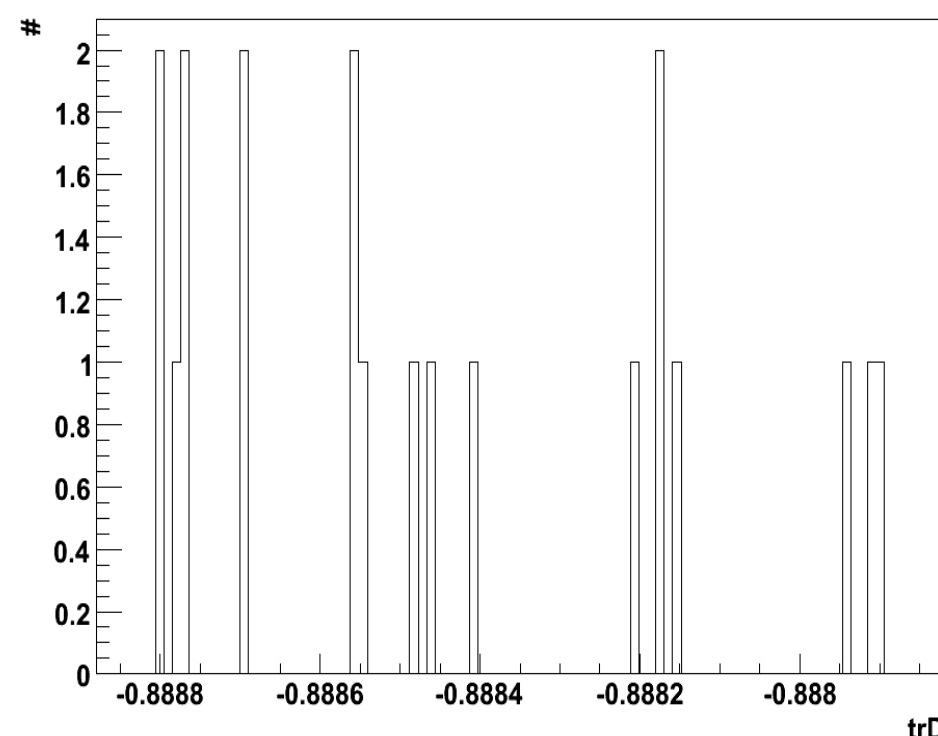
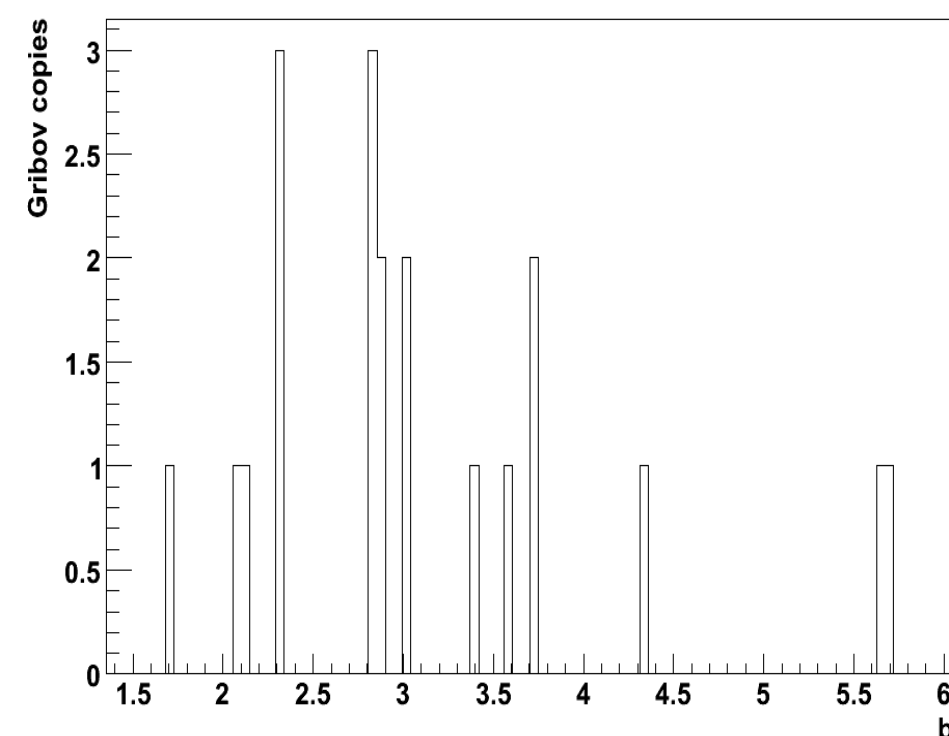
$$b = G(0.395 \text{ GeV}) / G(\infty \text{ GeV}) \text{ for } V = (3.1 \text{ fm})^3$$

$$\text{TrD for } V = (3.1 \text{ fm})^3$$


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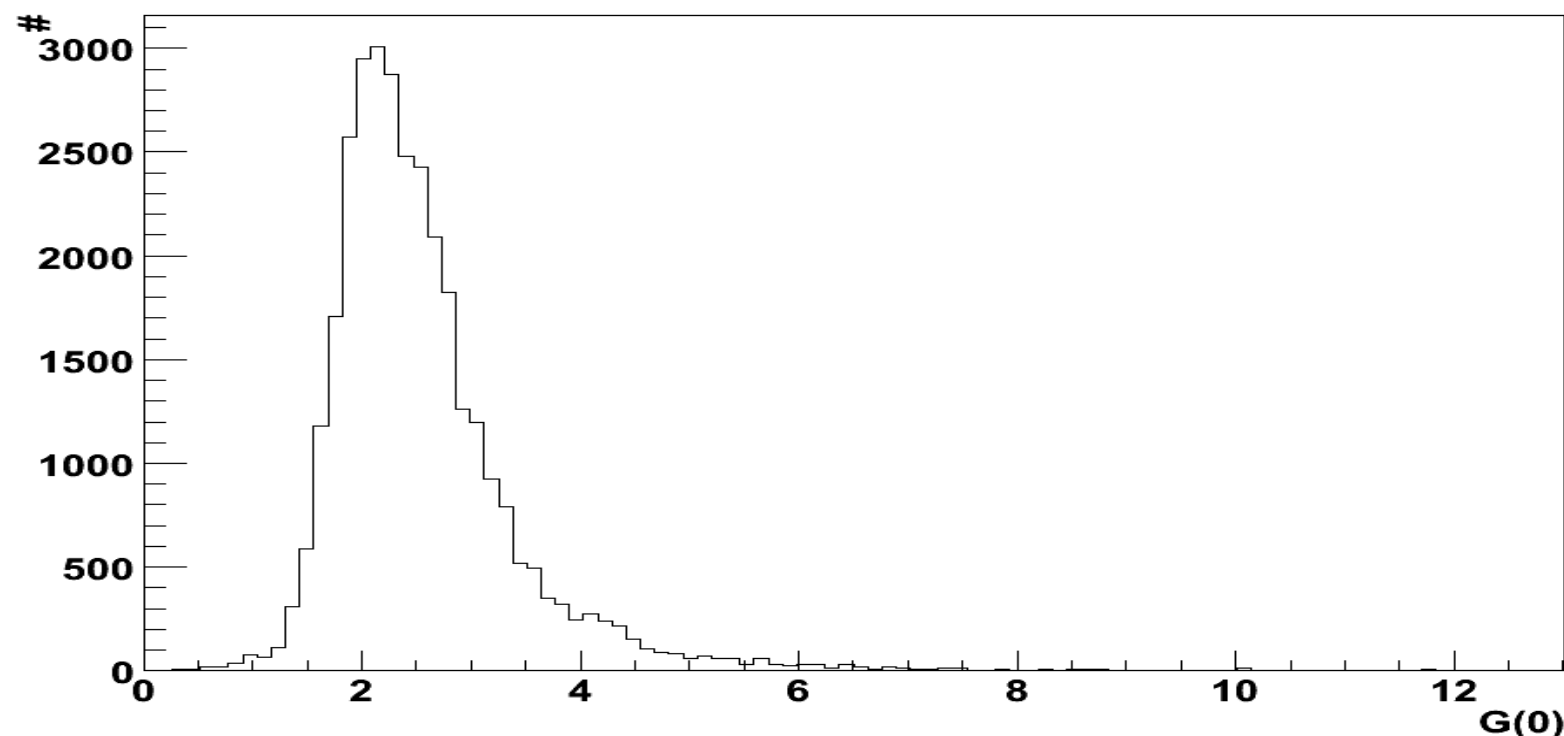
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- At small volumes: Small number of **Gribov copies**
- b significantly different, if **TrD** is significantly different
 - Correspond to different **continuum Gribov copies**
 - b is relatively more different than **TrD**

Distribution of b on average [3d, beta=3.46, Maas, unpublished]

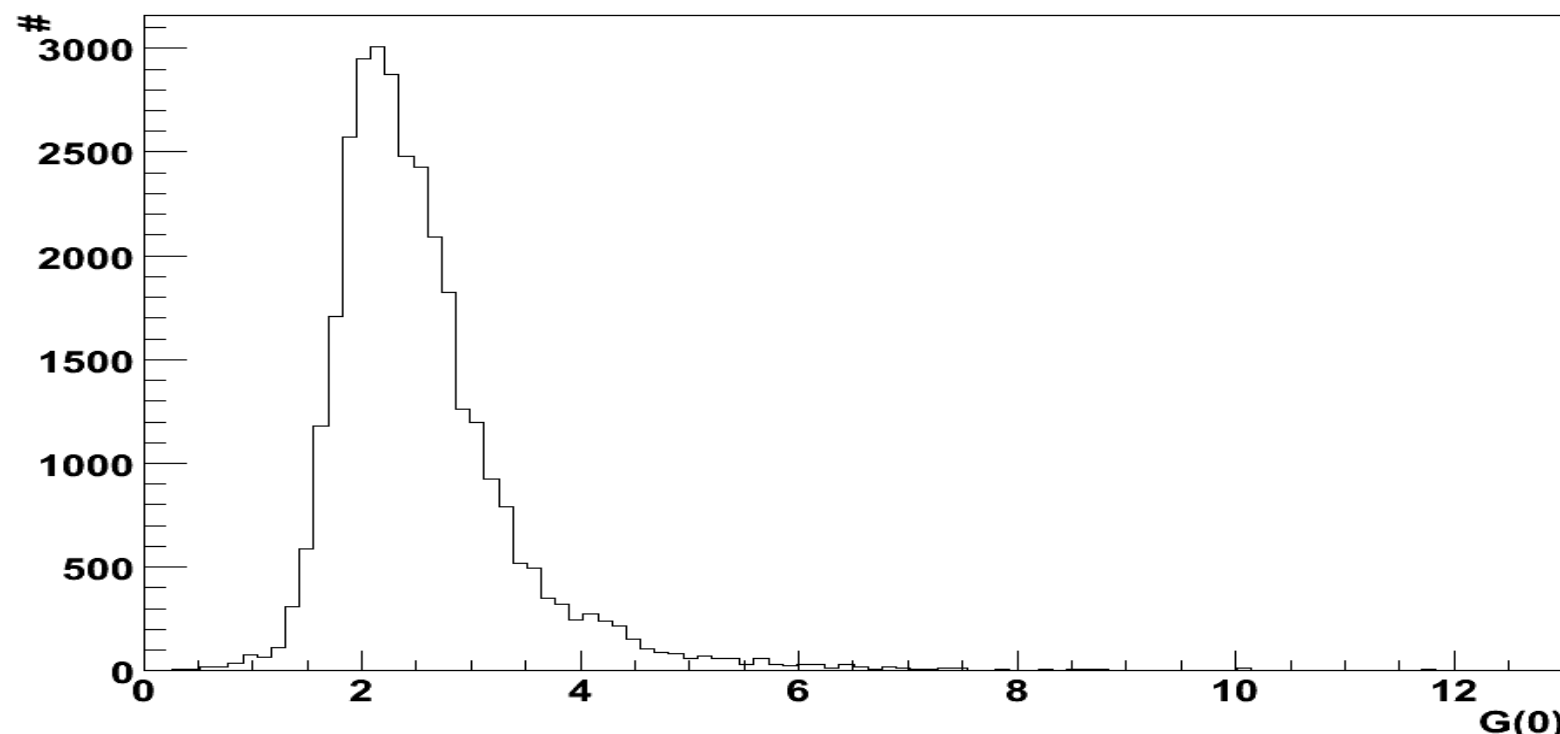
$G(0.395 \text{ GeV})/G(\infty \text{ GeV})$ for $V=(3.1 \text{ fm})^3$ from 1622 configurations



- Distribution is asymmetric with tail to large values
 - Extreme values for the **ghost propagator** are possible

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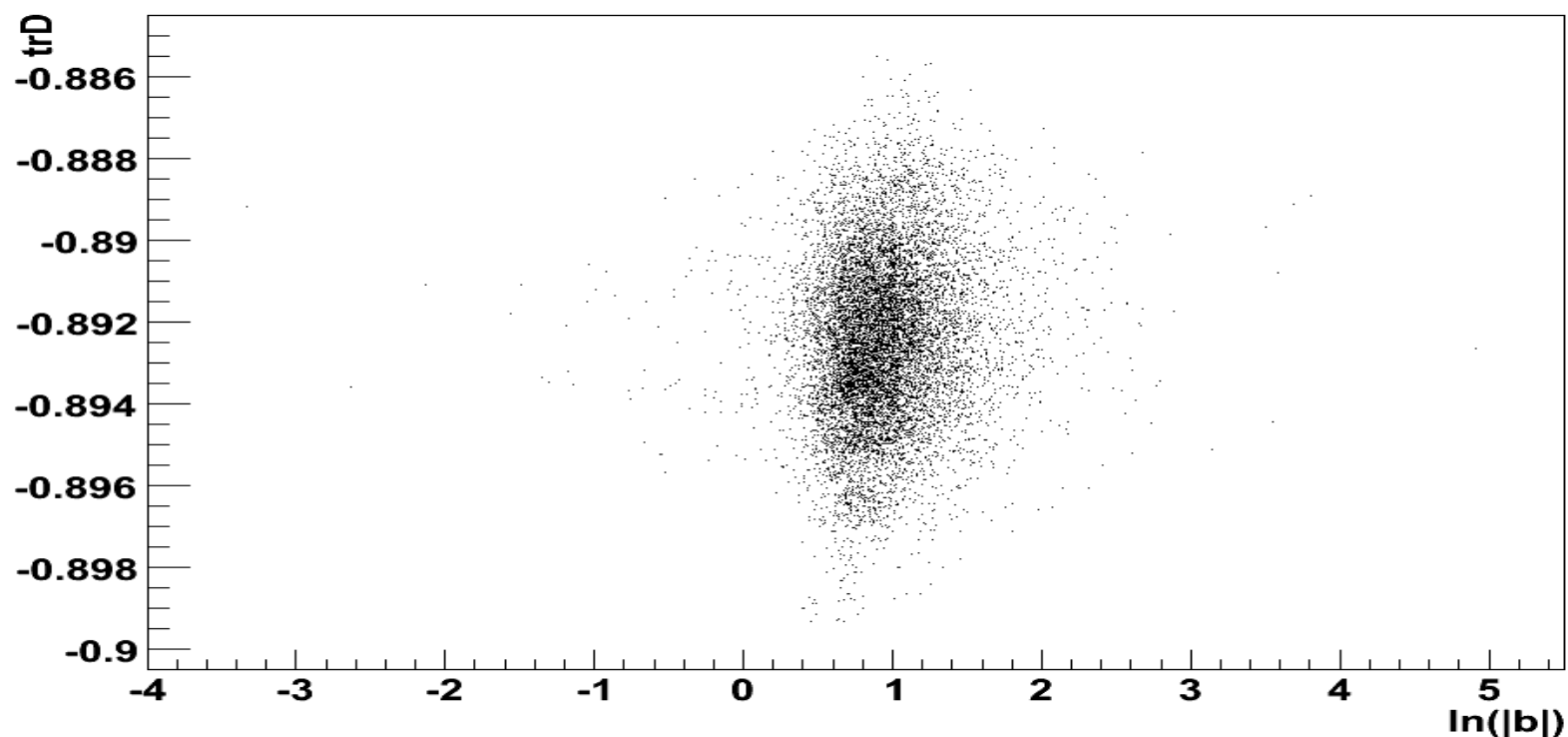
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- Is there a unique correlation of **b** and **$\text{Tr}D$** ?

Distribution of b on average [3d, beta=3.46, Maas, unpublished]

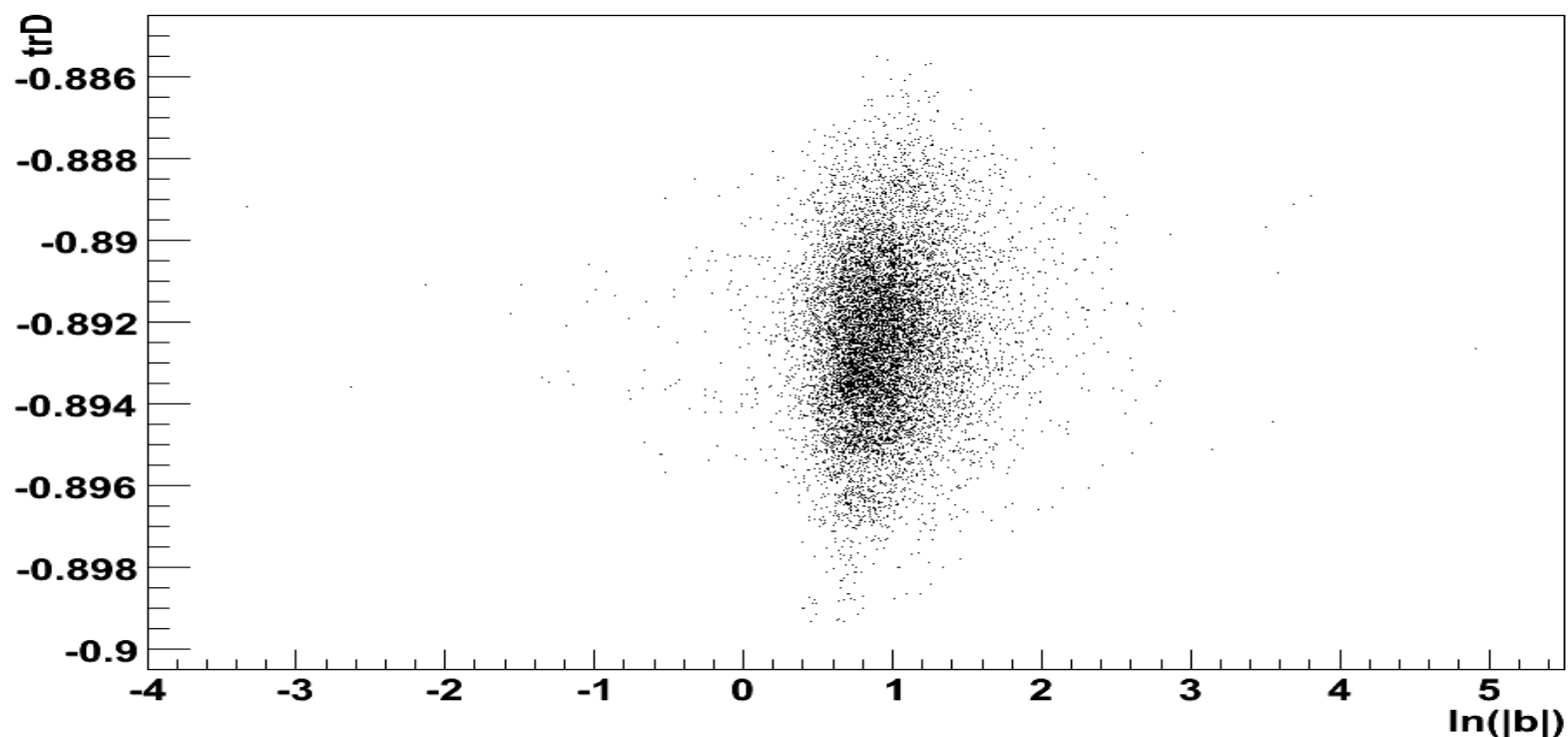
TrD vs. b for $V=(3.1 \text{ fm})^3$



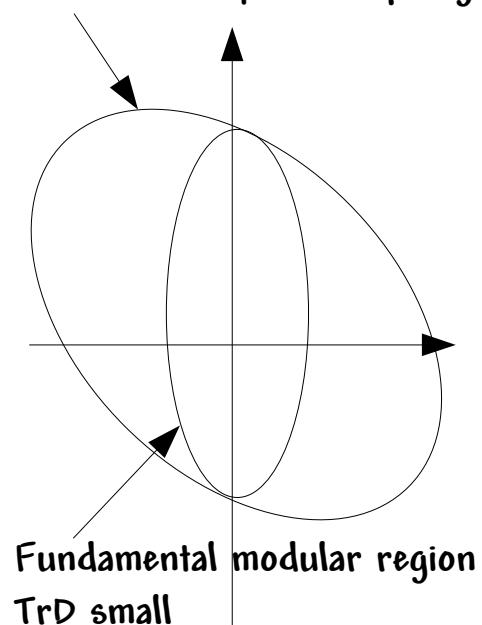
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Distribution of b on average [3d, beta=3.46, Maas, unpublished]

TrD vs. b for $V=(3.1 \text{ fm})^3$



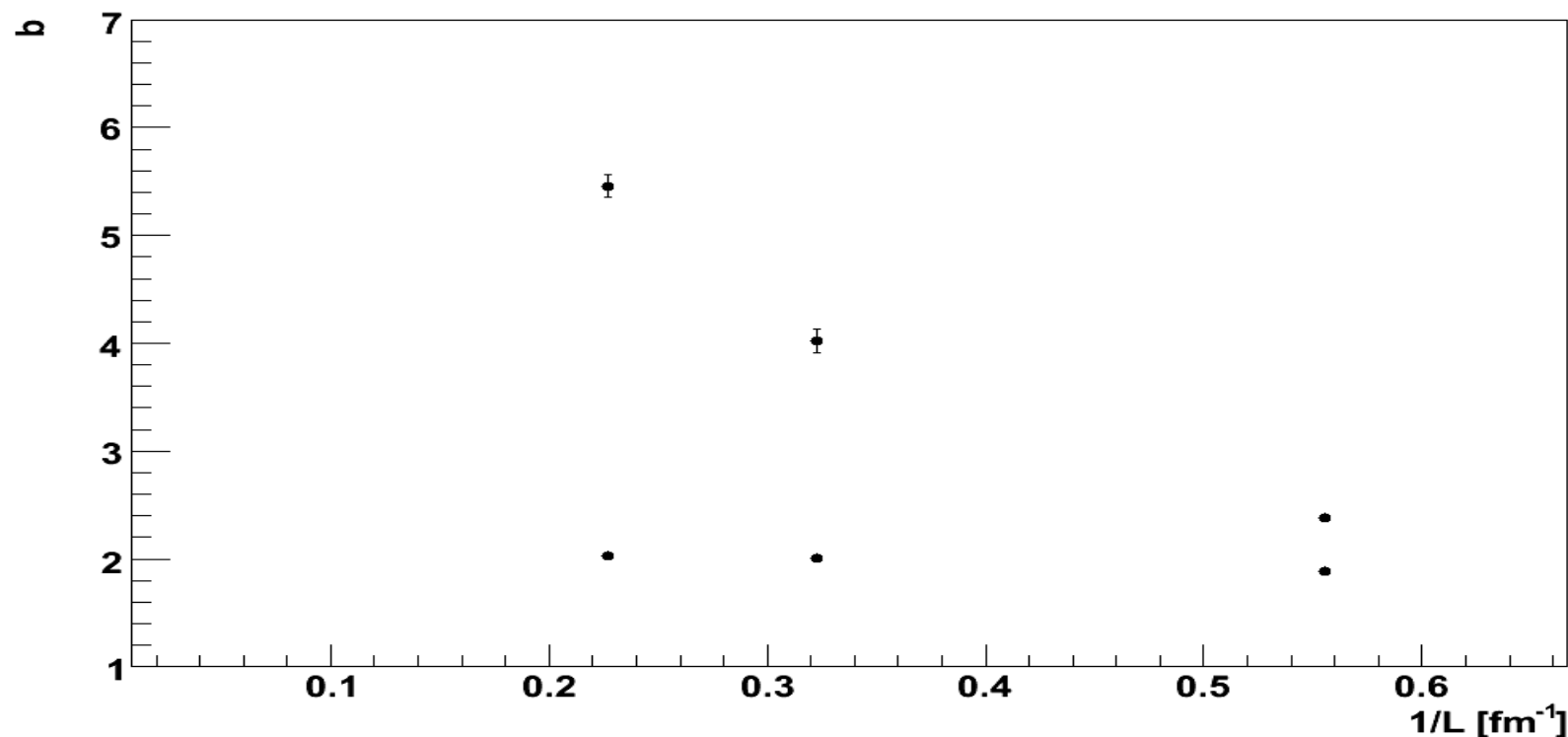
Gribov horizon: b potentially large



- No obvious correlation between both quantities
 - Only a small tendency that small b (far away from the Gribov horizon) correlate with small TrD (fundamental modular region)

The b gauge corridor [3d, beta=3.45, Maas, unpublished]

b corridor



- The **corridor** [min b , max b] opens with volume
 - Overall scale partly a **finite-volume effect**
 - **max b /min b** not: Possible range of b values increases
 - Can **max b** diverge in the thermodynamic limit?

Constructing a gauge parameter from b

- Volumes increases: b distribution becomes dense

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 - Select from the range $[\min b, \max b]$ a value B (for a given volume)
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 - Average over these copies to obtain the gauge-fixed correlation functions in the Landau- B gauge

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 - Select from the range $[\min b, \max b]$ a value B (for a given volume)
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 - Average over these copies to obtain the gauge-fixed correlation functions in the Landau- B gauge
 - In addition: $\min B$ and $\max B$ gauges, selecting the extremal values of b on each residual gauge orbit

Compare correlation functions between Landau gauges

- Landau- \mathcal{B} gauges
 - Here: $\min\mathcal{B}$ und $\max\mathcal{B}$ -gauges

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Compare correlation functions between Landau gauges

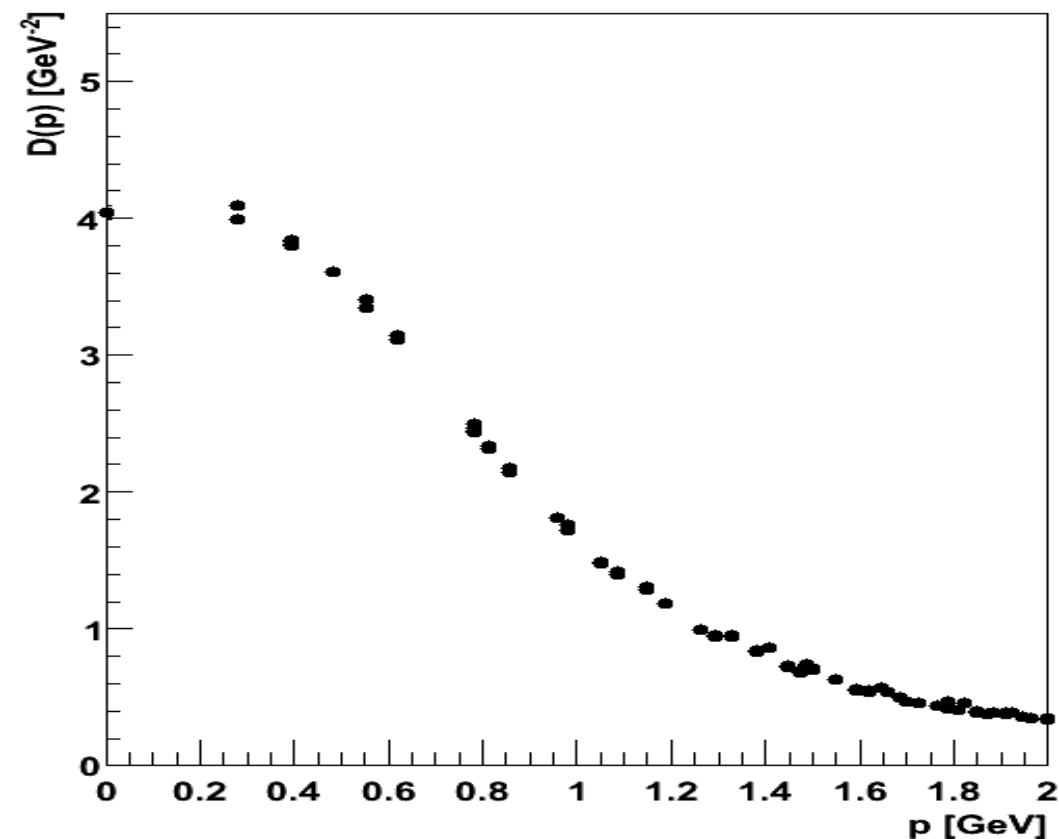
- Landau- \mathcal{B} gauges
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Compare correlation functions between Landau gauges

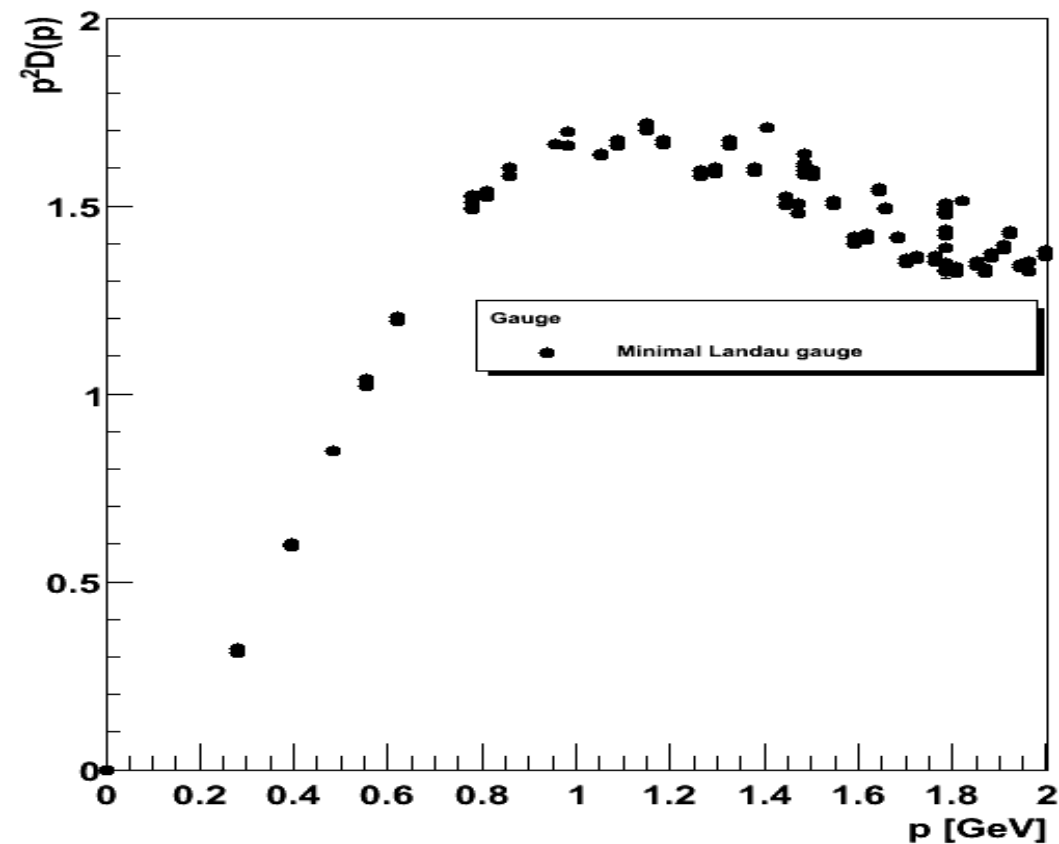
- Landau- \mathcal{B} gauges
 - Here: $\min\mathcal{B}$ und $\max\mathcal{B}$ -gauges
- Absolute Landau gauge
- Inverse Landau gauge
 - Maximize $\text{Tr}D$
- Minimal Landau gauge
 - Select a random copy on each residual gauge orbit
 - Probability distribution is determined by the **underlying algorithm**
 - **Differ for different algorithms**, but cheap compared to all other gauges
 - No direct analogue in or translation to continuum field theory
 - Current standard in **lattice calculations**

Gluon propagator in the minimal Landau gauge [3d, Maas, unpublished]

Gluon propagator



Gluon dressing function



- Mild infrared suppression in 3d (strong in 2d, none in 4d)

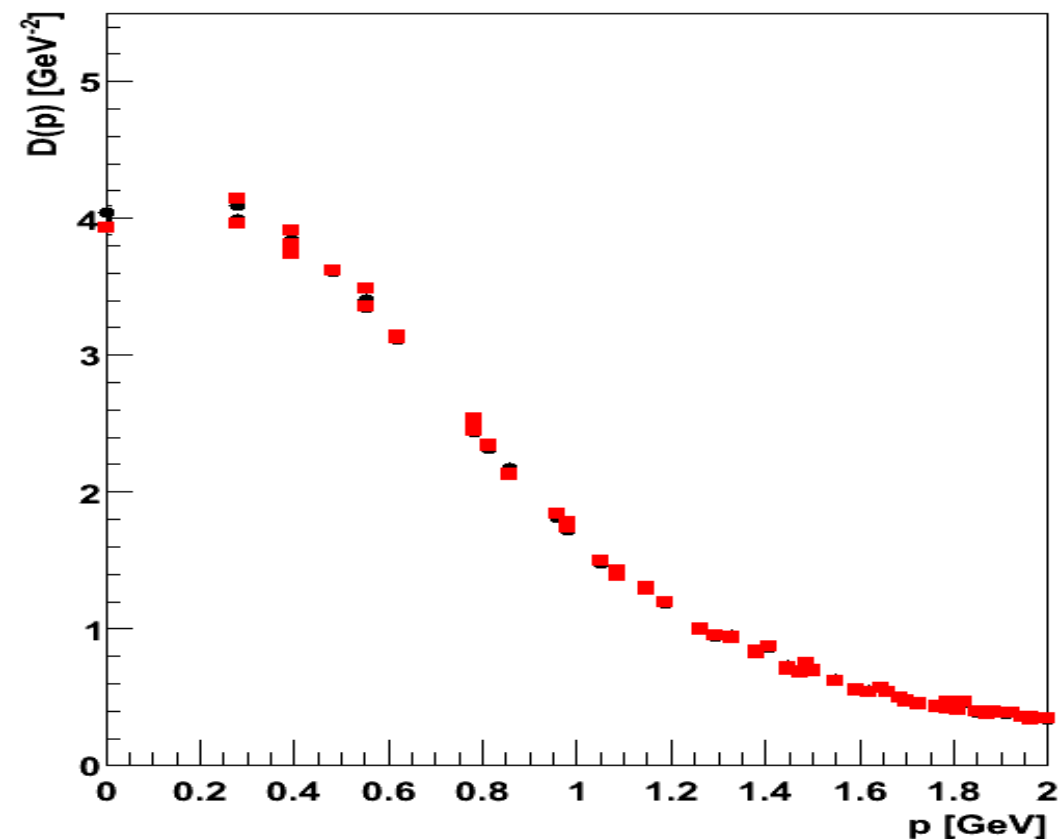
- Infrared constant in the infinite-volume/continuum limit

[Cucchieri et al. 2007/8, von Smekal et al. 2008]

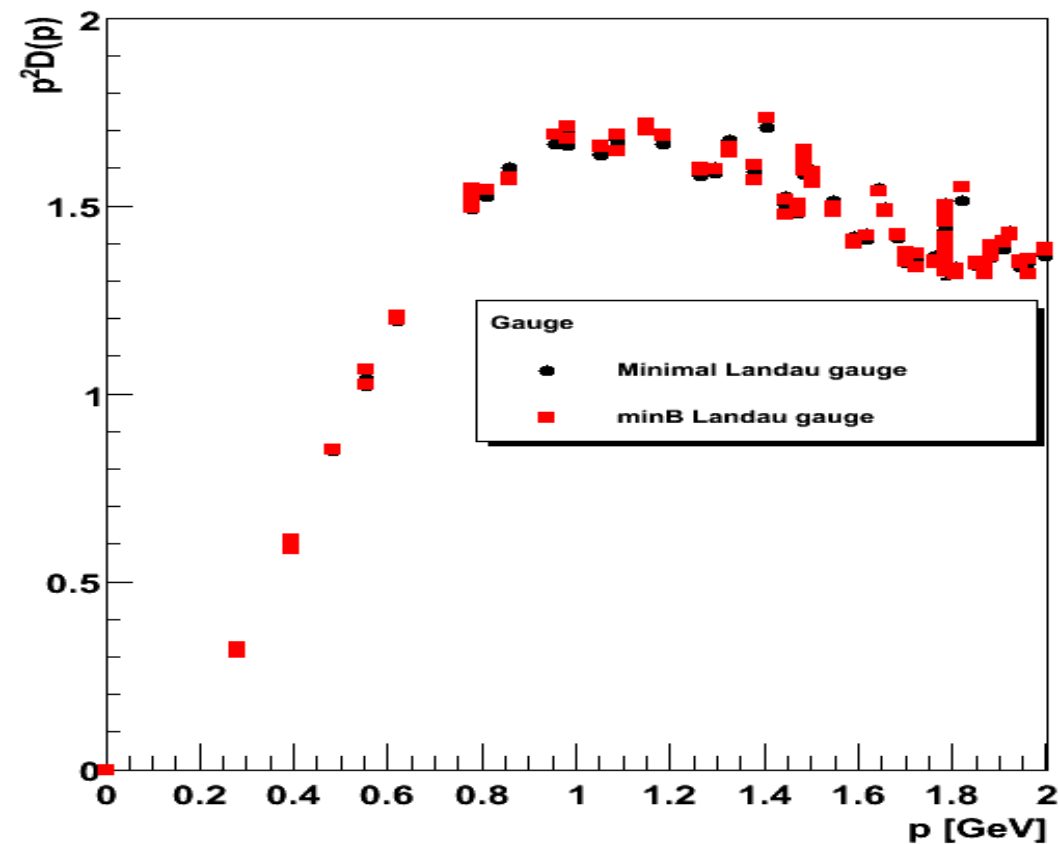
- Possibly vanishing in 2d [Maas, 2007, Cucchieri 2007/8]

Gluon propagator in the minB Landau gauge [3d, Maas, unpublished]

Gluon propagator



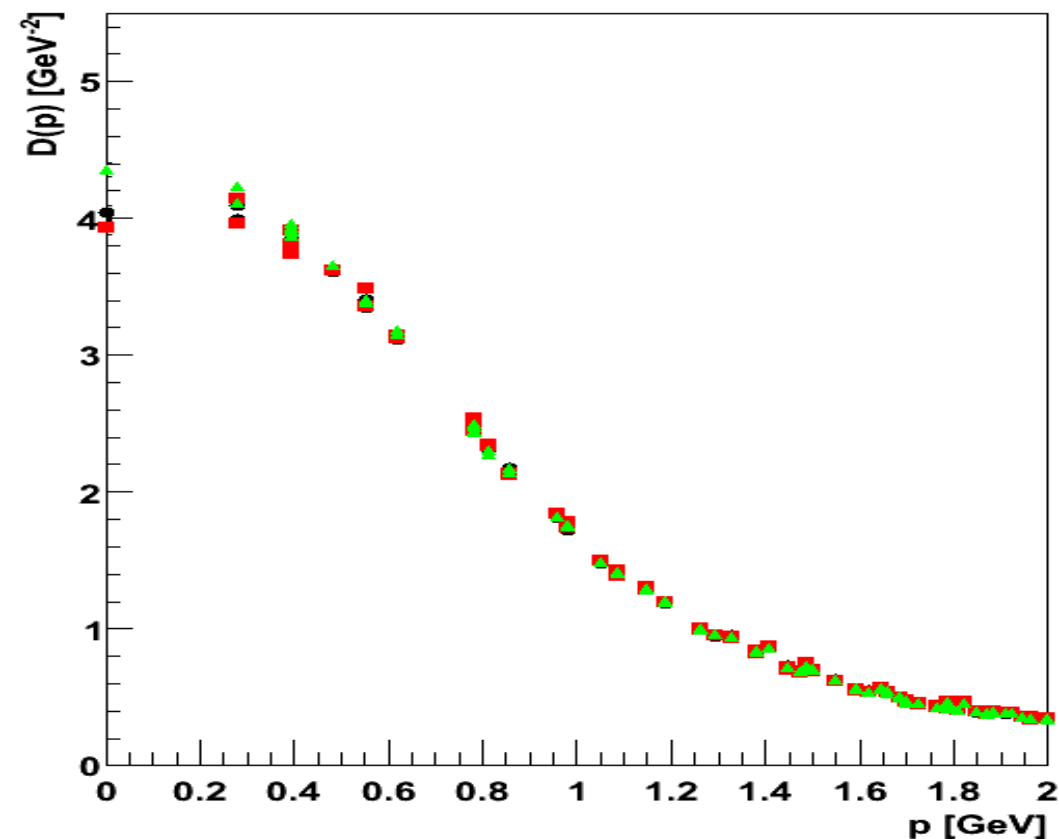
Gluon dressing function



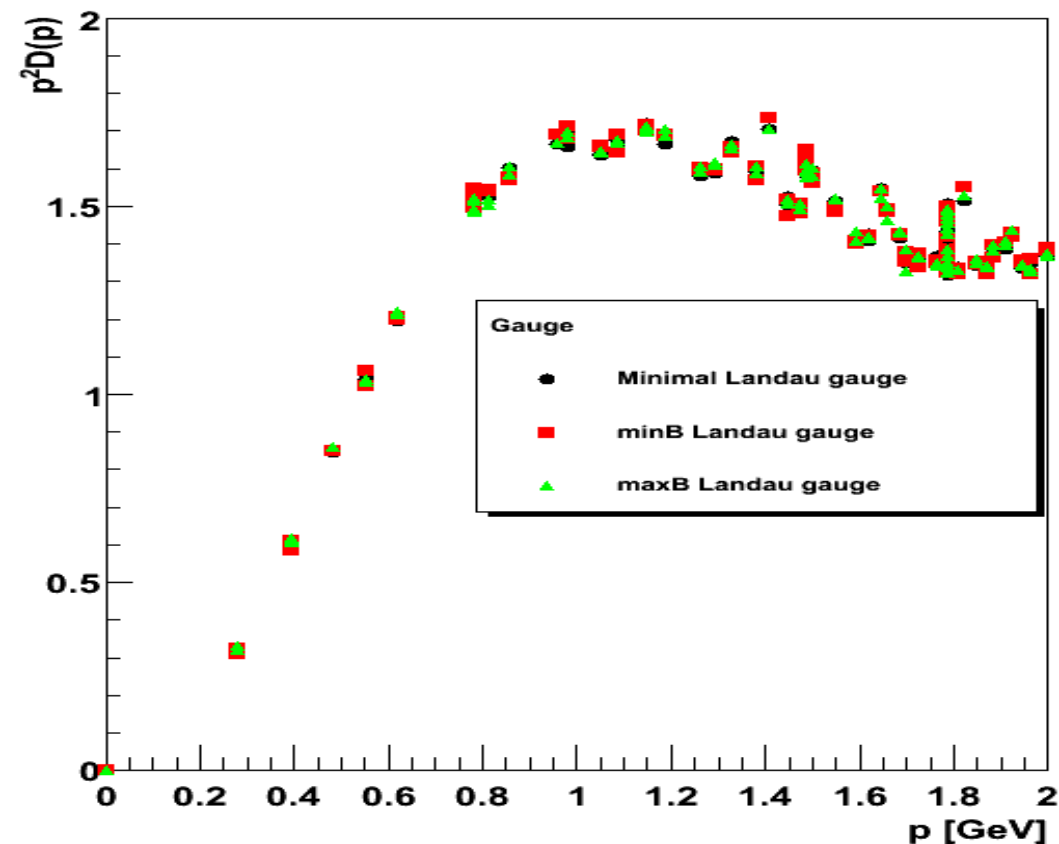
- At current volumes little difference to minimal Landau gauge

Gluon propagator in the maxB Landau gauge [3d, Maas, unpublished]

Gluon propagator



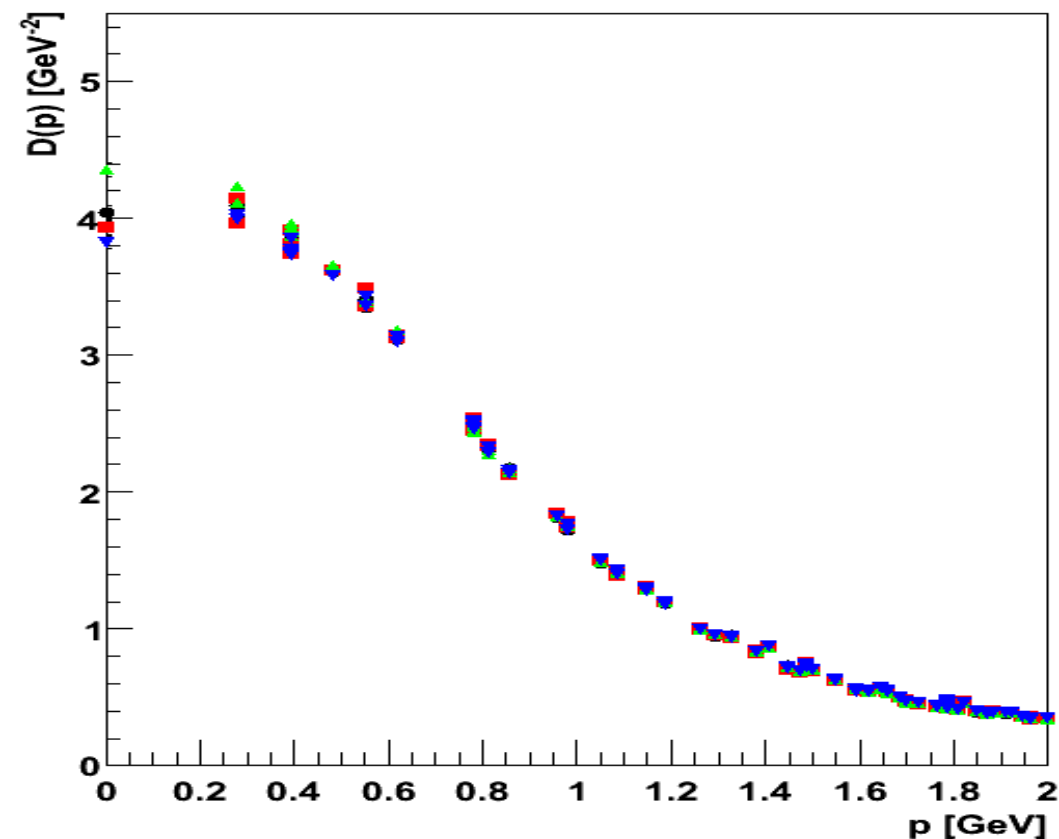
Gluon dressing function



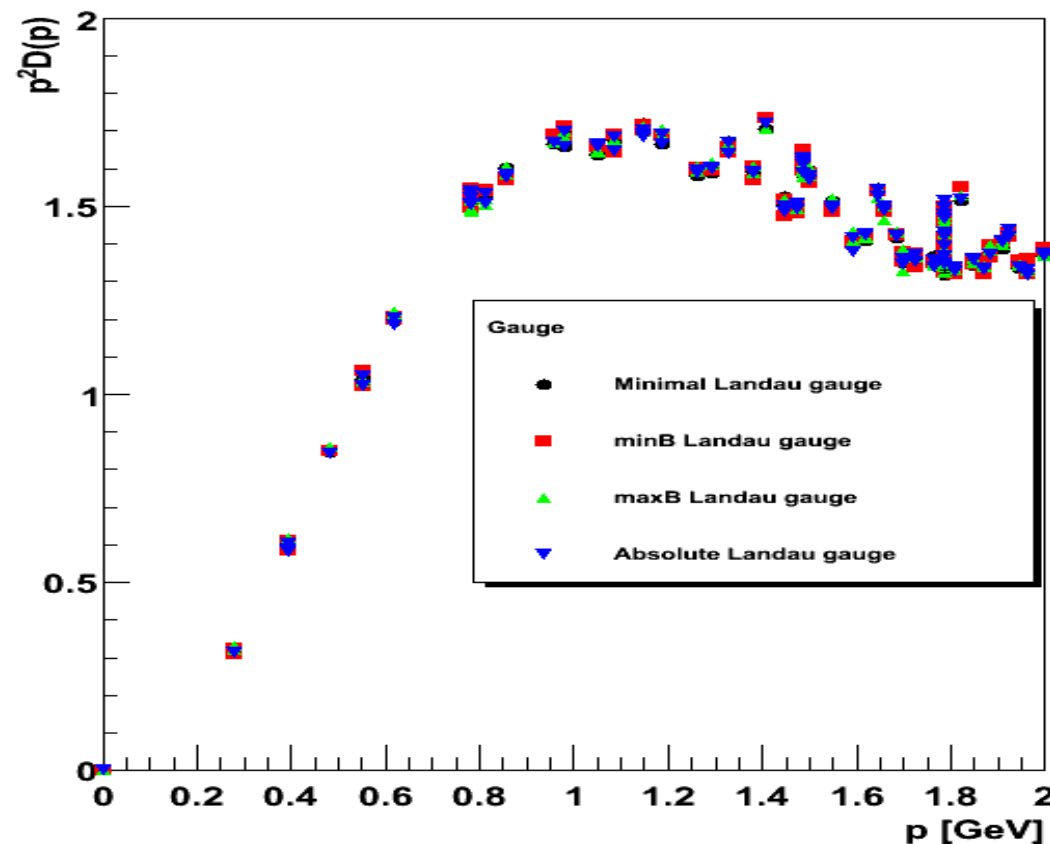
- At current volumes little difference to minimal Landau gauge
 - A little less infrared suppressed

Gluon propagator in the absolute Landau gauge [3d, Maas, unpublished]

Gluon propagator



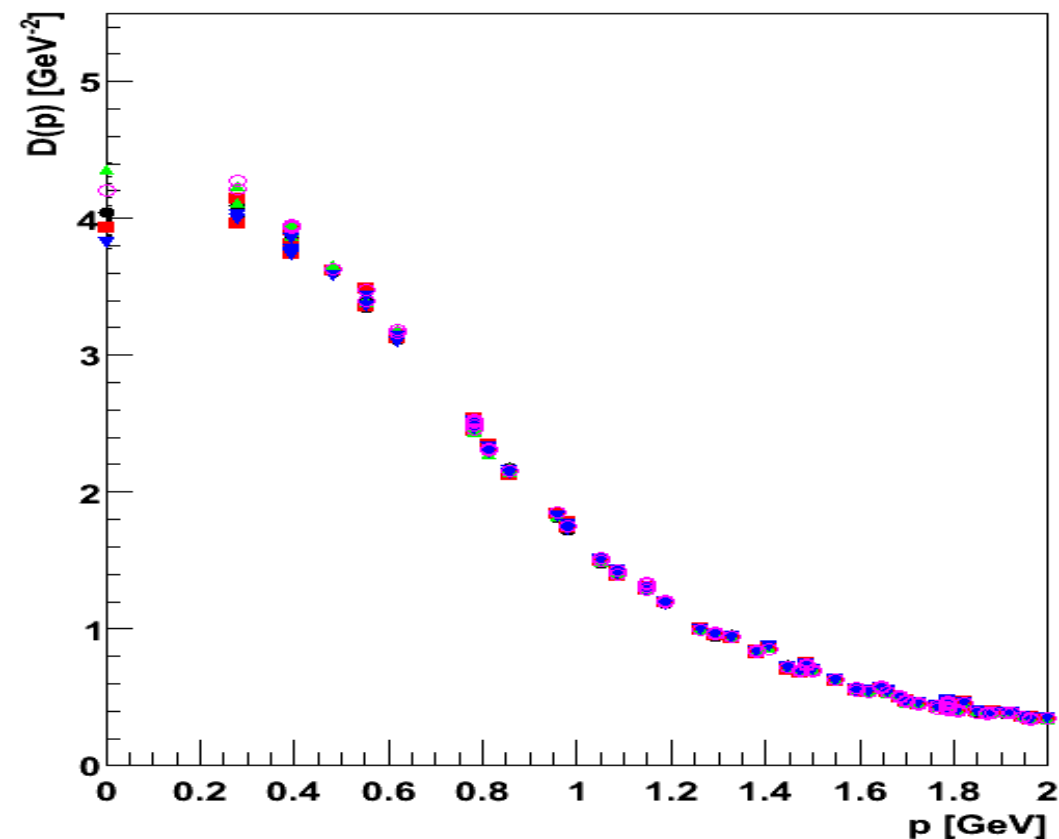
Gluon dressing function



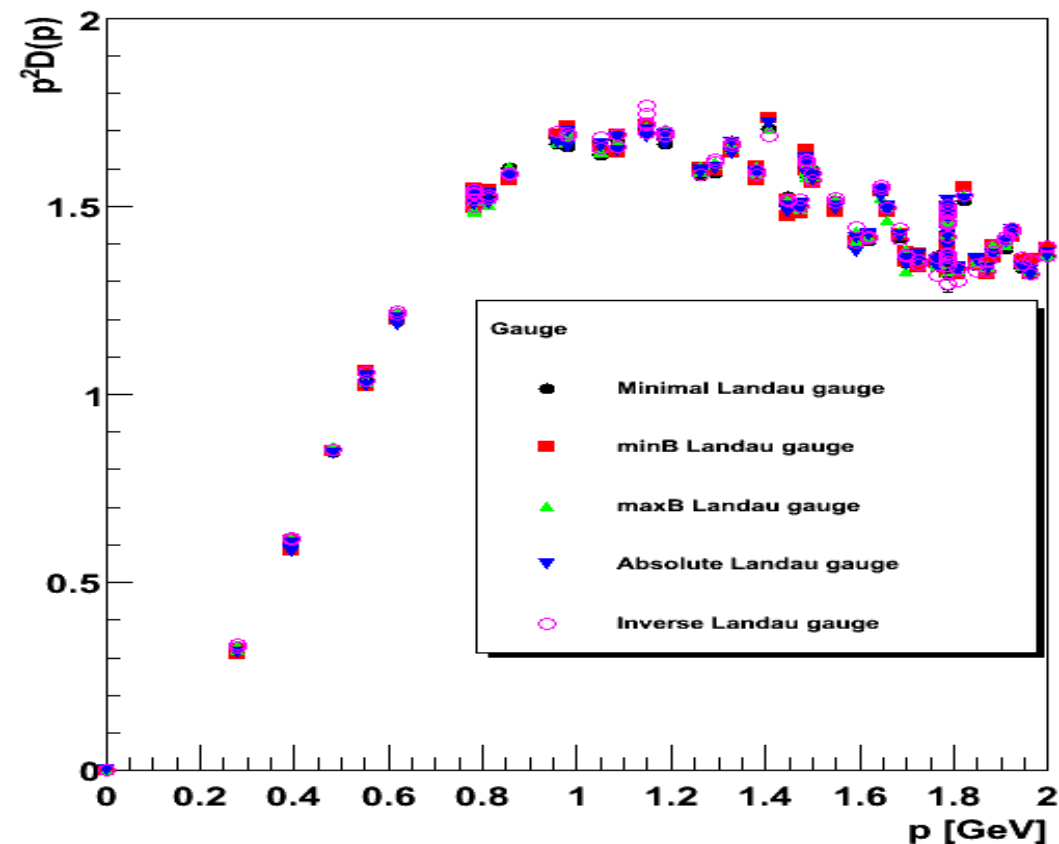
- At current volumes little difference to minimal Landau gauge
 - Somewhat stronger infrared suppressed
 - Unclear whether it vanishes

Gluon propagator in the inverse Landau gauge [3d, Maas, unpublished]

Gluon propagator



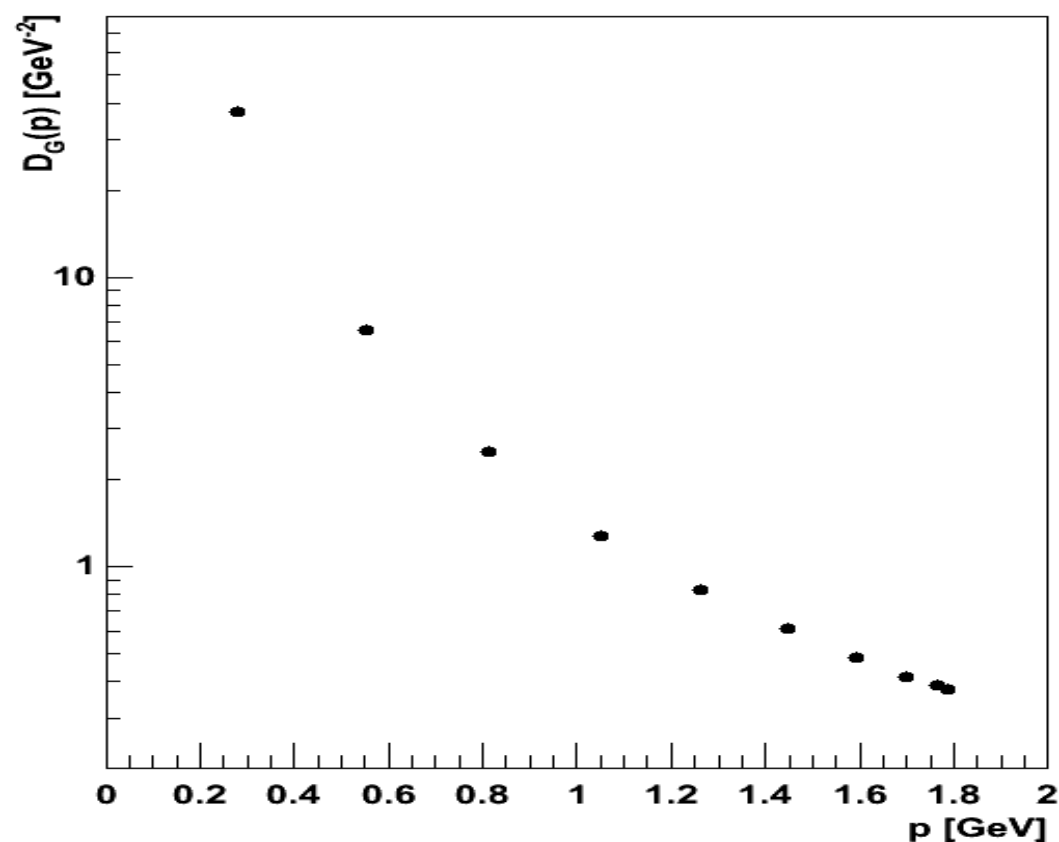
Gluon dressing function



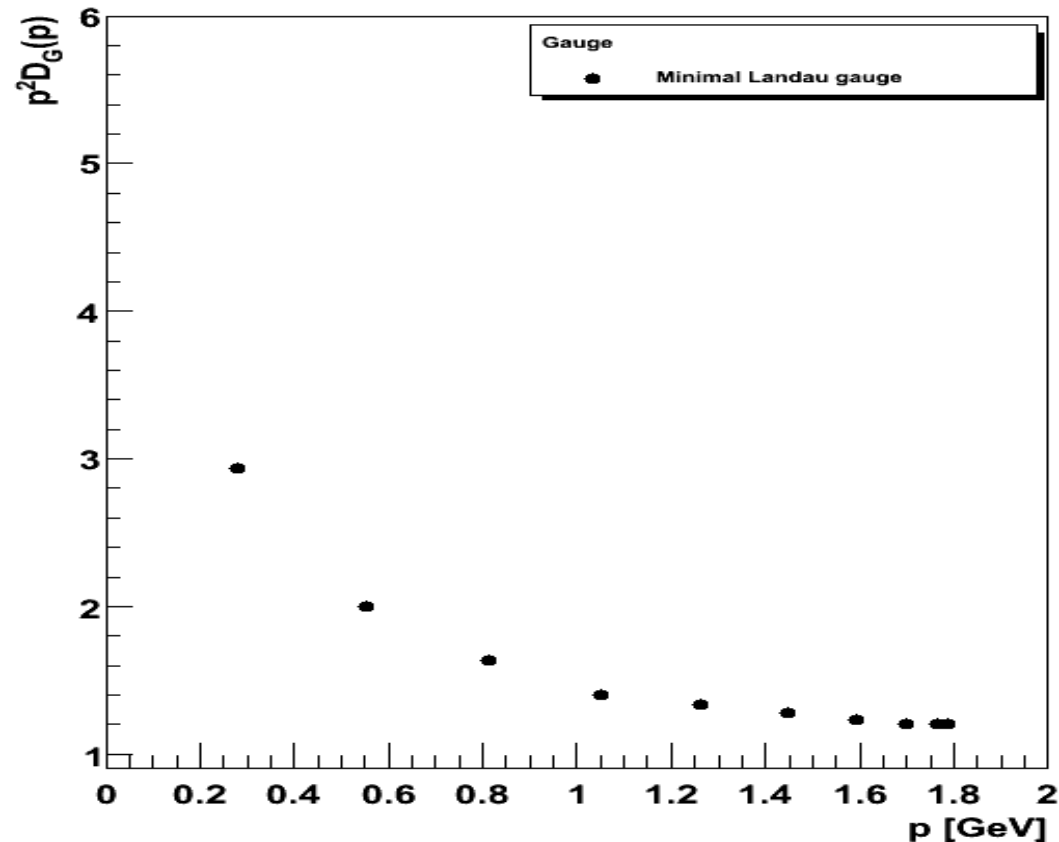
- At current volumes little difference to minimal Landau gauge

Ghost propagator in the minimal Landau gauge [3d, Maas, unpublished]

Ghost propagator



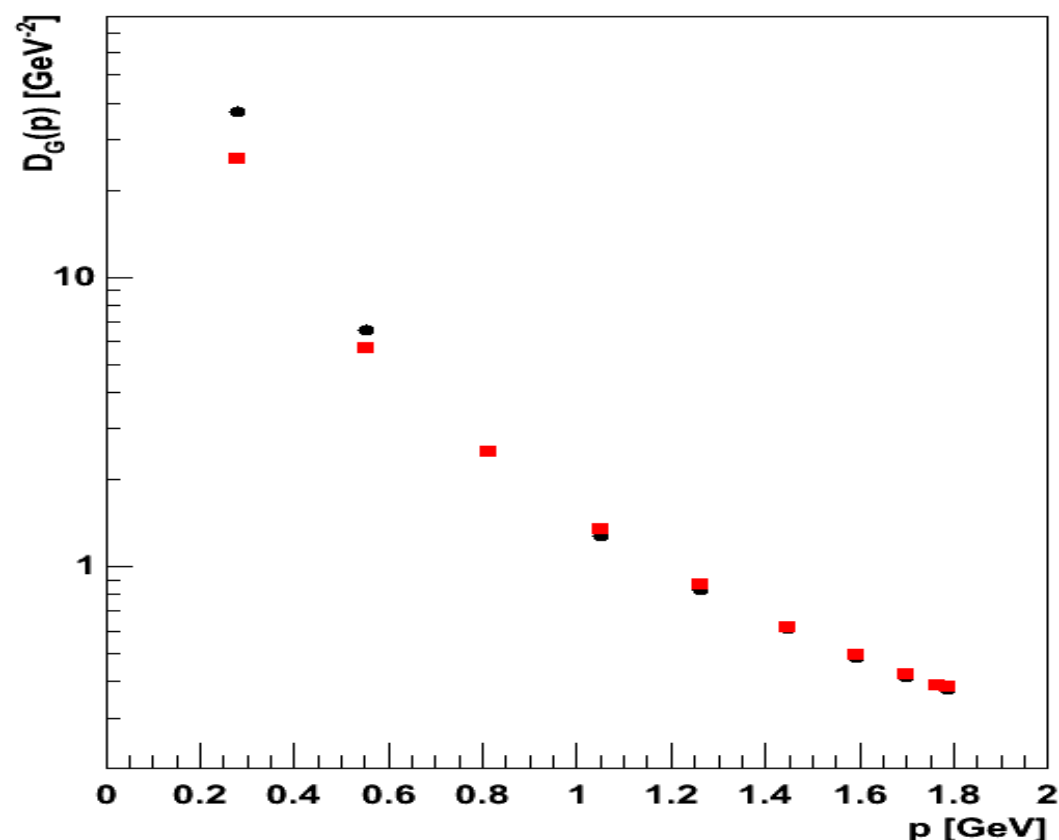
Ghost dressing function



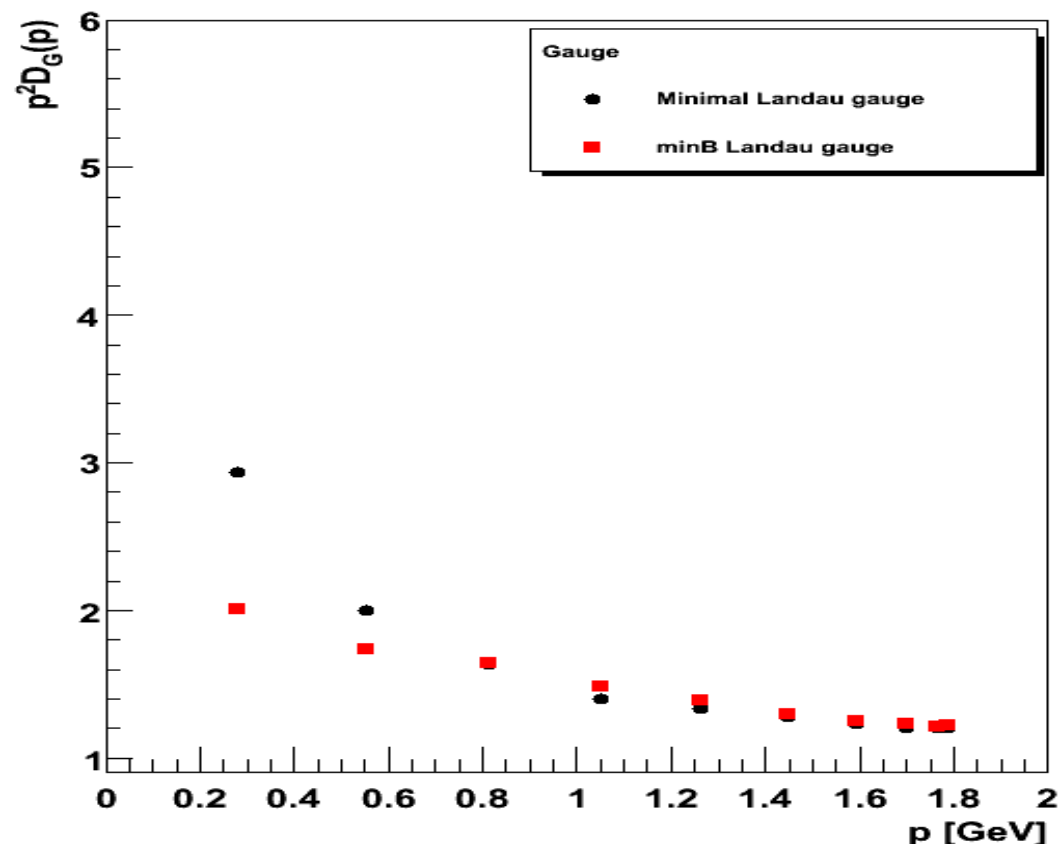
- Infrared enhanced (in all dimensions)
 - In 3d and 4d infrared finite dressing function [Cucchieri et al. 2007/8, von Smekal 2007]
 - In 2d possibly diverging [Maas 2007, Cucchieri et al. 2007/8]

Ghost propagator in the minB Landau gauge [3d, Maas, unpublished]

Ghost propagator



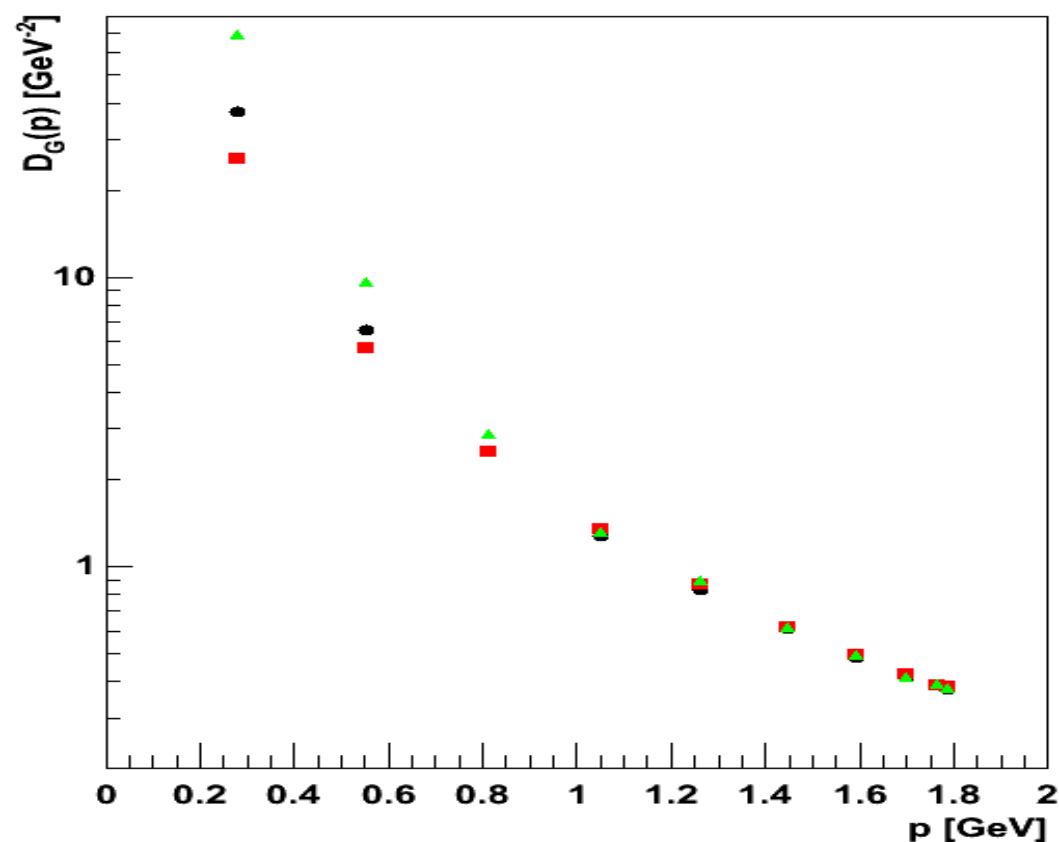
Ghost dressing function



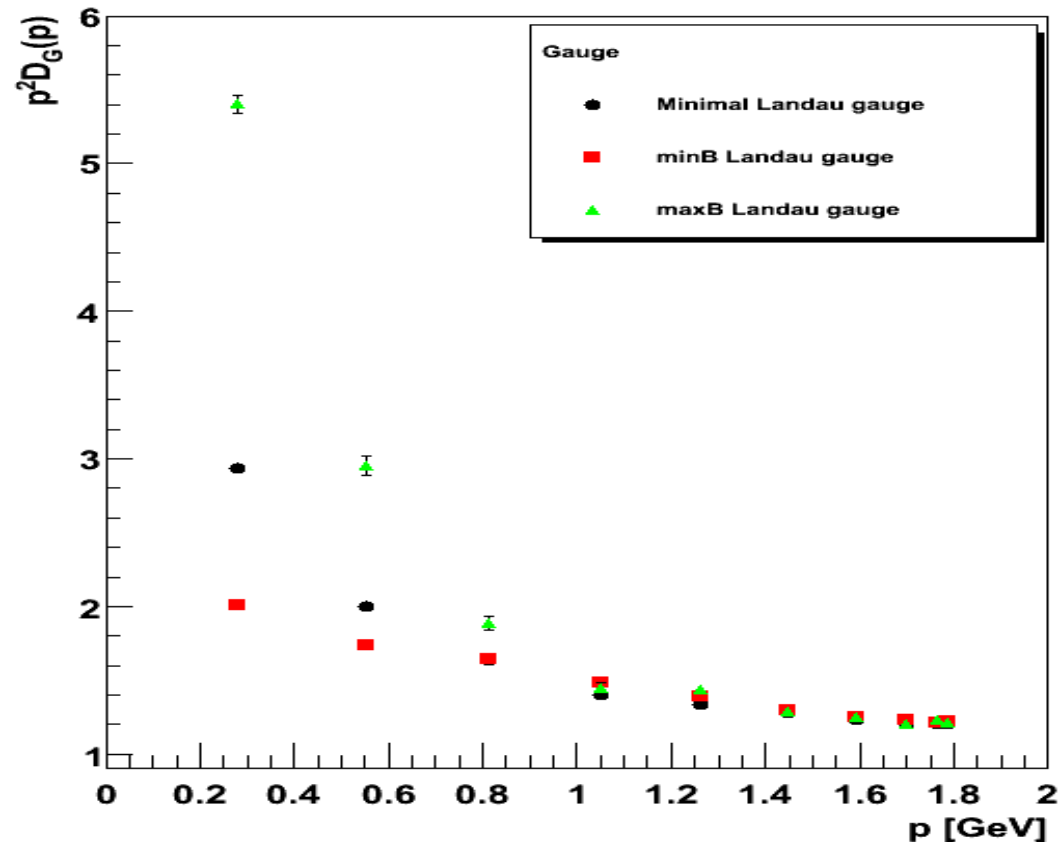
- Significantly less enhanced
 - Will have in 3d and 4d infrared finite dressing function
 - Fate in 2d yet unclear

Ghost propagator in the minimal Landau gauge [3d, Maas, unpublished]

Ghost propagator



Ghost dressing function

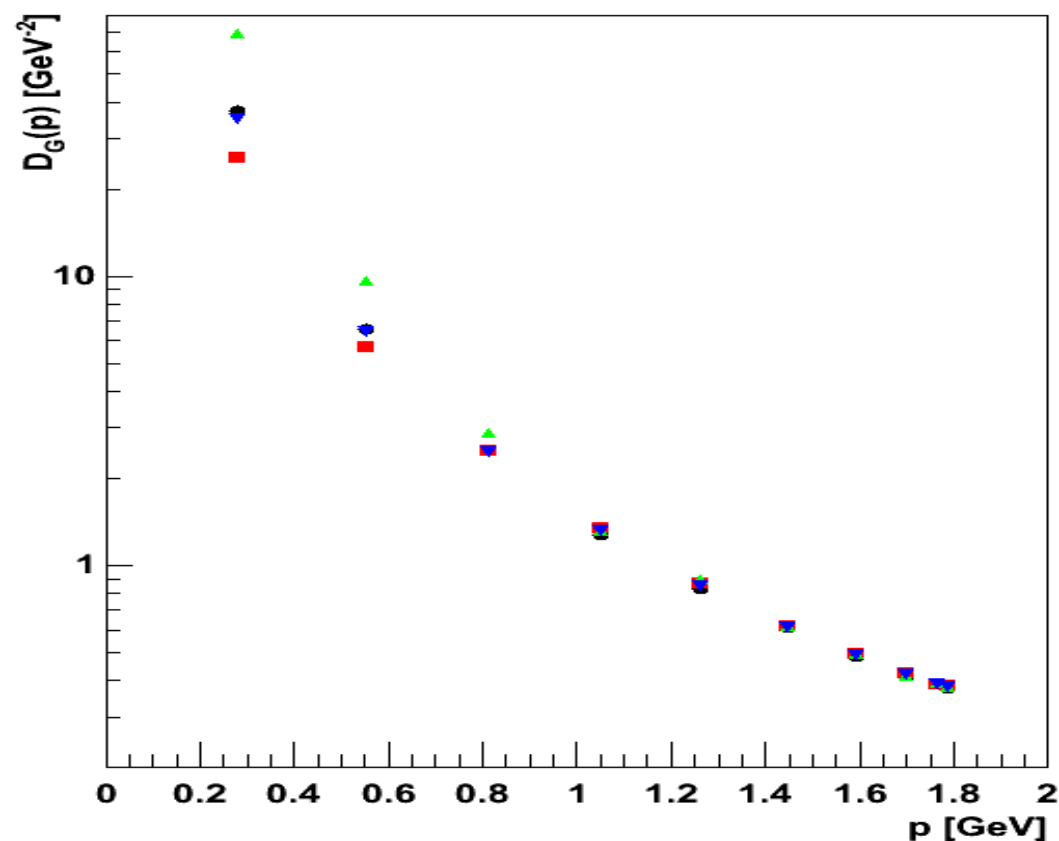


- Strongly enhanced

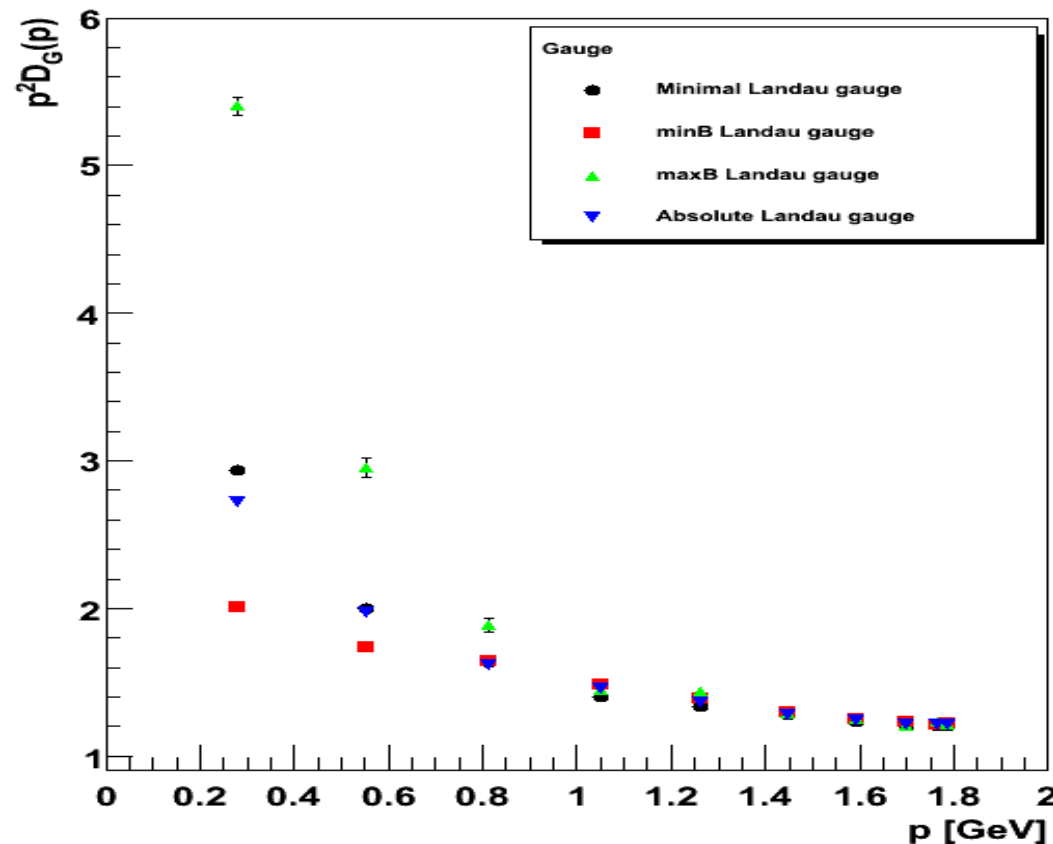
- Could have infrared divergent dressing function in all d
- Effects up to almost 1 GeV

Ghost propagator in the minimal Landau gauge [3d, Maas, unpublished]

Ghost propagator



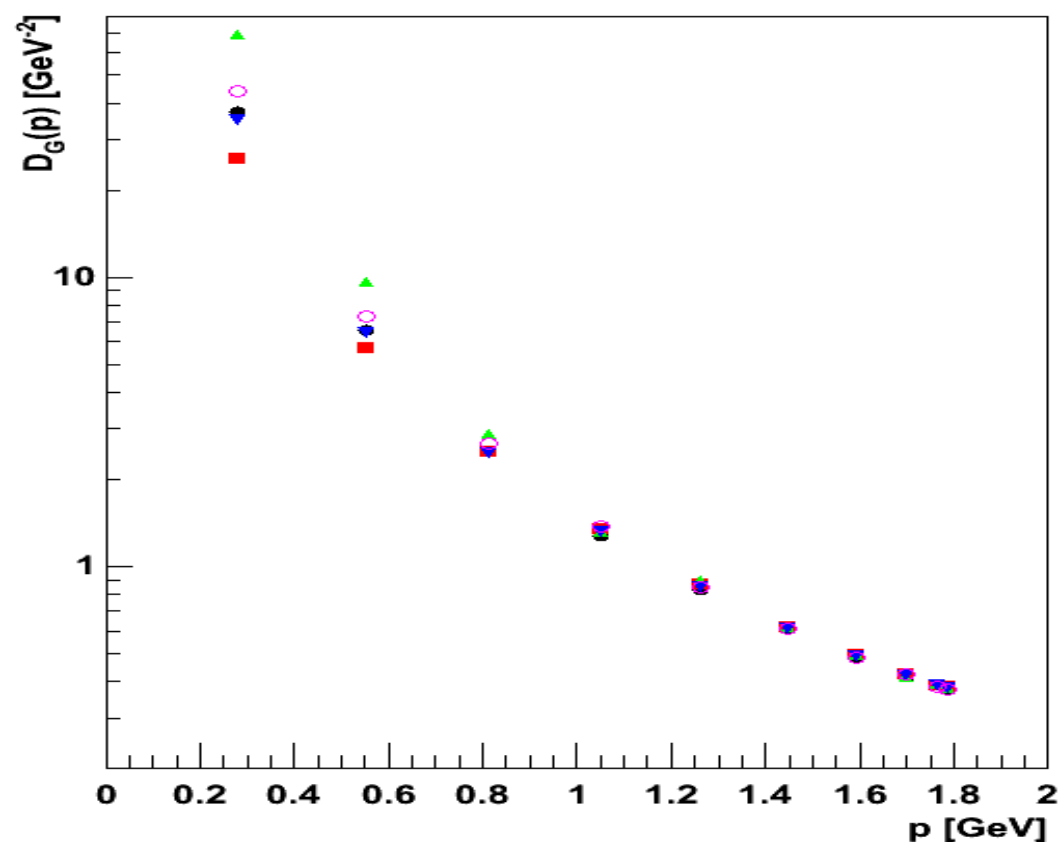
Ghost dressing function



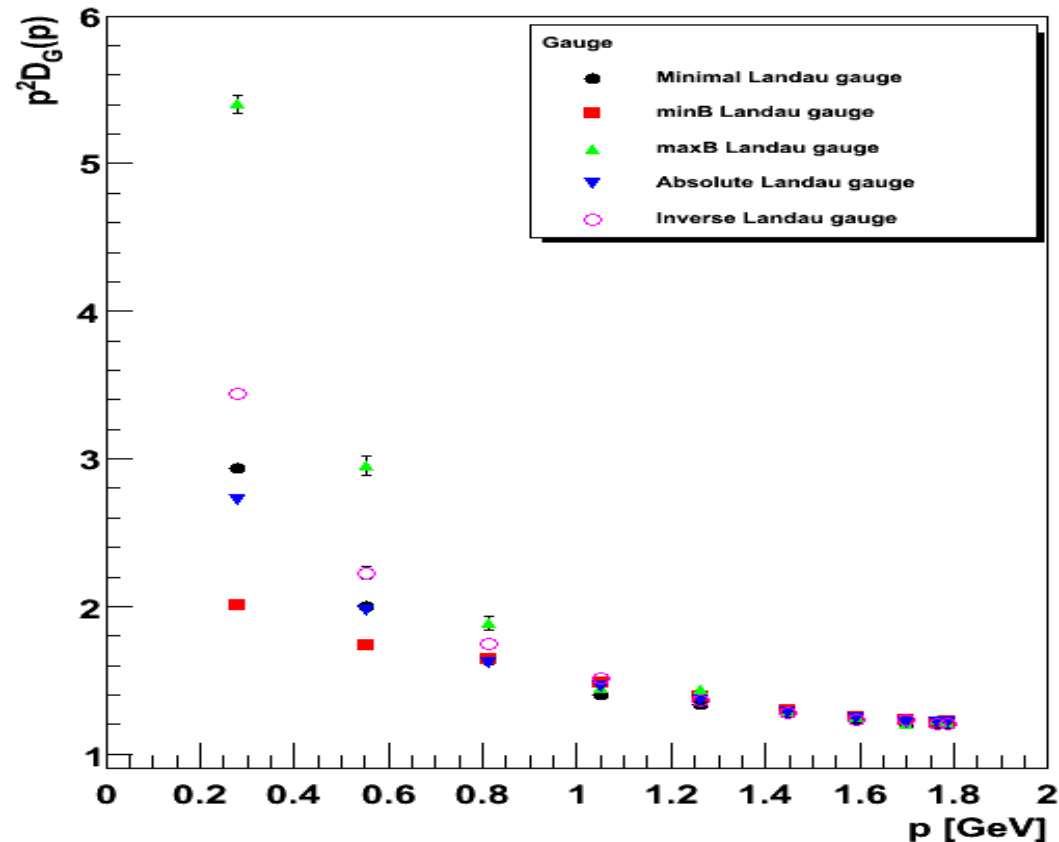
- Less strongly enhanced than minimal Landau gauge
 - But has different finite volume evolution [Maas, 2008]

Ghost propagator in the minimal Landau gauge [3d, Maas, unpublished]

Ghost propagator



Ghost dressing function



- Mildly stronger enhanced than minimal Landau gauge
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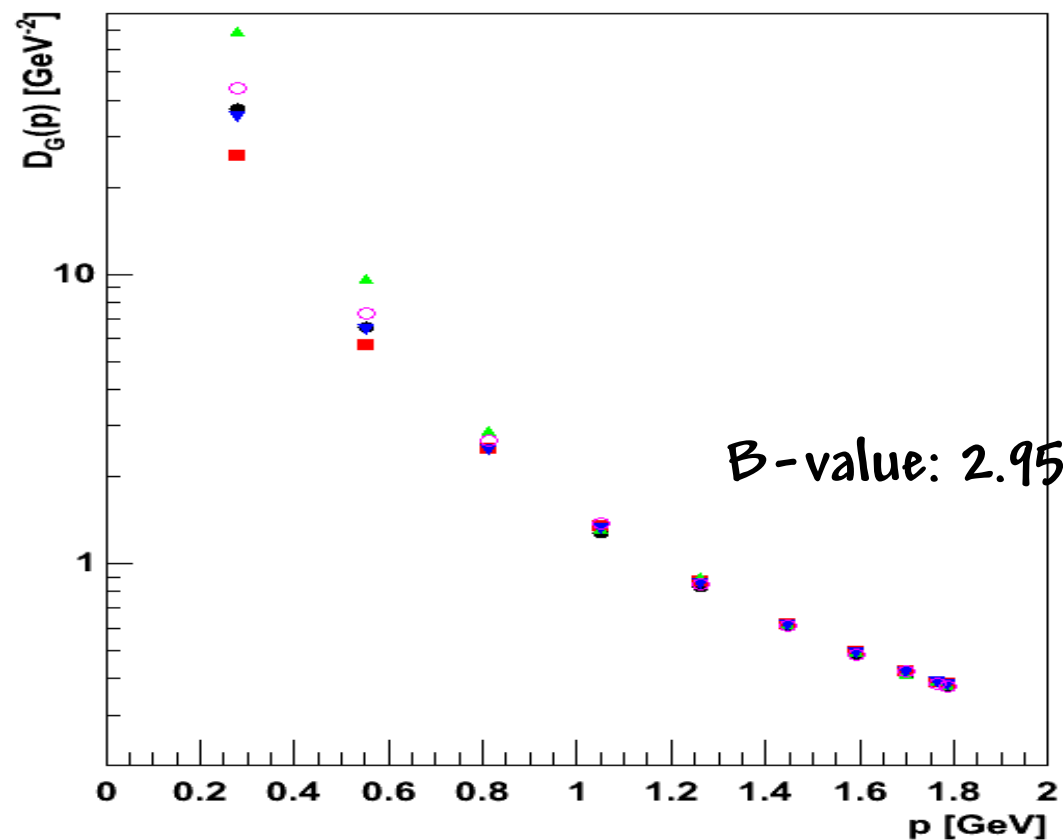
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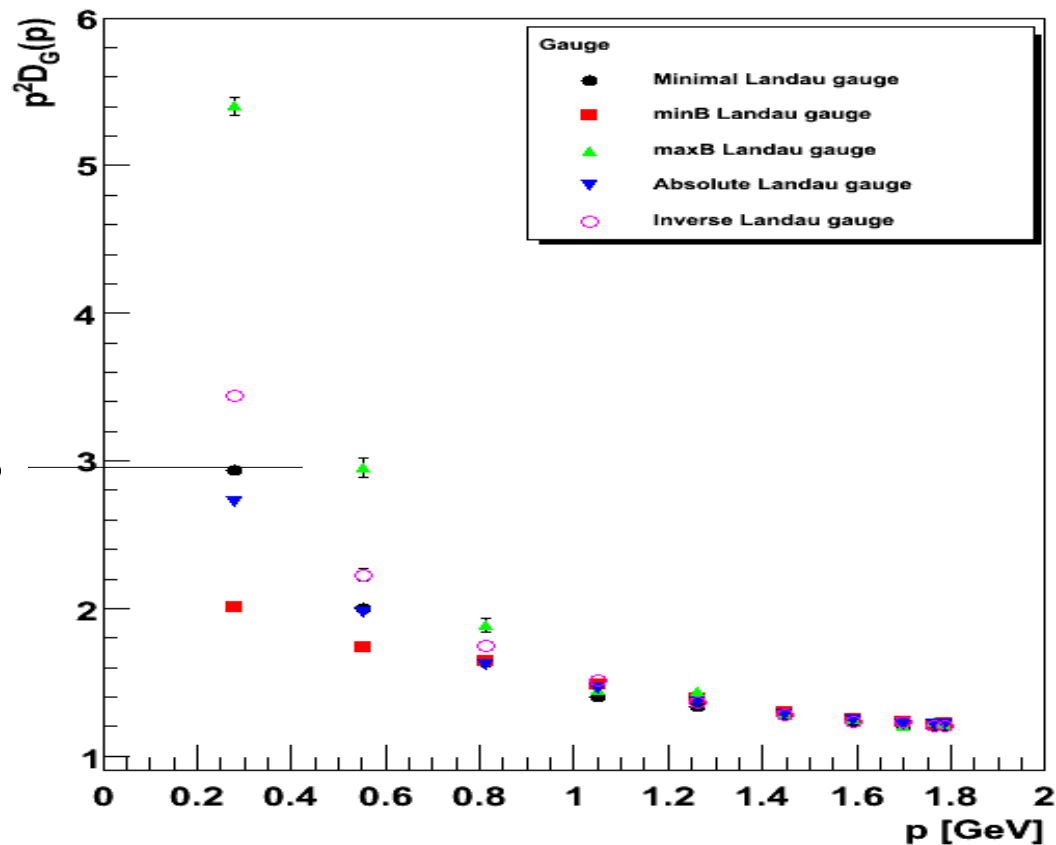
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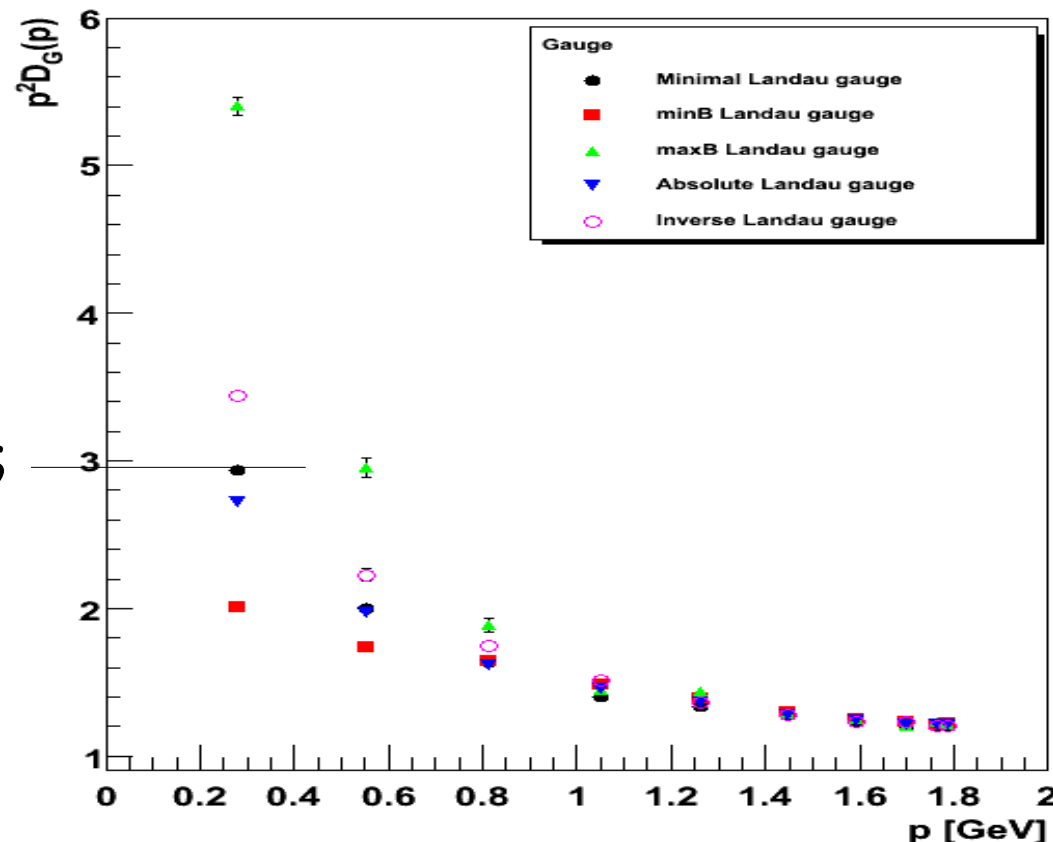
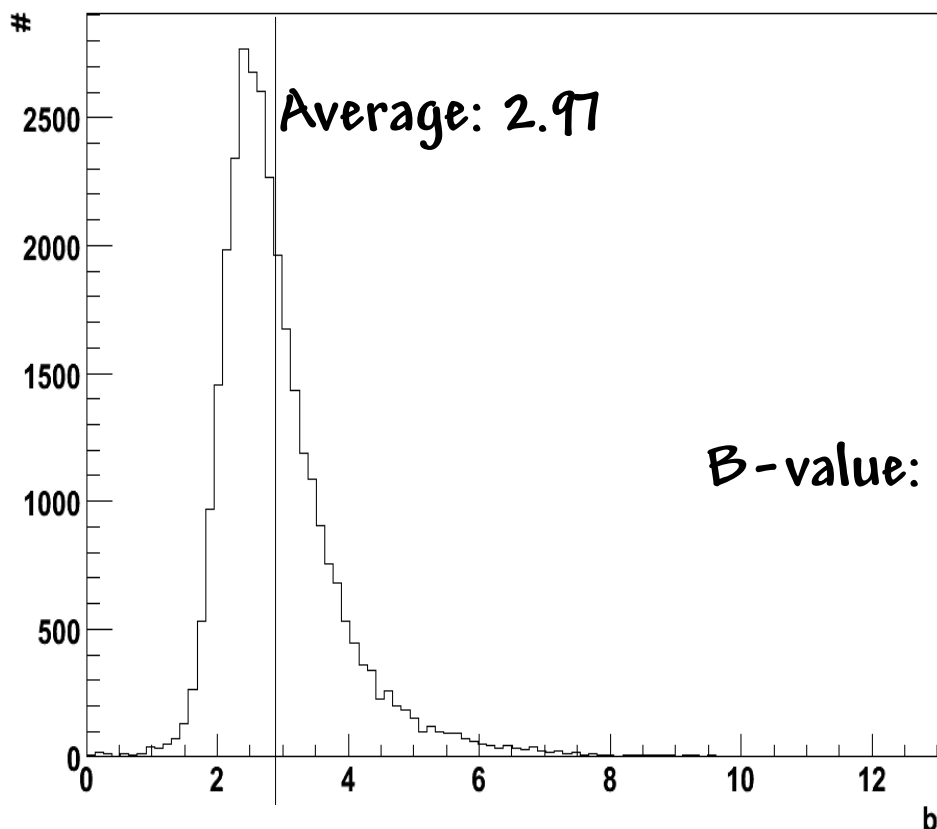
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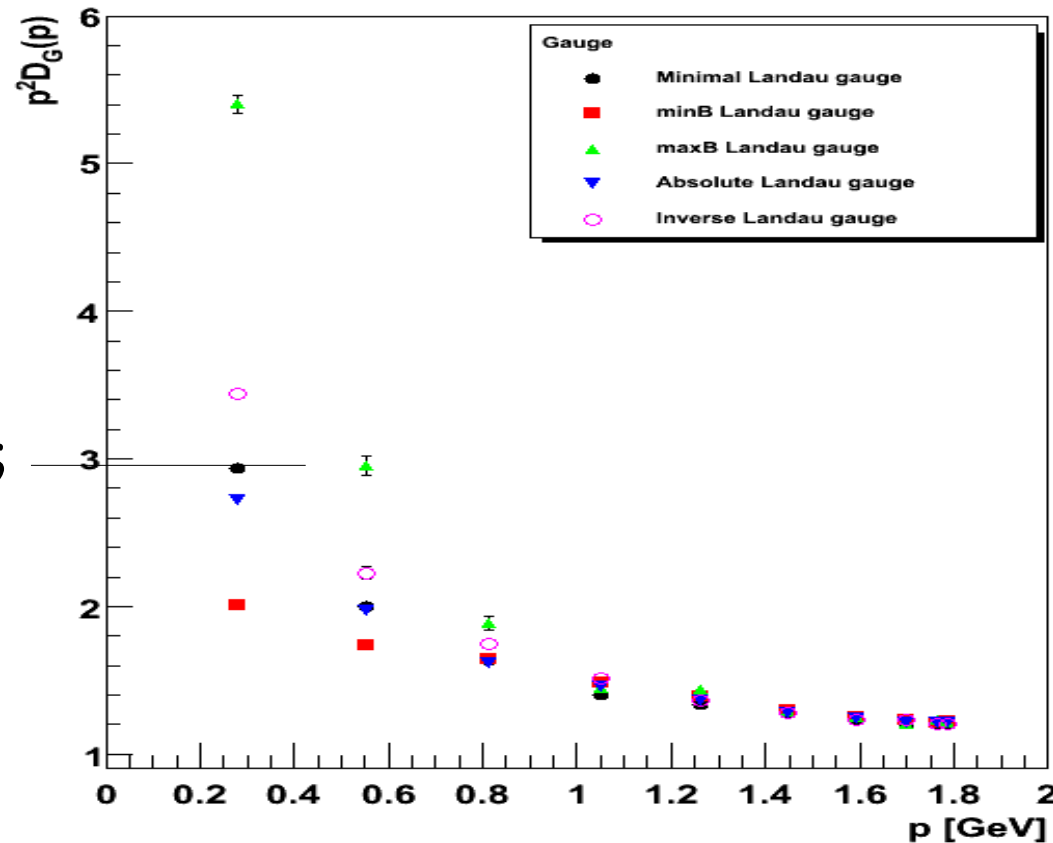
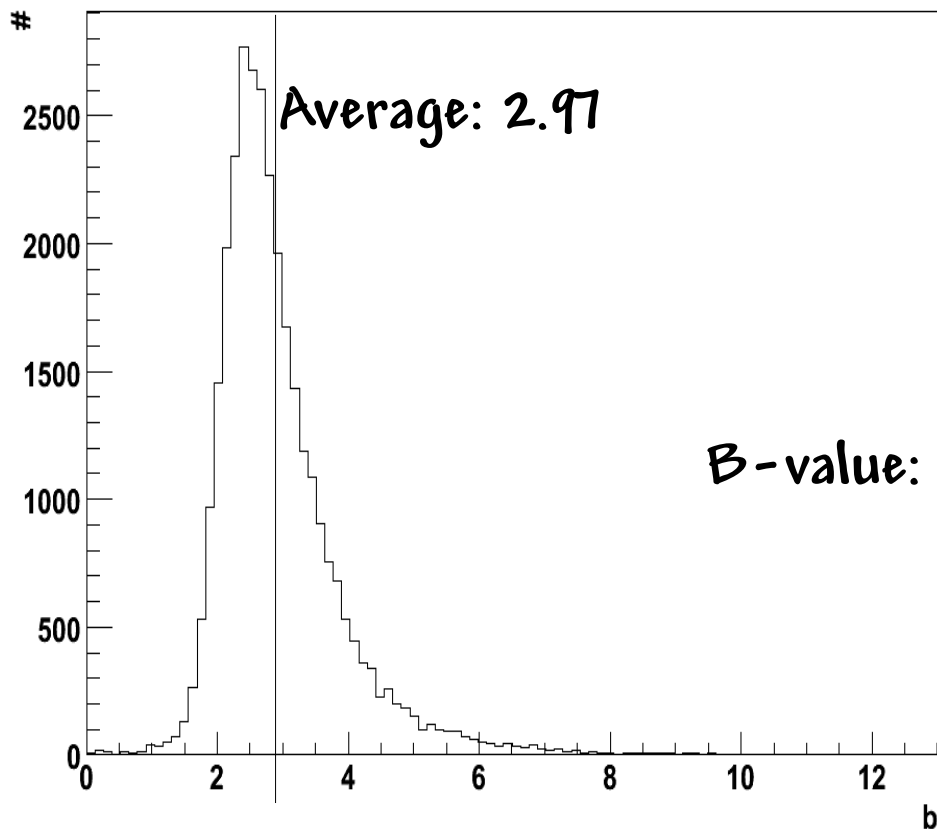


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 - Also for absolute or inverse Landau gauge, but with a deformed map
 - All Landau gauges should be map* B -gauge

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 - But how to identify the **first Gribov horizon** outside the **lattice**?
 - Positivity and monotonicity of the **ghost propagator**? Assume

[Reinhardt et al. 2008]

Methods

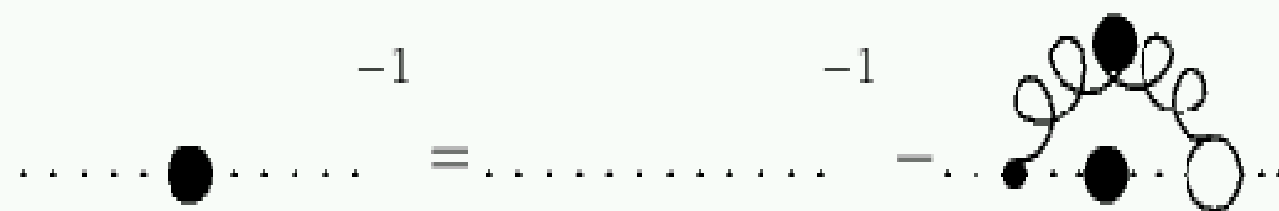
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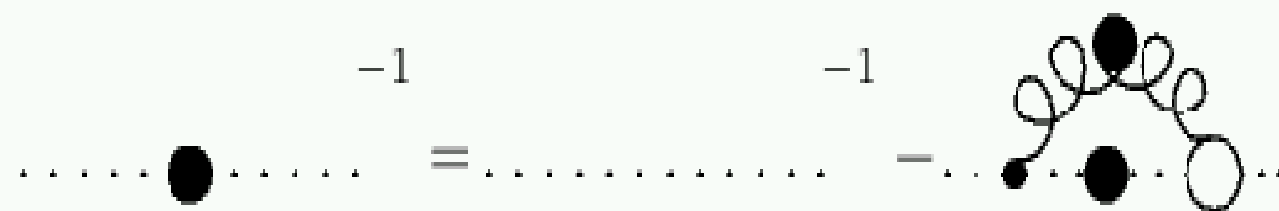
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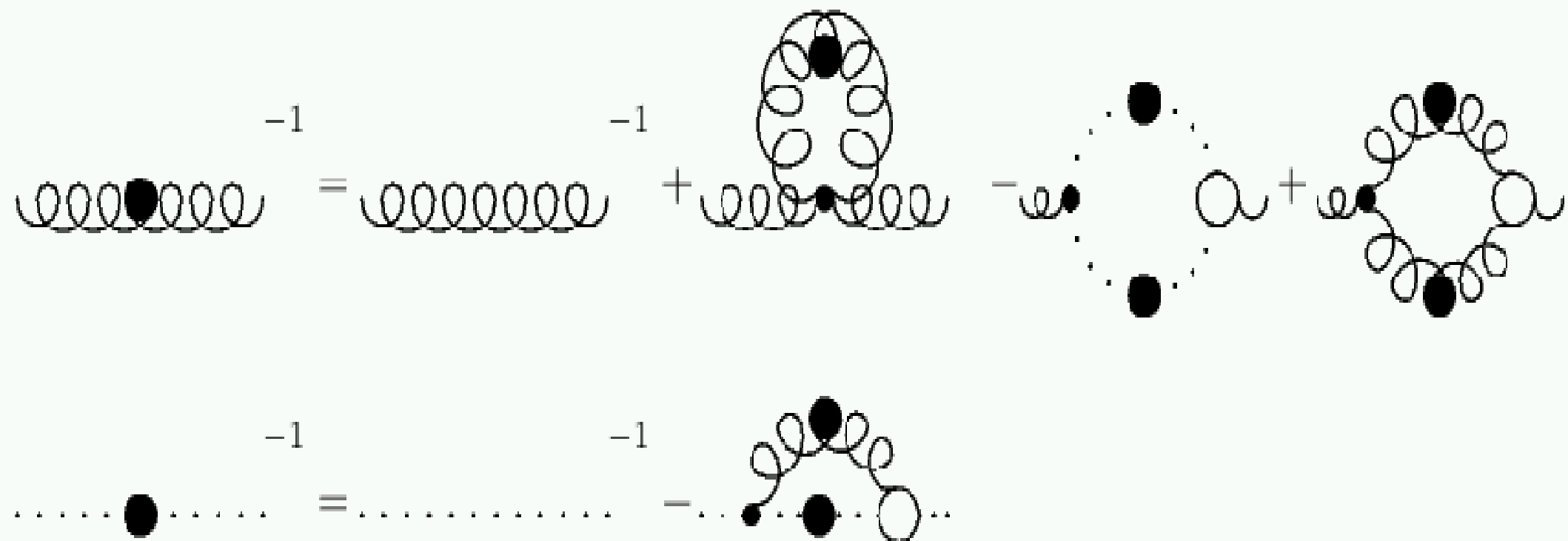
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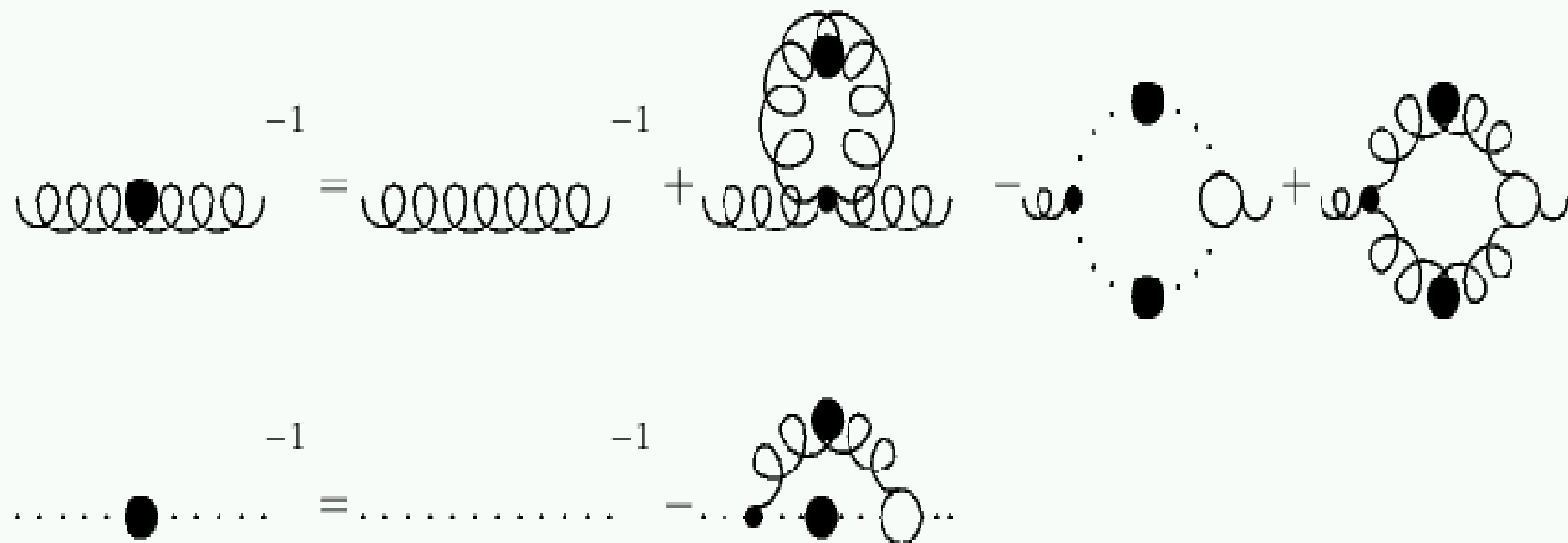
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- Consistent and self-consistent truncation schemes can be developed

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- Combination of all methods most successful!

Implementing \mathcal{B} Landau gauges in the continuum

- \mathcal{B} -Landau gauges are implemented via the **ghost dressing function equation**

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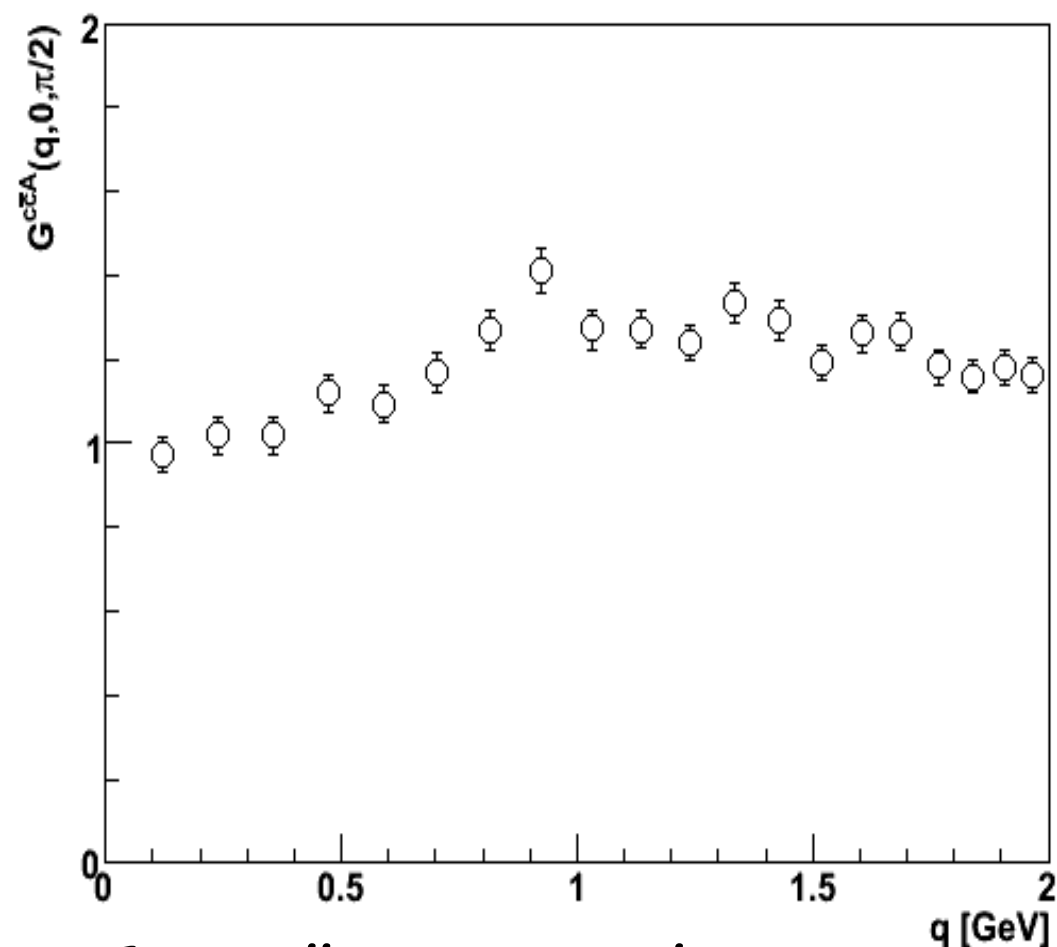
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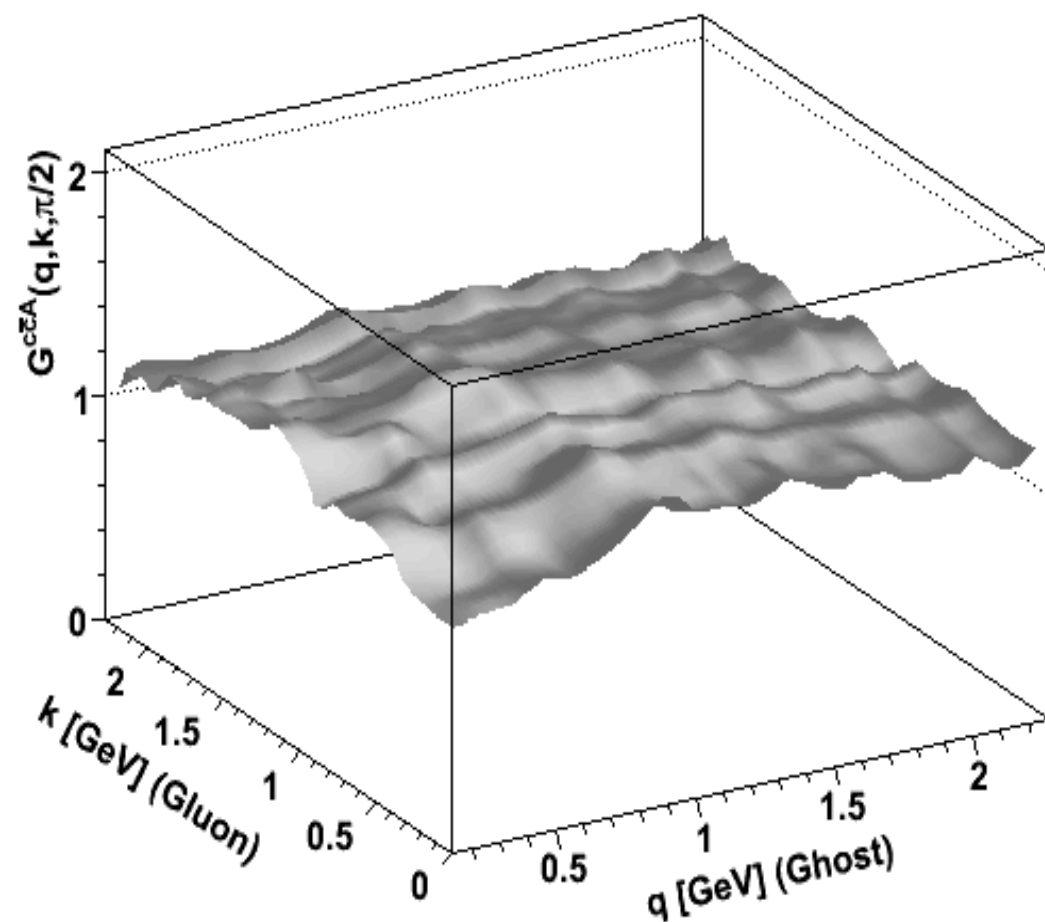
Ghost-gluon vertex in 3d

[60³, beta=4.2, Cucchieri et al., 2008]

Ghost-gluon vertex, one momentum vanishing



Ghost-gluon vertex, orthogonal momenta



- Essentially constant, only some structure at 1 GeV
- Same in 2d and 4d
- In agreement with DSE predictions [Schleifenbaum et al., PRD 2005]

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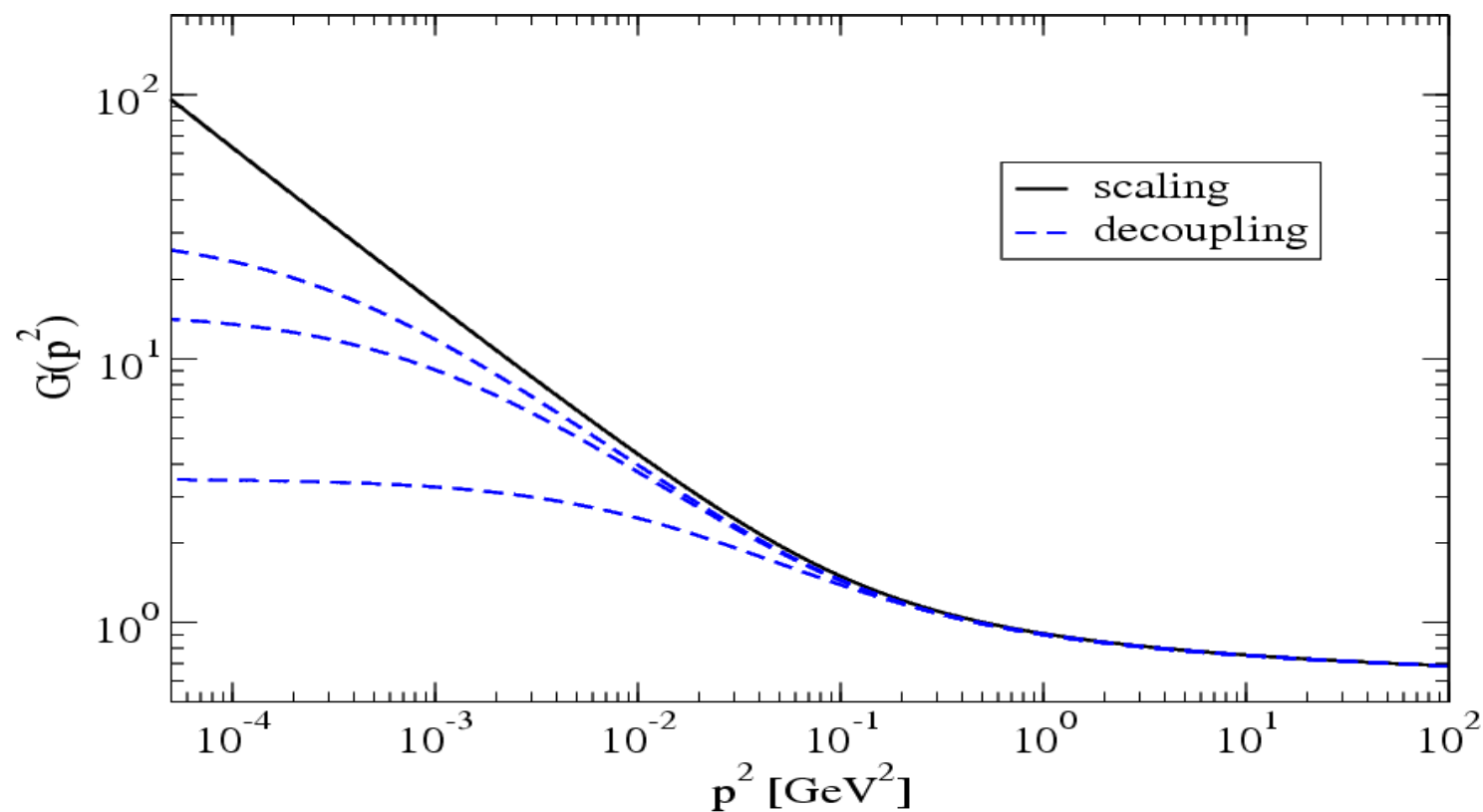
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- Select wave-function renormalization Z such as to implement a particular \mathcal{B} -value $\mathcal{B} = 1/(Z+A)$ and thus **\mathcal{B} -Landau gauge**

Ghost dressing function in the continuum [Fischer et al., 2008]

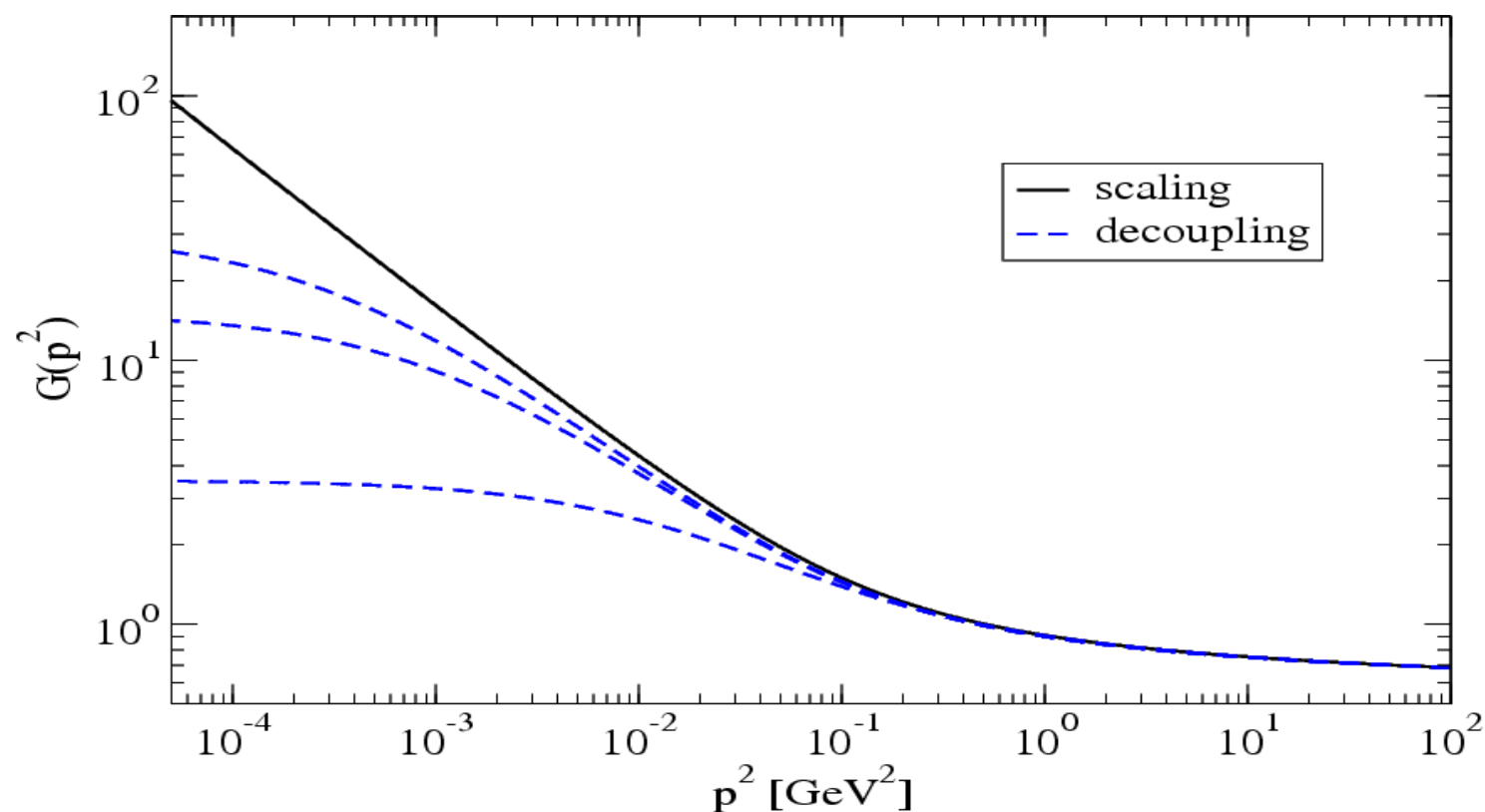
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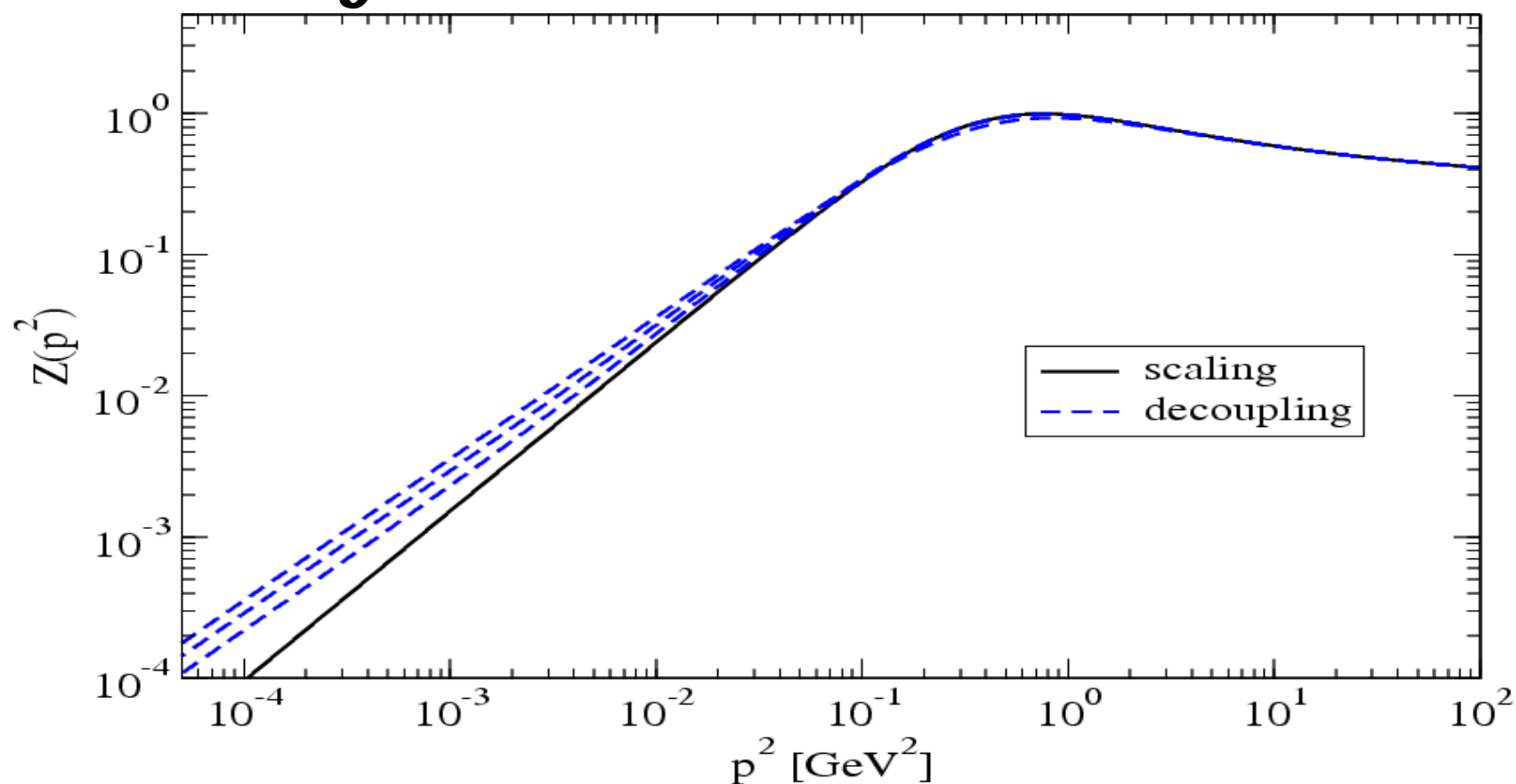
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 - Continuum implementation of the B -Landau gauges

Gluon dressing function in the continuum [Fischer et al., 2008]



- **Decoupling gauges** yield a decoupling (infrared massive) **gluon propagator** - consistent with the lattice result
- The **scaling gauge** yields an infrared vanishing **gluon propagator**

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 - This needs a more formal formulation, volume studies...etc.

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- Can be expanded to the case with matter fields [Alkofer et al., 2007/8]

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- However: Yet unknown, to which **Landau-B gauge** this corresponds, or if it applies generally to all **Landau gauges**
 - Will require always large (not infinitesimal) transformations

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- However: Yet unknown, to which **Landau-B gauge** this corresponds, or if it applies generally to all **Landau gauges**
 - Will require always large (not infinitesimal) transformations
- Assume it to be valid in the **scaling gauge**

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 - Very attractive, if correct

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- Additional freedom: Select a **gauge** which suits the needs
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- But a formal formulation and more details are still required...